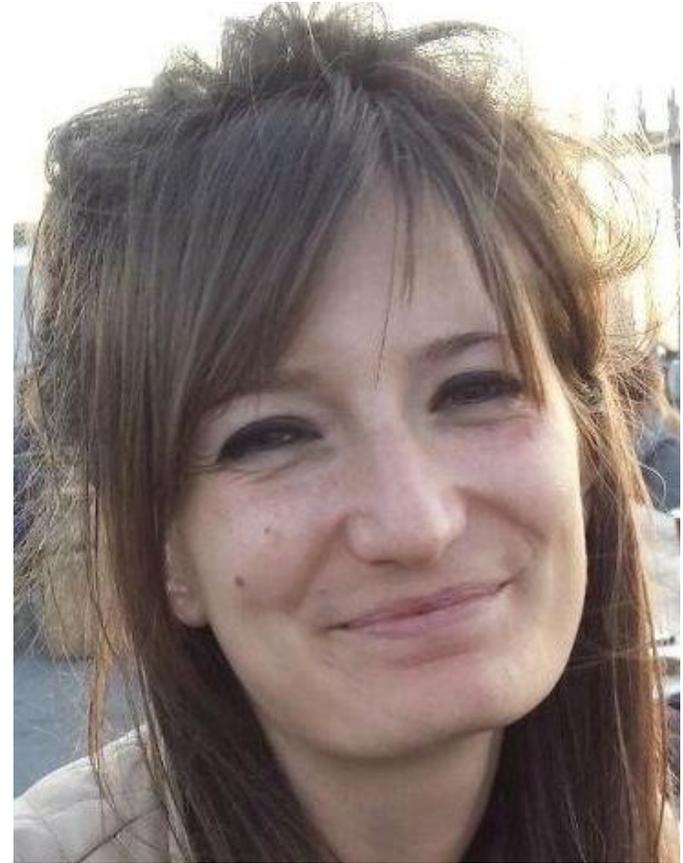


# Field Patterns: A new type of Wave

Ornella Mattei and Graeme Milton,

Department of Mathematics, The University of Utah





# Formulation of the problem

- Generic wave equation:

$$\frac{\partial}{\partial x} \left( \alpha(x, t) \frac{\partial u(x, t)}{\partial x} \right) - \frac{\partial}{\partial t} \left( \beta(x, t) \frac{\partial u(x, t)}{\partial t} \right) = 0$$

The coefficients are **time – dependent** → **DYNAMIC MATERIALS**

- **Boundary conditions:** The medium is infinite in the  $x$ -direction
- **Initial conditions:**

$$u(x, 0) = g(x)$$
$$\frac{\partial u(x, t)}{\partial t} \Big|_{t=0} = f(x)$$

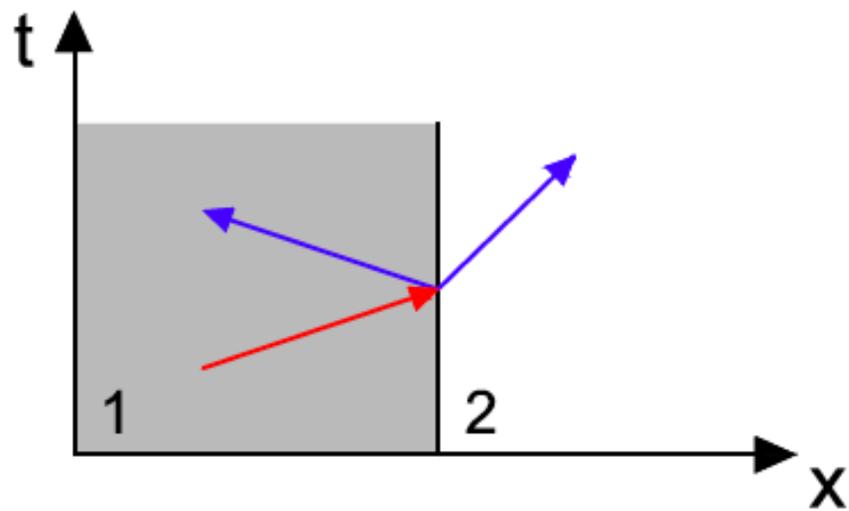
[see, e.g., Lurie, 2007]

# Realization of dynamic materials

- Liquid crystals
- Ferroelectric, ferromagnetic materials
- Pump wave + small amplitudes waves [e.g. Louisell & Quate (1958)]
- Transmission line with modulated inductance [e.g. Cullen (1958)]
- Experiments and more references in [Honey & Jones (1958)]
- Dynamic modulation + photons [e.g. Fang et al. (2012), Boada et al. (2012), Celi et al. (2014), Yuan et al. (2016)]
- ...
- Walking droplets [e.g. Couder et al. (2005), Couder & Fort (2006), Bush (2015)]
- Time reversal [e.g. Fink (2016), Goussev et al. (2016)]

# Dynamic composites

Pure space interface

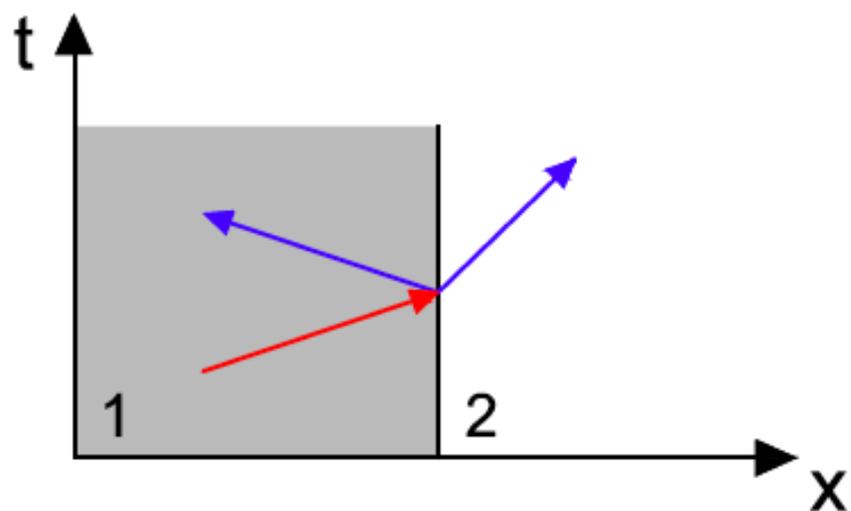


Pure time interface

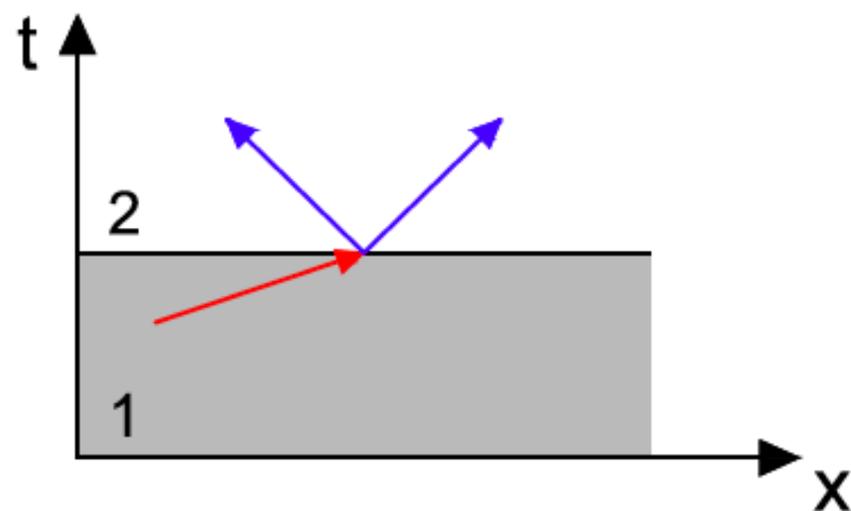


# Dynamic composites

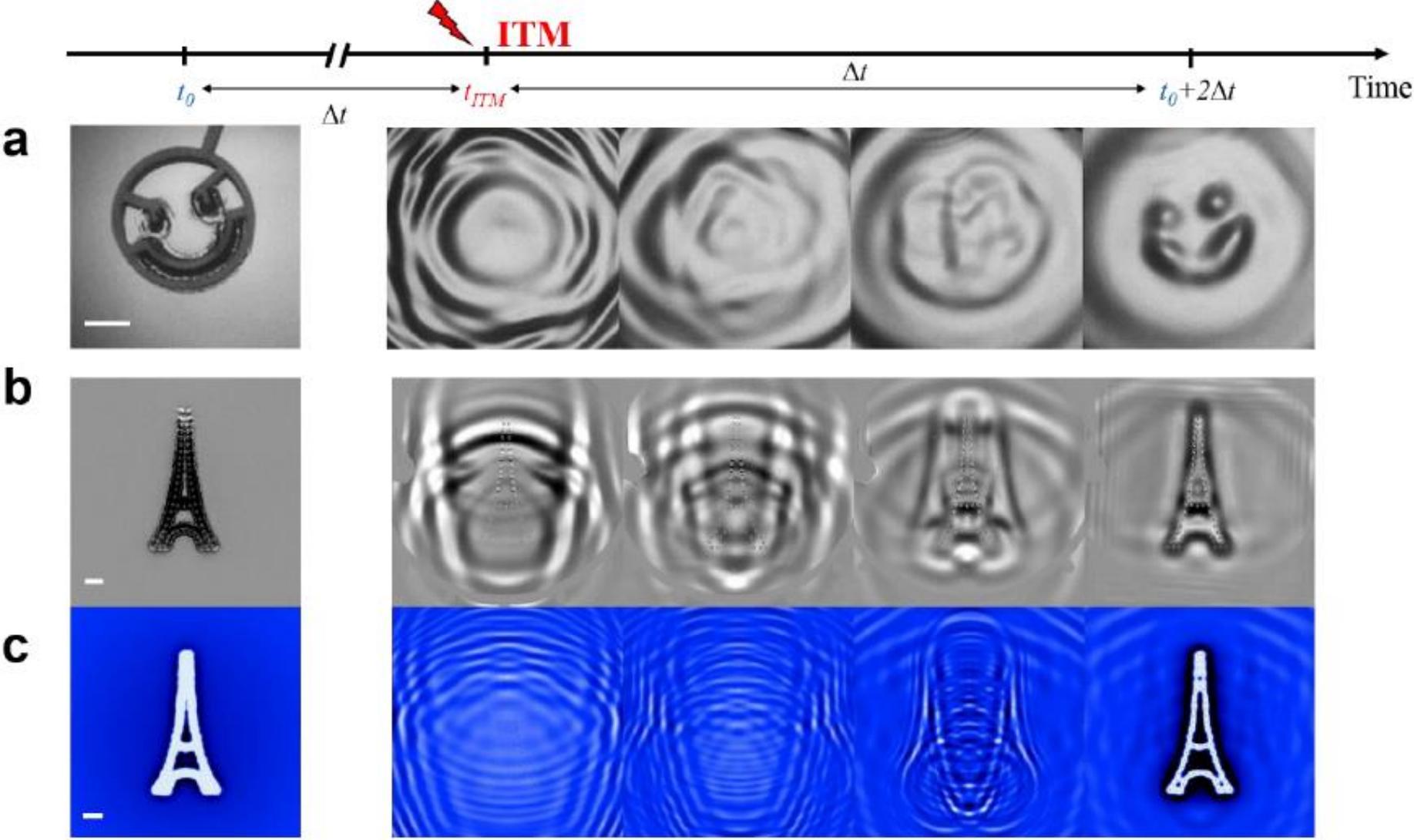
Pure space interface



Pure time interface



# What happens at a time interface?



# Thinking of the wave equation as a conductivity problem

$$\mathbf{j}(\mathbf{x}) = \boldsymbol{\sigma}(\mathbf{x})\mathbf{e}(\mathbf{x}), \quad \text{where } \nabla \cdot \mathbf{j} = 0, \quad \mathbf{e} = -\nabla V,$$

$$\boldsymbol{\sigma}(\mathbf{x}) = \begin{pmatrix} \alpha(\mathbf{x}) & 0 \\ 0 & -\beta(\mathbf{x}) \end{pmatrix}, \quad \begin{array}{ll} \text{material 1} & \rightarrow \alpha_1, \beta_1 \\ \text{material 2} & \rightarrow \alpha_2, \beta_2 \end{array}$$

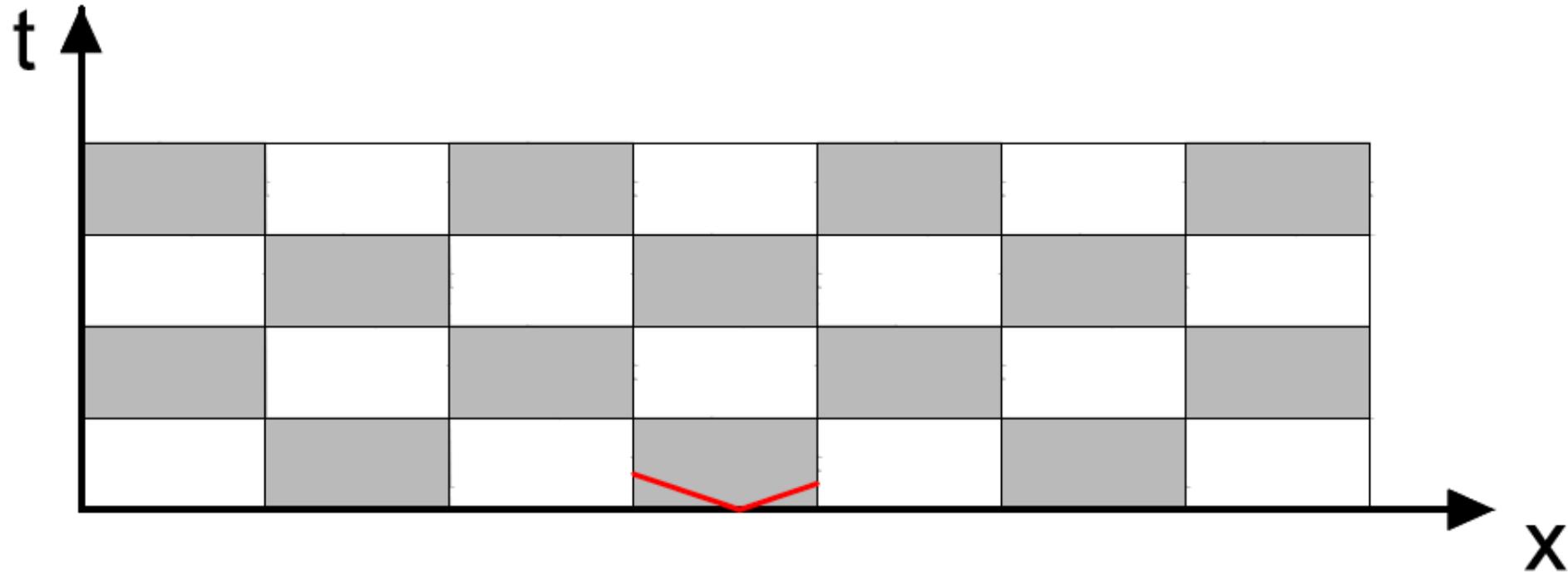
$$\frac{\partial}{\partial x_1} \left( \alpha(x_1, x_2) \frac{\partial V(x_1, x_2)}{\partial x_1} \right) - \frac{\partial}{\partial x_2} \left( \beta(x_1, x_2) \frac{\partial V(x_1, x_2)}{\partial x_2} \right) = 0$$

N.B. Hyperbolic materials!! [See, e.g. the review Poddubny, Iorsh, Belov, Kivshar, 2013]

$$\alpha_i \frac{\partial^2 V_i}{\partial x^2} - \beta_i \frac{\partial^2 V_i}{\partial t^2} = 0, \quad i = 1, 2$$

D'Alembert solution:  $V_i(x, t) = V_i^+(x - c_i t) + V_i^-(x + c_i t) \quad c_i = \sqrt{\frac{\alpha_i}{\beta_i}}$

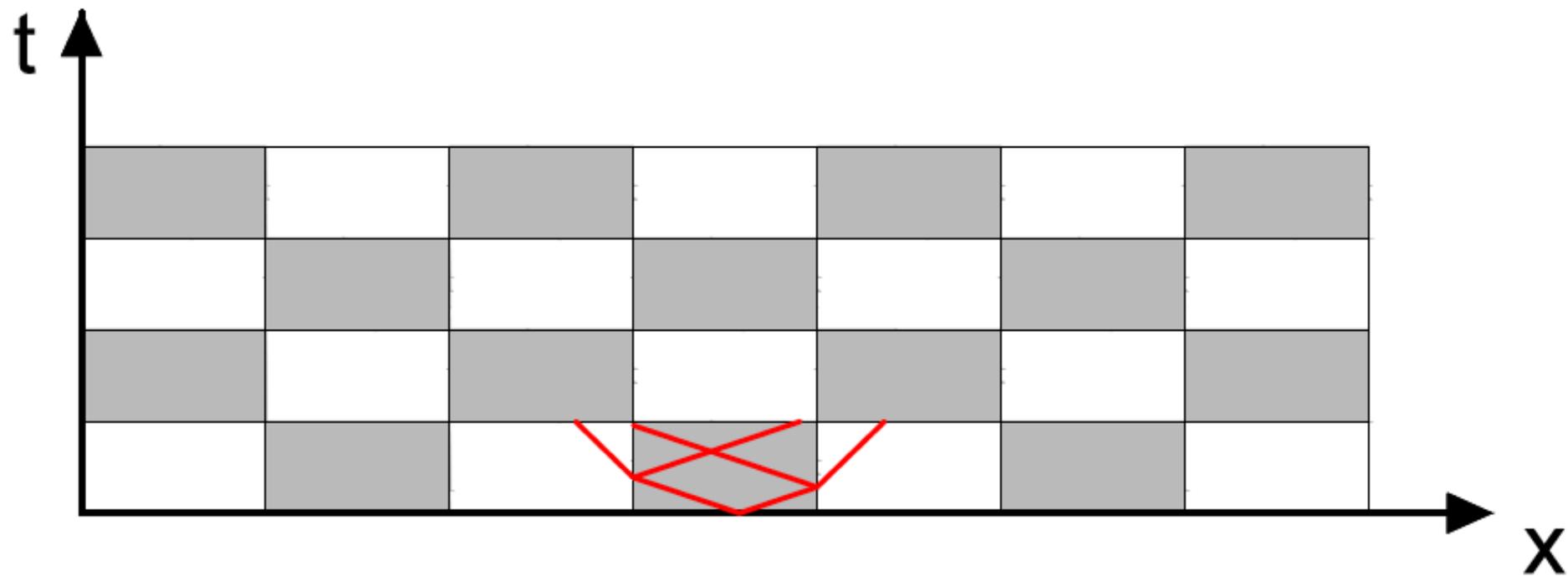
# Evolution of a disturbance in a space-time checkerboard



Transmission conditions:

$$\begin{cases} V_1 = V_2 \\ \mathbf{n} \cdot \boldsymbol{\sigma}_1 \nabla V_1 = \mathbf{n} \cdot \boldsymbol{\sigma}_2 \nabla V_2 \end{cases}$$

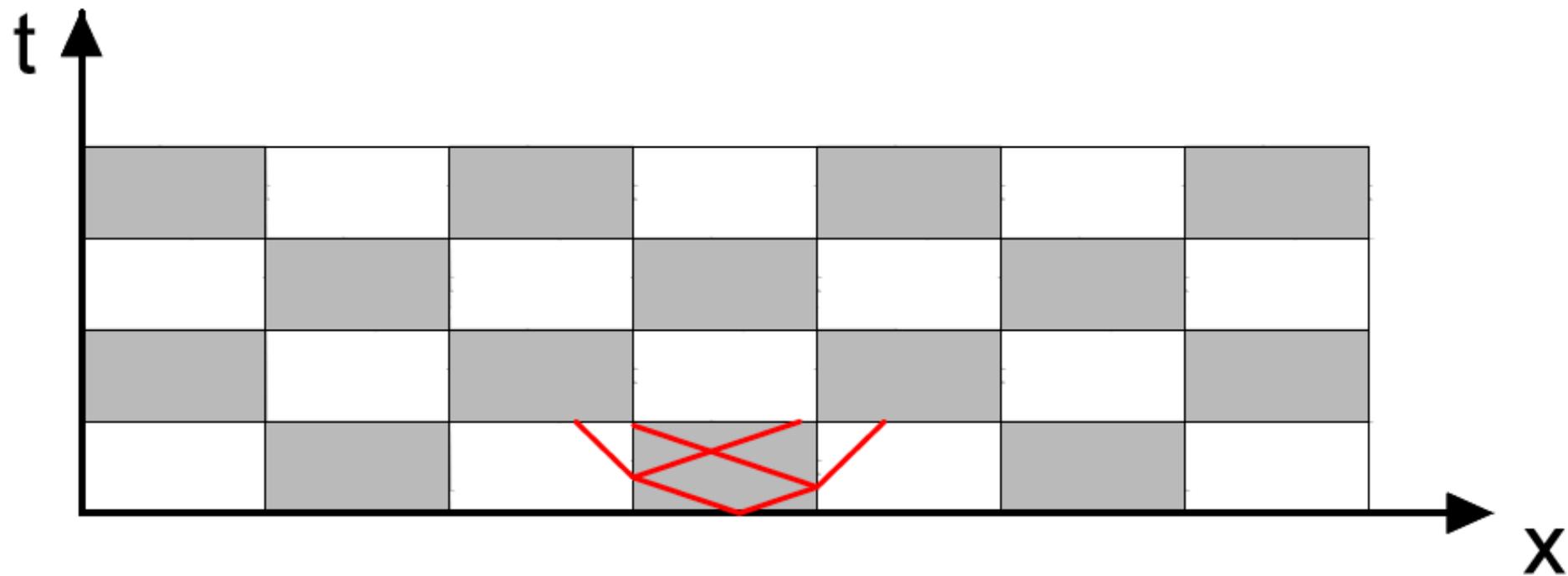
# Evolution of a disturbance in a space-time checkerboard



Transmission conditions:

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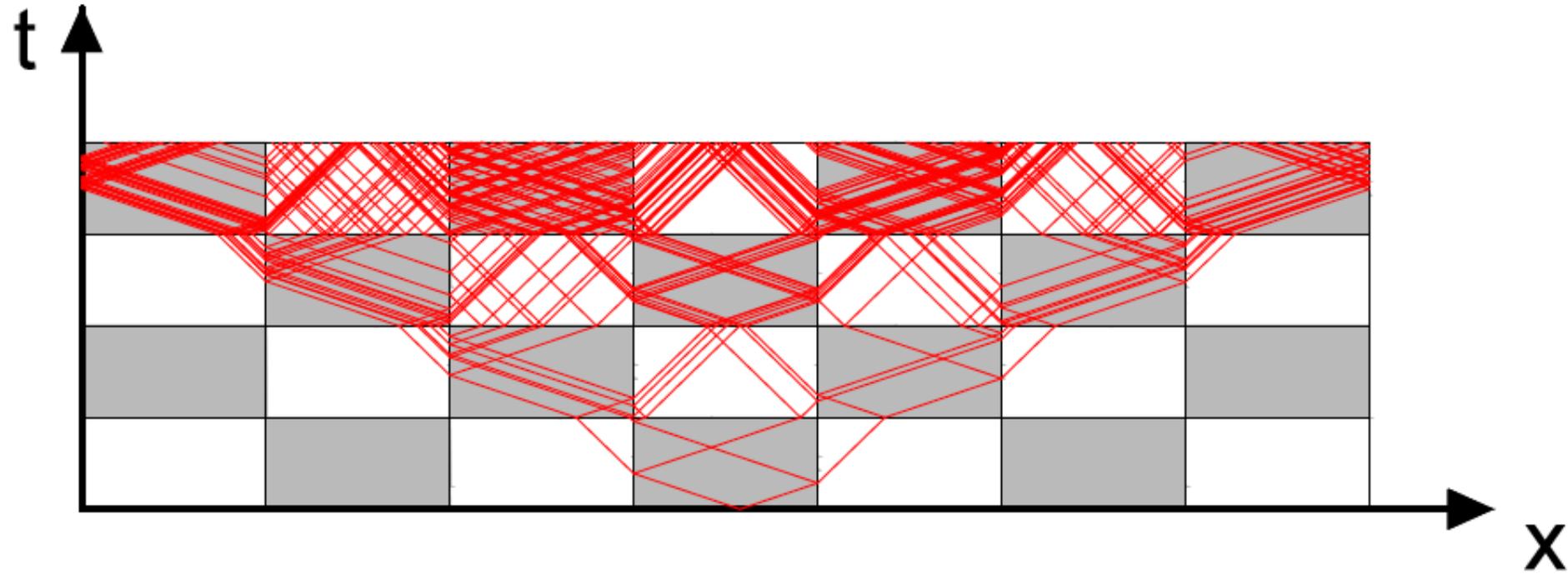
# Evolution of a disturbance in a space-time checkerboard



Transmission conditions:

$$\begin{cases} V_1 = V_2 \\ \mathbf{n} \cdot \boldsymbol{\sigma}_1 \nabla V_1 = \mathbf{n} \cdot \boldsymbol{\sigma}_2 \nabla V_2 \end{cases}$$

# Evolution of a disturbance in a space-time checkerboard

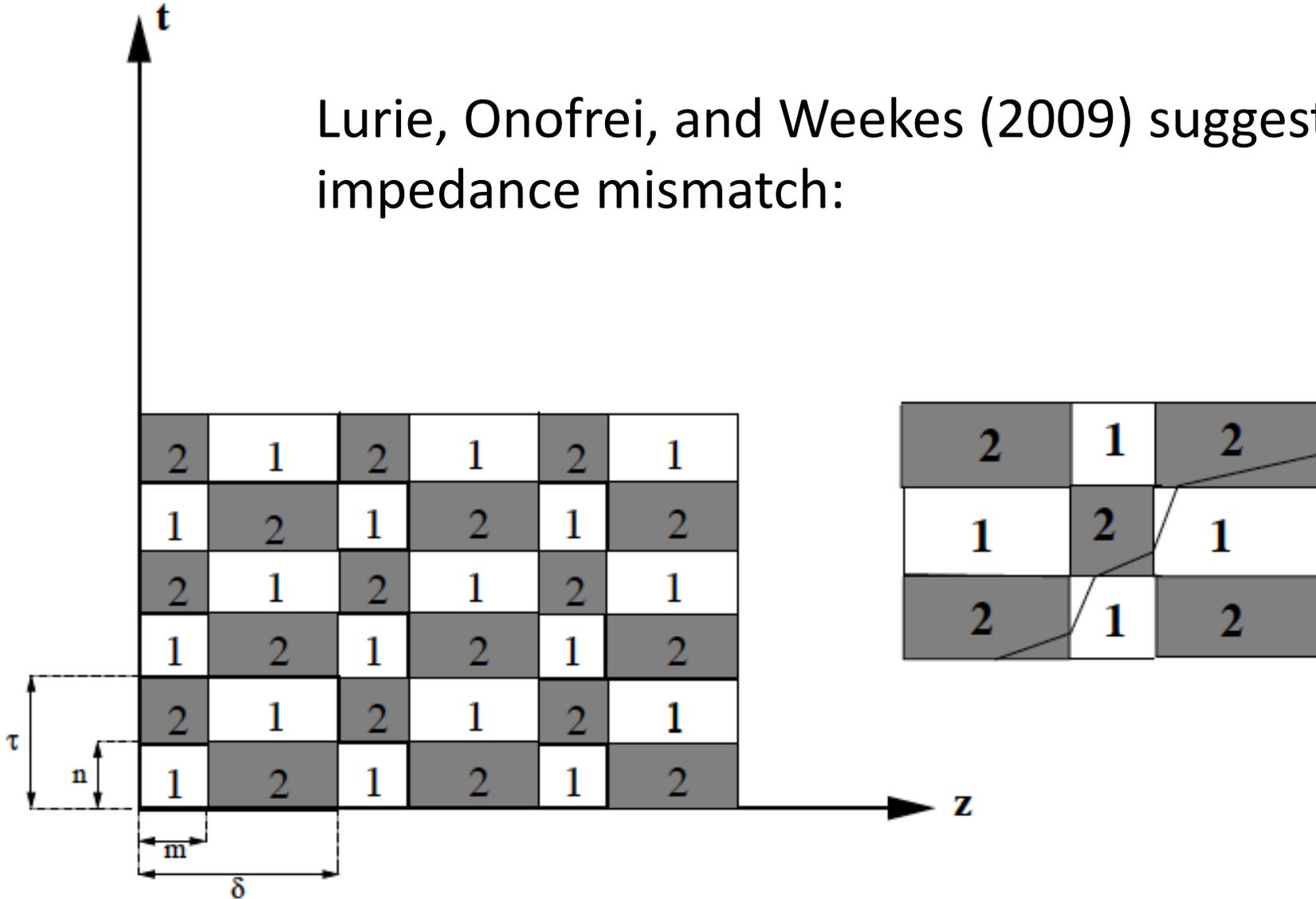


Transmission conditions:

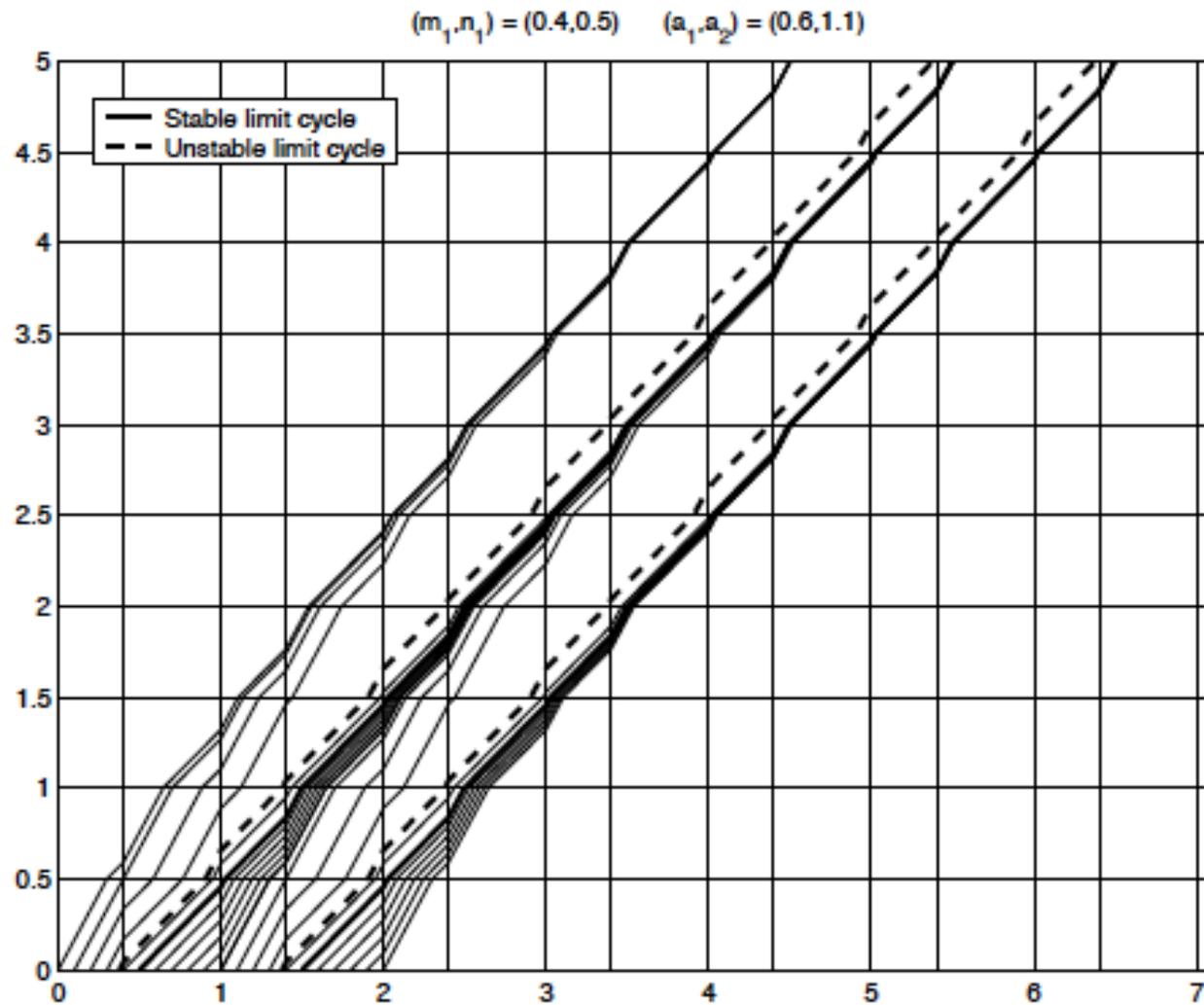
$$\begin{cases} V_1 = V_2 \\ \mathbf{n} \cdot \boldsymbol{\sigma}_1 \nabla V_1 = \mathbf{n} \cdot \boldsymbol{\sigma}_2 \nabla V_2 \end{cases}$$

# How to avoid this complicated cascade?

Lurie, Onofrei, and Weekes (2009) suggested having a zero impedance mismatch:



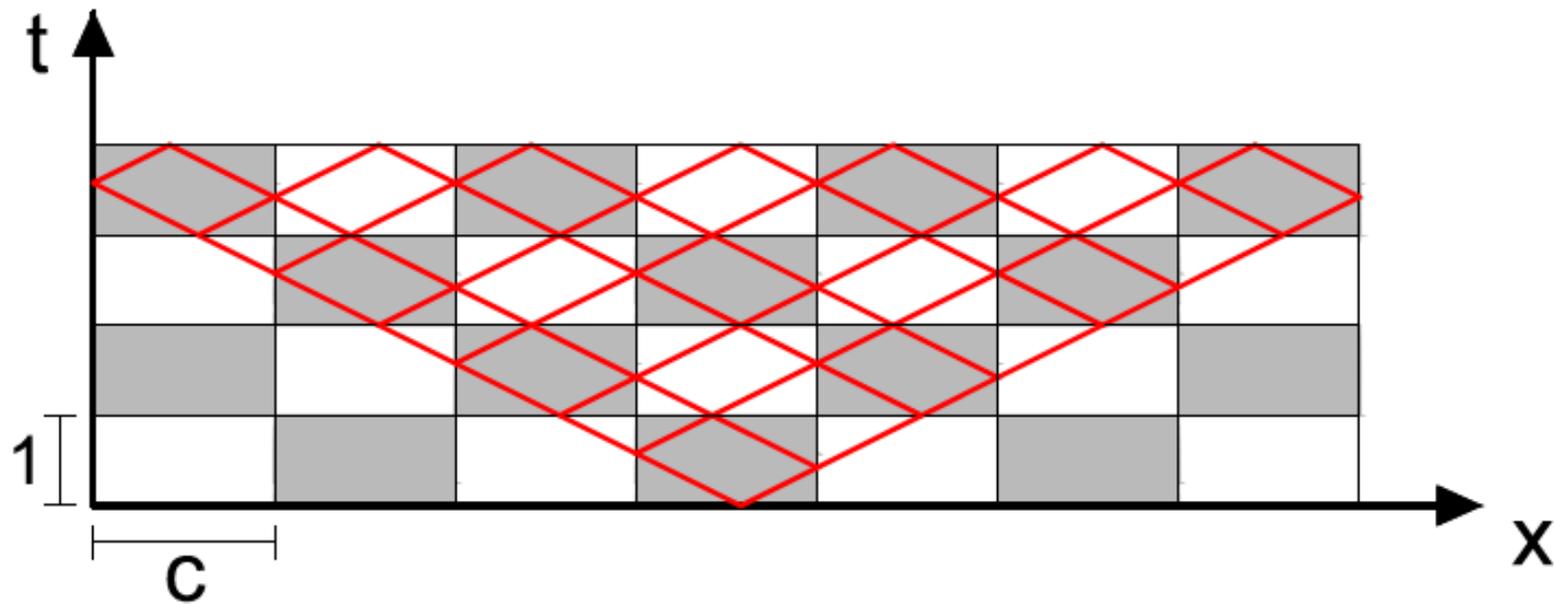
Curiously they found accumulations of the characteristic lines:



A bit like a shock but in a linear medium!

# Field patterns in a space-time checkerboard

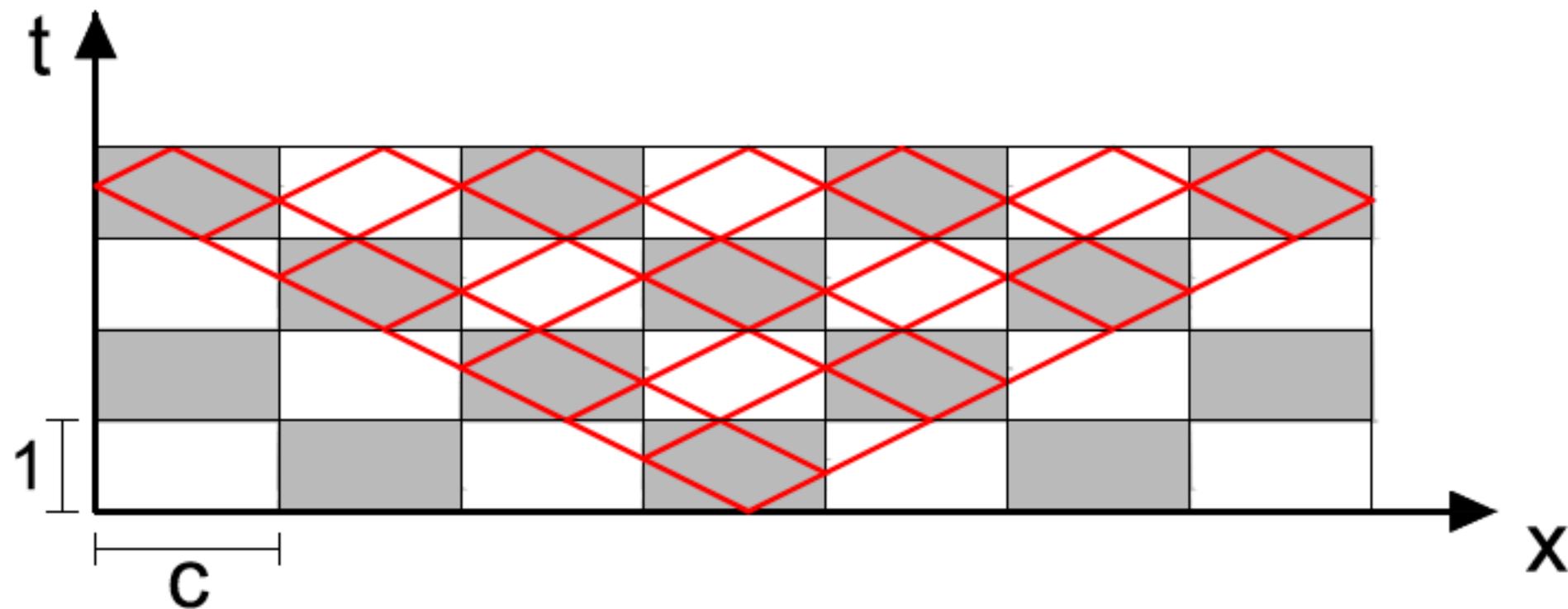
$$c_1 = c_2 = c \quad \Rightarrow \quad \frac{\alpha_1}{\beta_1} = \frac{\alpha_2}{\beta_2}$$



Field patterns are a new type of wave propagating along orderly patterns of characteristic lines!!!

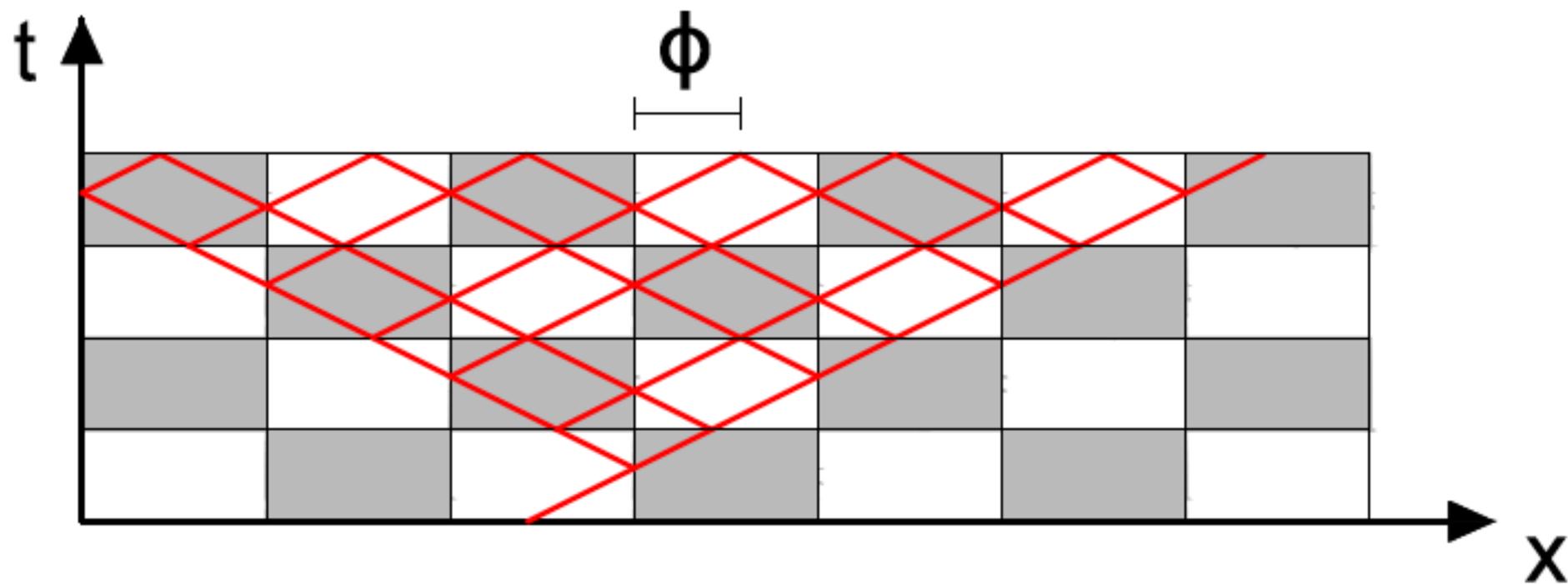
# Field patterns in a space-time checkerboard

$$c_1 = c_2 = c \quad \Rightarrow \quad \frac{\alpha_1}{\beta_1} = \frac{\alpha_2}{\beta_2}$$



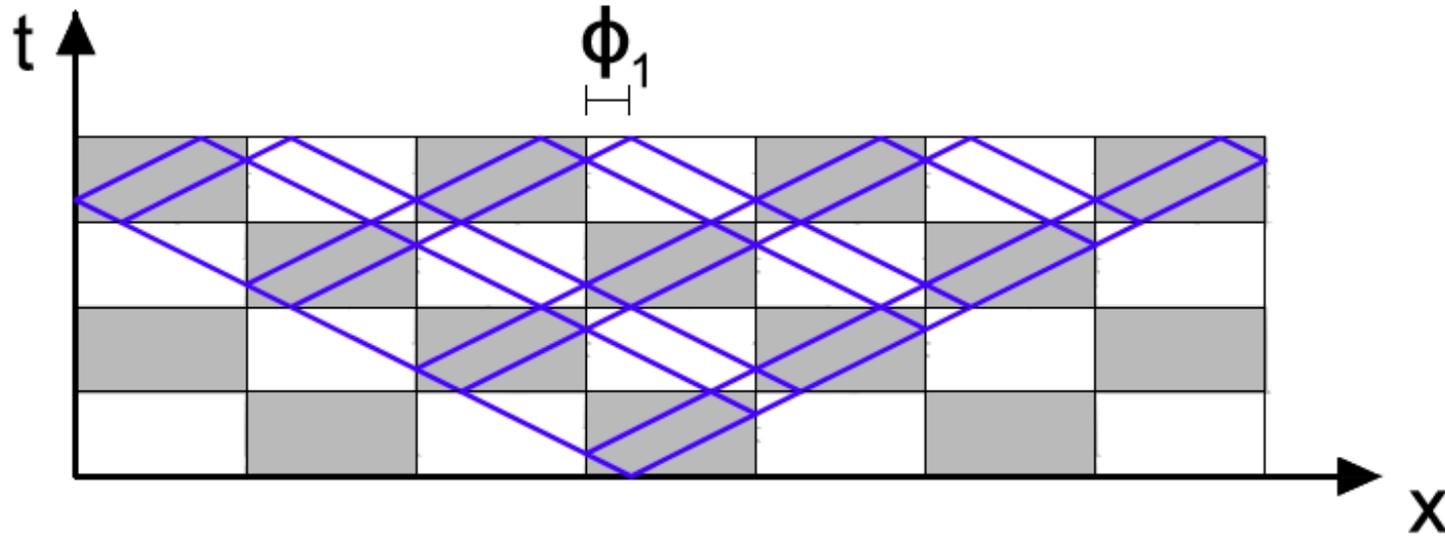
# Families of field patterns

$$c_1 = c_2 = c \quad \Rightarrow \quad \frac{\alpha_1}{\beta_1} = \frac{\alpha_2}{\beta_2}$$



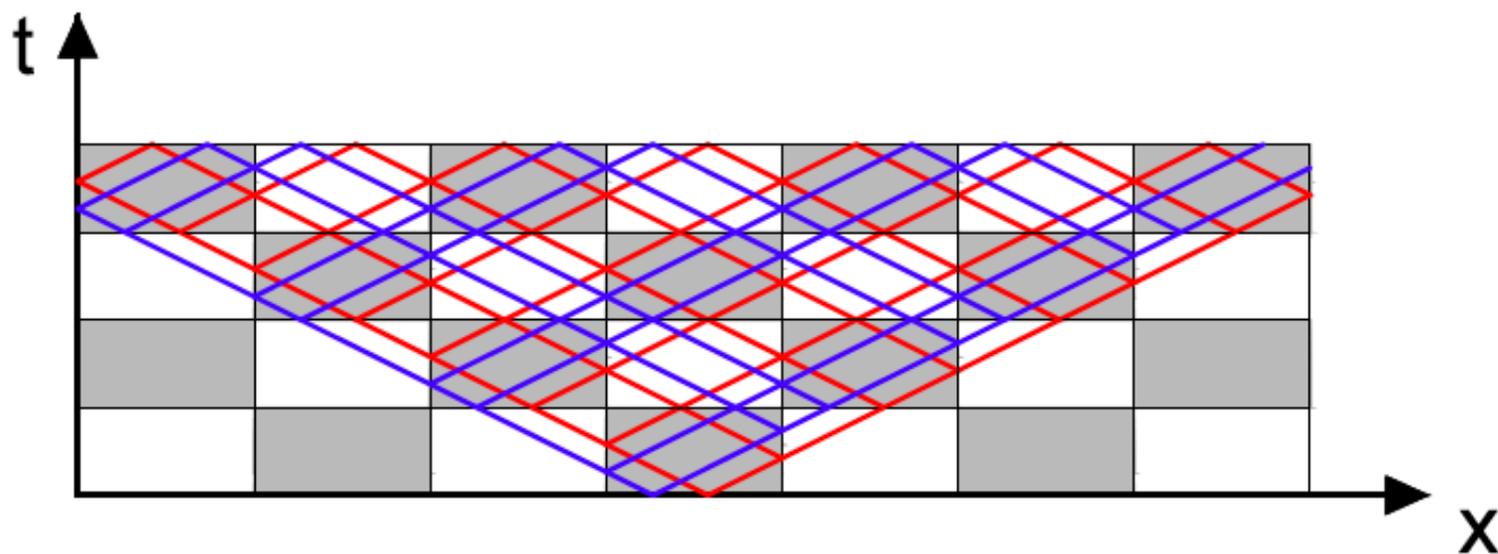
# Families of field patterns

$$c_1 = c_2 = c \quad \Rightarrow \quad \frac{\alpha_1}{\beta_1} = \frac{\alpha_2}{\beta_2}$$



Field patterns are a new type of wave propagating along orderly patterns of characteristic lines which arise in specific space-time microstructures whose geometry in one spatial dimension plus time is somehow commensurate to the slope of the characteristic lines.

## Multidimensional nature of field patterns

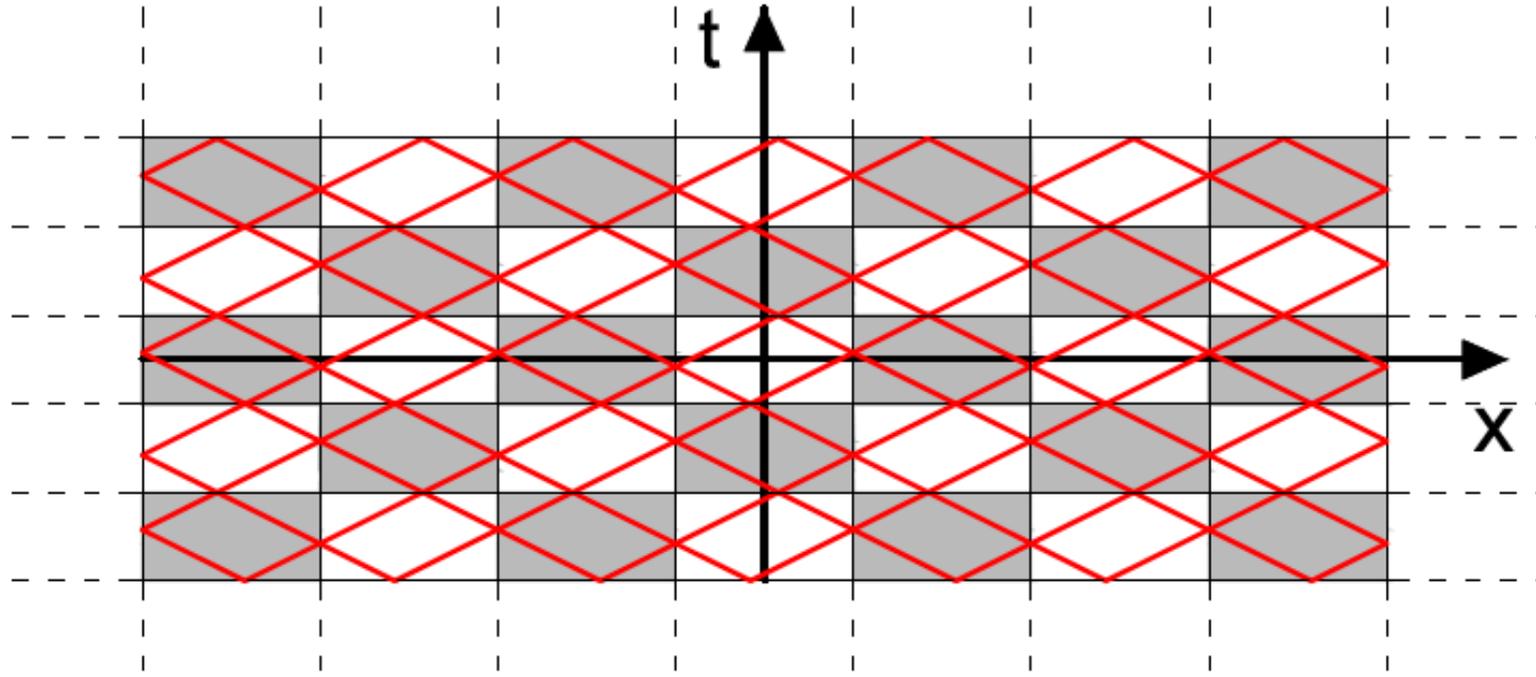


$$V(x, t) = \sum_{i=1}^m V_{\phi_i}(x, t)$$

Multidimensional space:  $V(x_1, x_2, \dots, x_m) = \sum_{i=1}^m V_{\phi_i}(x_i, t)$

Multicomponent potential:  $\mathbf{V}(x, t)$

# PT-symmetry of field patterns

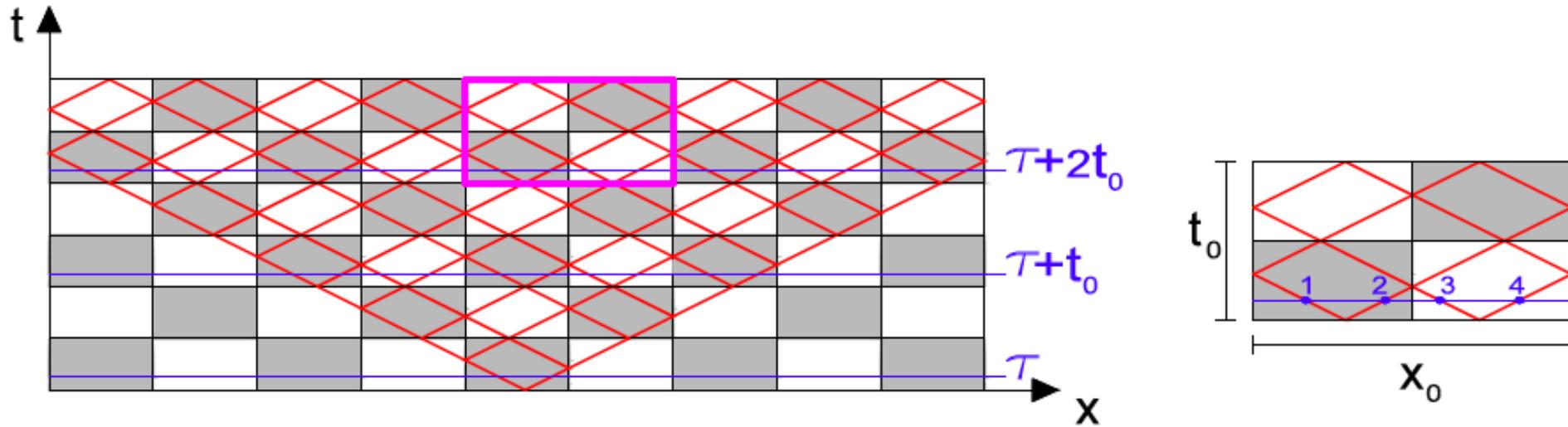


[Quantum physics, e.g., Bender and Boettcher, 1998,  
Optics, e.g., Zyablovsky et al., 2014]

Unbroken PT-symmetry  $\rightarrow$  real eigenvalues

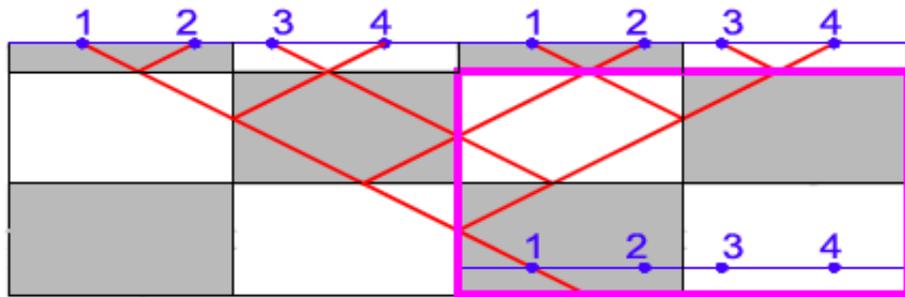
Broken PT-symmetry  $\rightarrow$  complex conjugate eigenvalues

# The transfer matrix



$$j(k, m, n + 1) = \sum_{k', m'} T_{(k, m), (k', m')} j(k', m', n)$$

$$T_{(k, m), (k', m')} = G_{k, k'}(m - m')$$

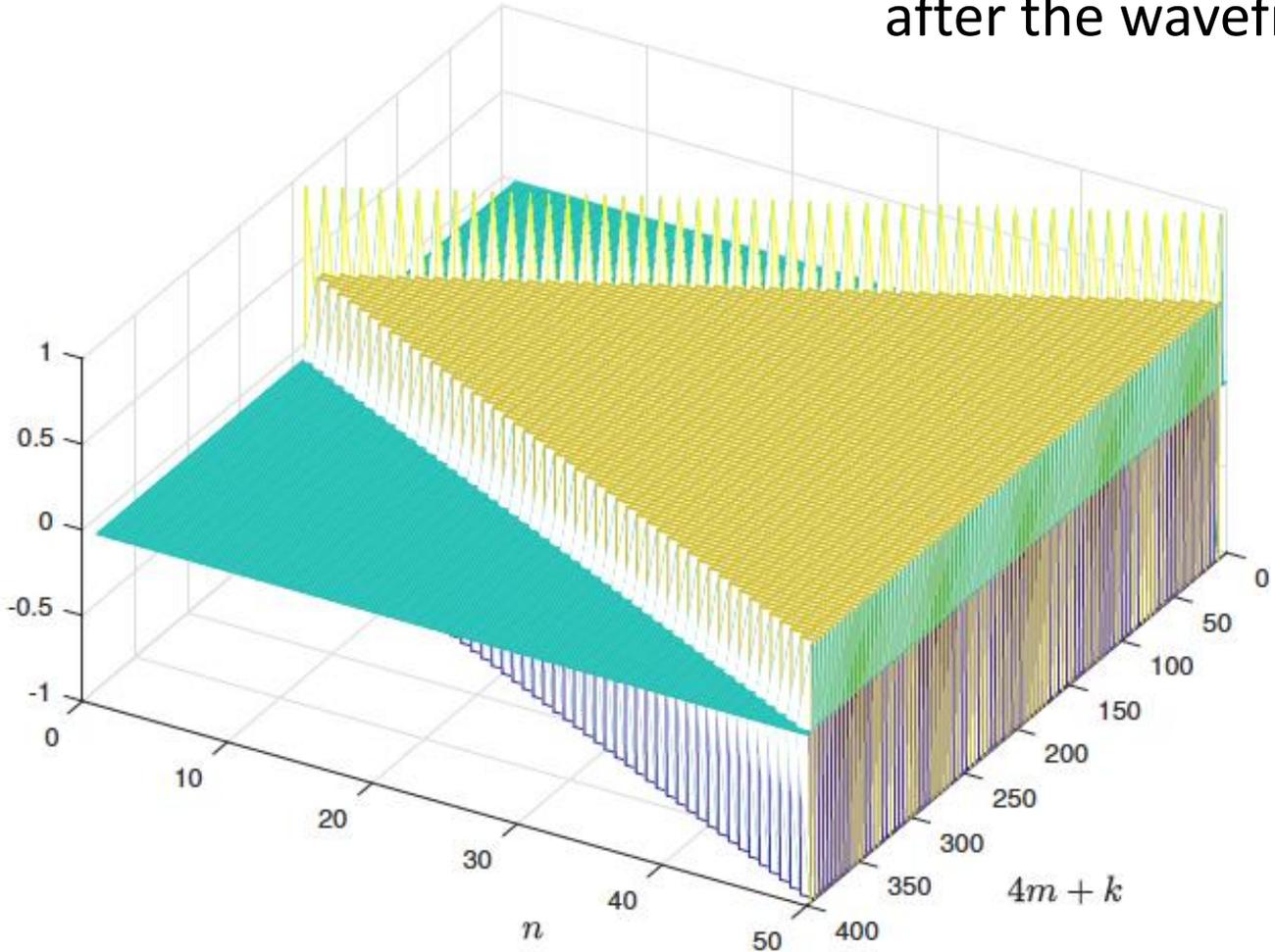


$T$  depends only on  $\gamma_i$

$T$  is  $PT$ -symmetric

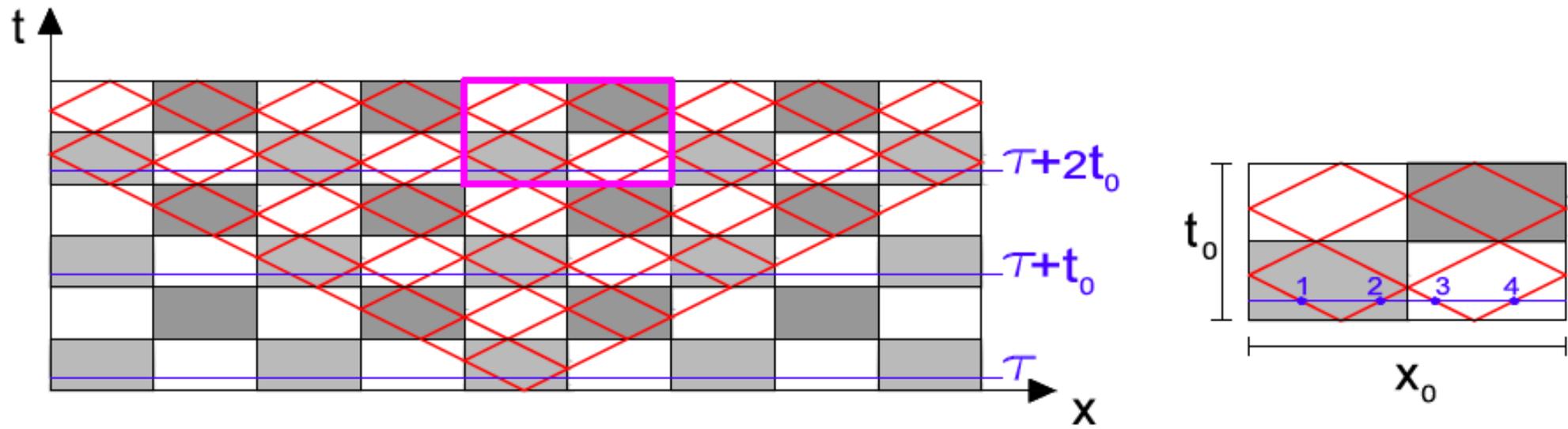
# Evolution in time of the current distribution

Note: oscillations continue after the wavefront!



# Three-component space-time checkerboard

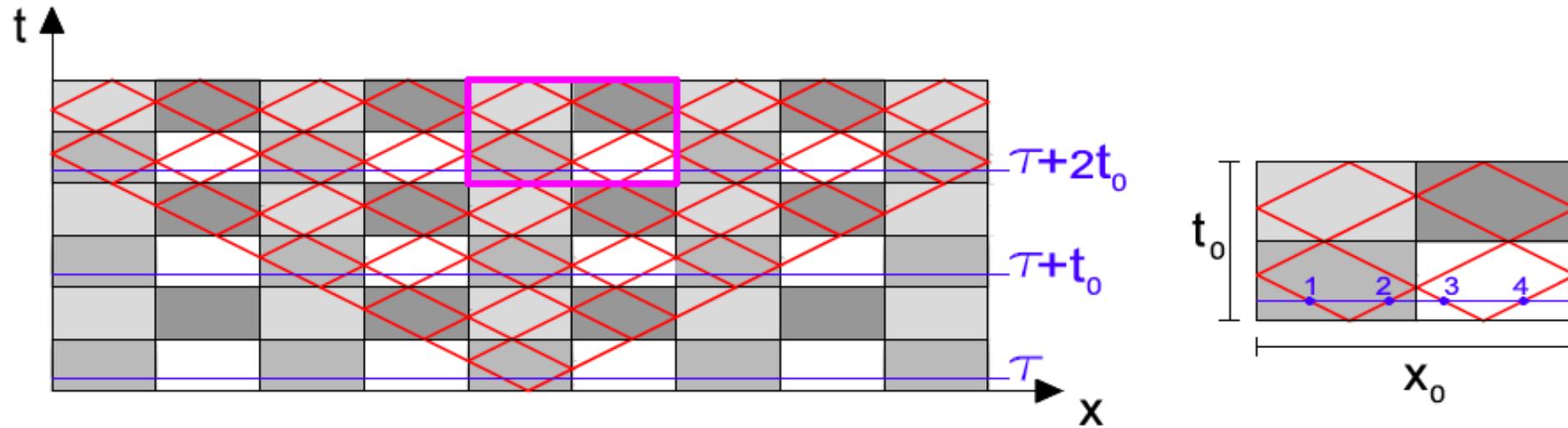
$$c_1 = c_2 = c_3$$



UNBROKEN PT-symmetry  $\Rightarrow$  Only propagating modes

# Four-component space-time checkerboard

$$c_1 = c_2 = c_3 = c_4$$

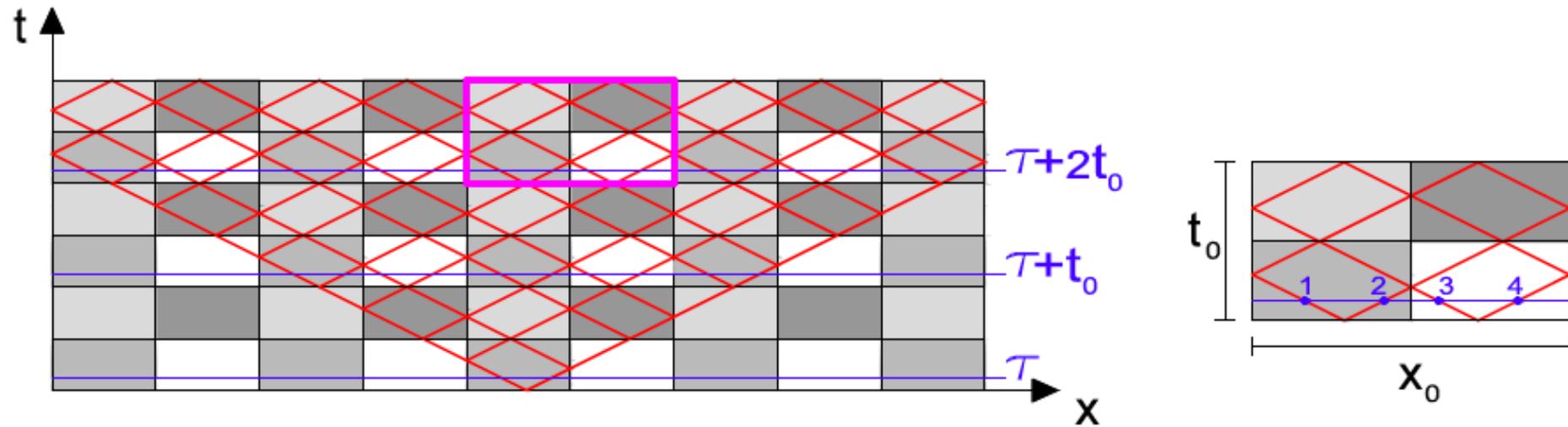


For some combinations of  $\gamma_1, \gamma_2, \gamma_3, \gamma_4$  : **UNBROKEN** PT-symmetry

For other combinations of  $\gamma_1, \gamma_2, \gamma_3, \gamma_4$  : **BROKEN** PT-symmetry

# Four-component space-time checkerboard

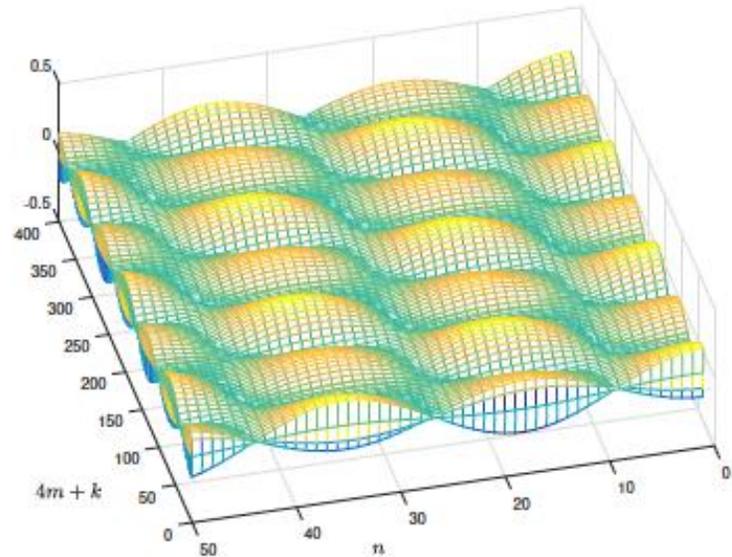
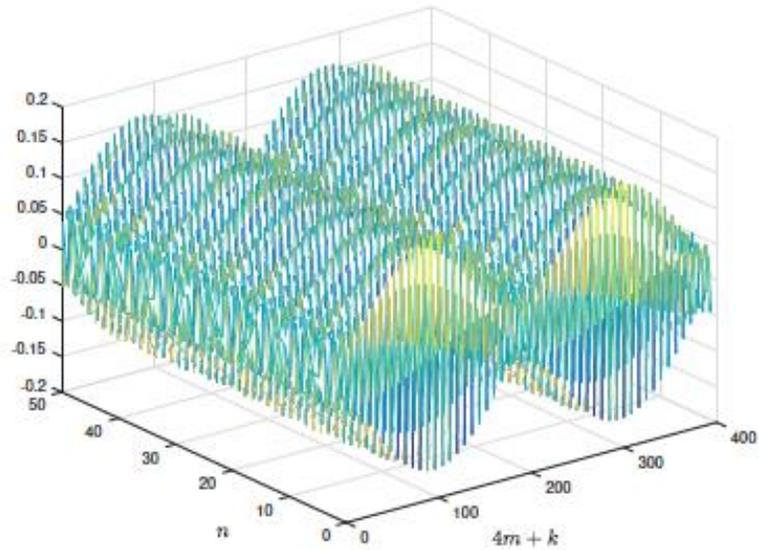
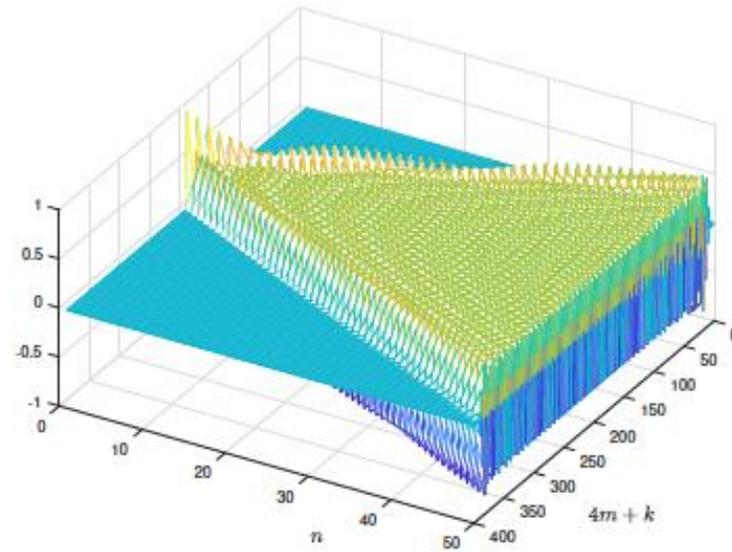
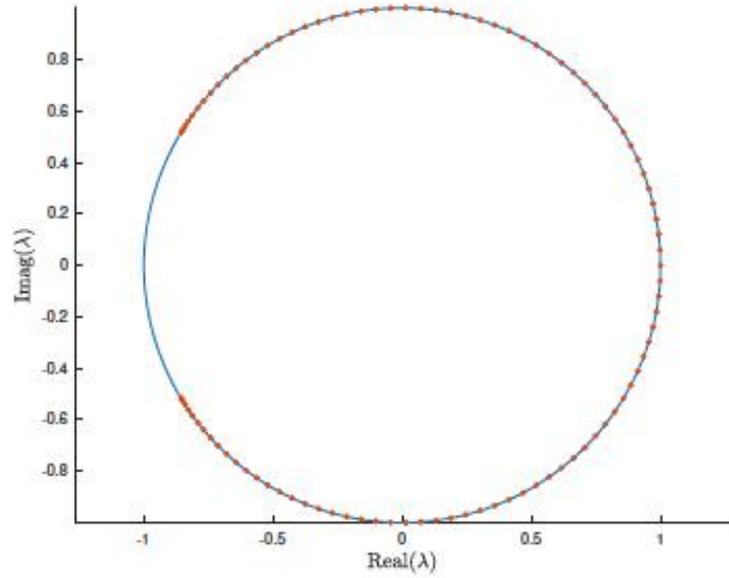
$$c_1 = c_2 = c_3 = c_4$$



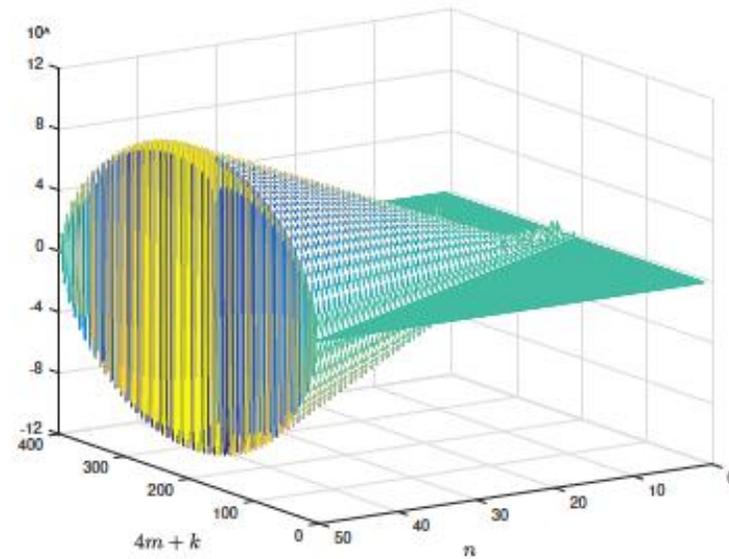
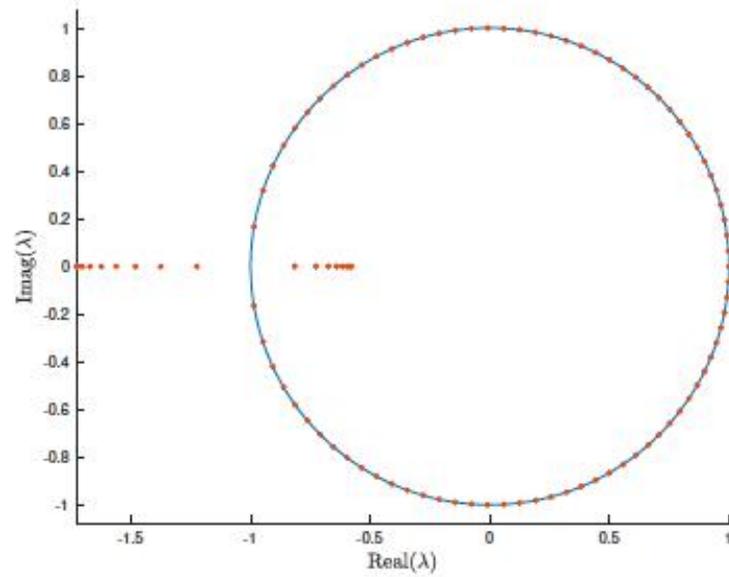
For some combinations of  $\gamma_1, \gamma_2, \gamma_3, \gamma_4$  : **UNBROKEN** PT-symmetry

For other combinations of  $\gamma_1, \gamma_2, \gamma_3, \gamma_4$  : **BROKEN** PT-symmetry

# Unbroken $PT$ -symmetry for the four-phase checkerboard

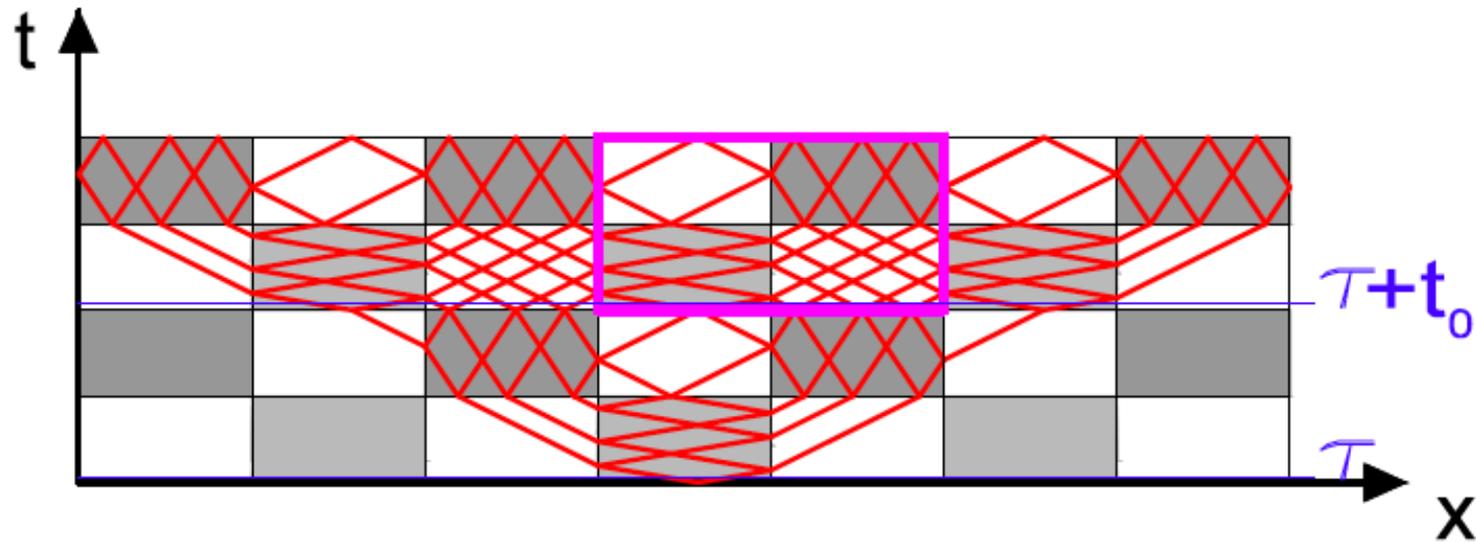


# Broken $PT$ -symmetry for the four-phase checkerboard



# Three-phase space-time checkerboard

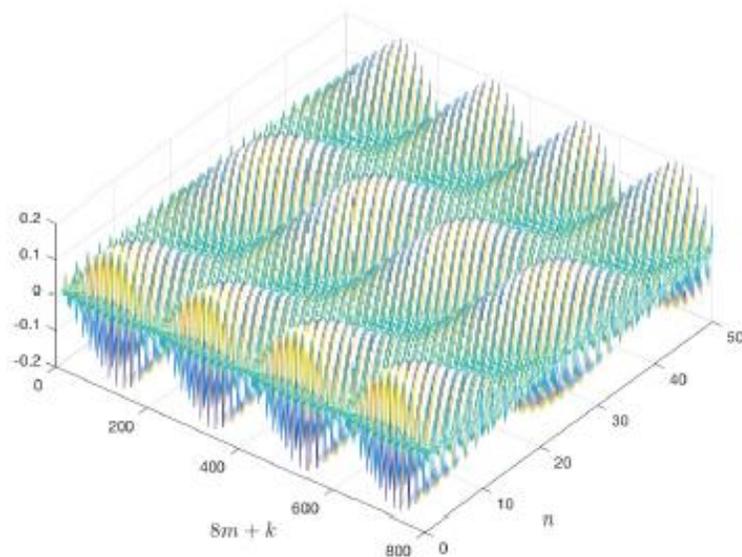
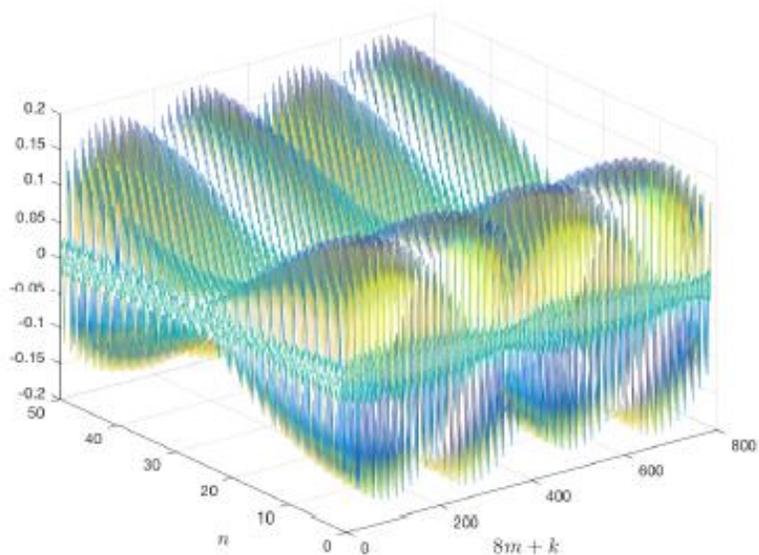
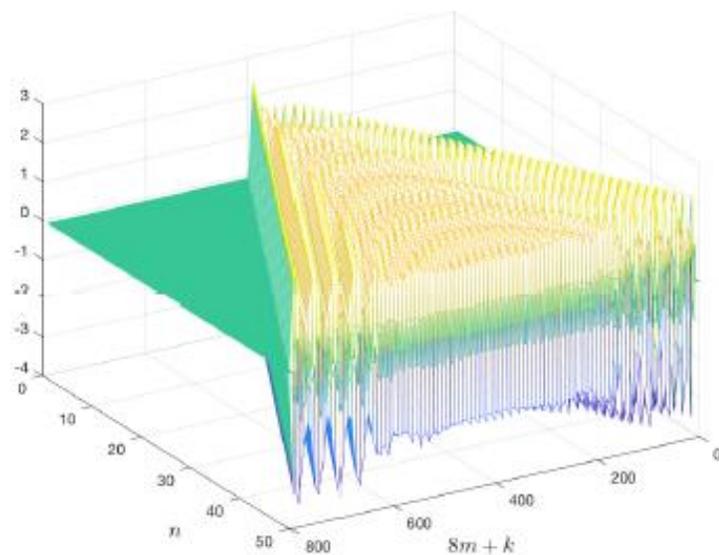
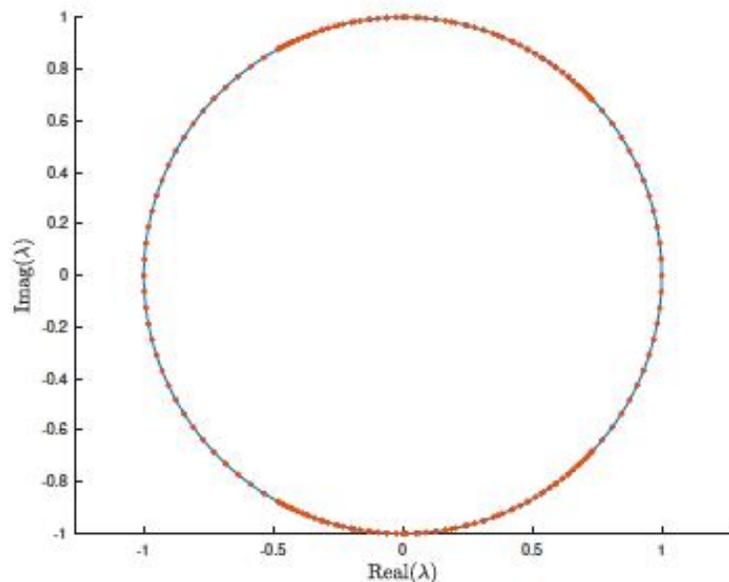
$$c_2/c_1 = c_1/c_3 = 3$$



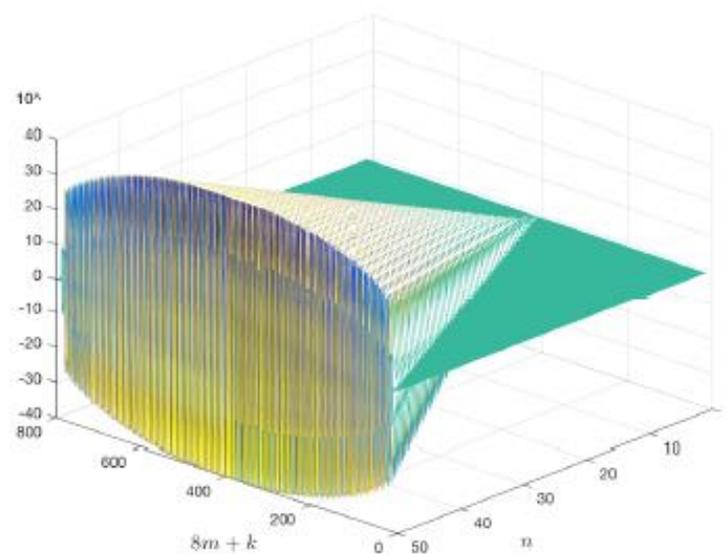
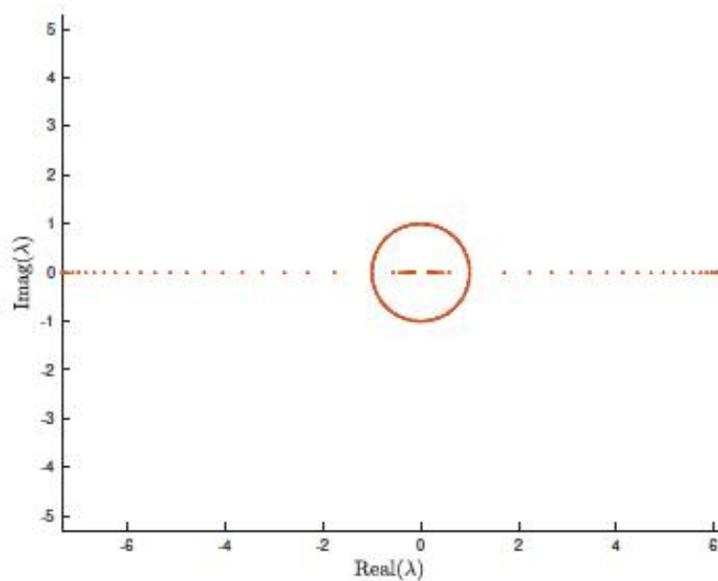
For some combinations of  $\gamma_1, \gamma_2, \gamma_3$  : **UNBROKEN** PT-symmetry

For other combinations of  $\gamma_1, \gamma_2, \gamma_3$  : **BROKEN** PT-symmetry

# Unbroken $PT$ -symmetry for the three-phase checkerboard



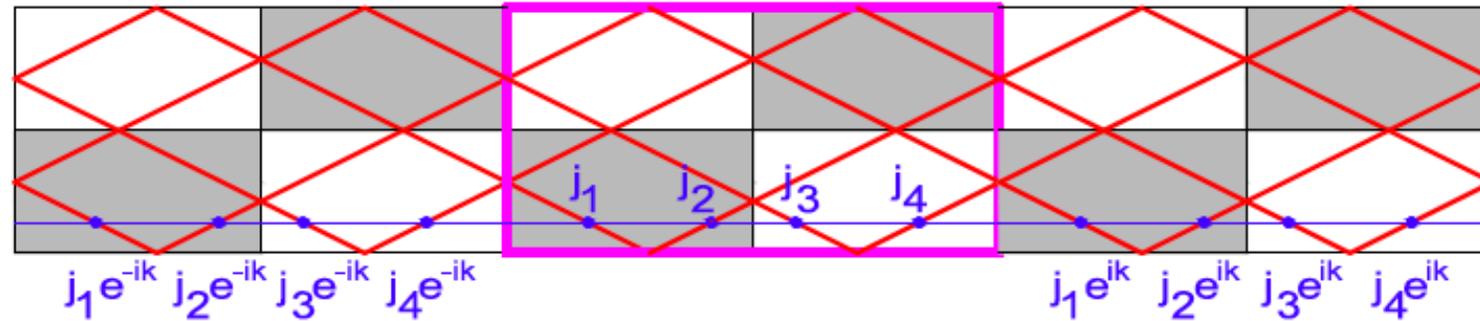
# Broken $PT$ -symmetry for the three-phase checkerboard



# Bloch–Floquet theory applied to field patterns

Periodicity with respect to  $x$ :

$$j(l, m + s, n) = \exp(iks) j(l, m, n)$$



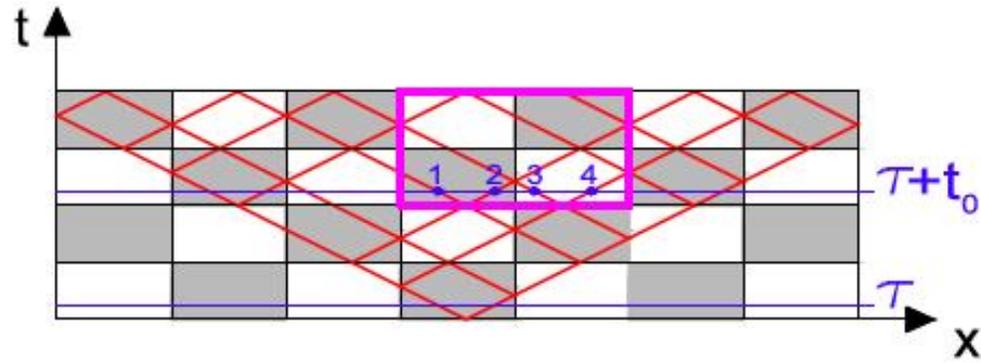
Periodicity with respect to  $t$ :

$$j(l, m, n + q) = \exp(i\omega q) j(l, m, n)$$

Recall:  $j(l, m, n + q) = \lambda^q j(l, m, n)$ , then

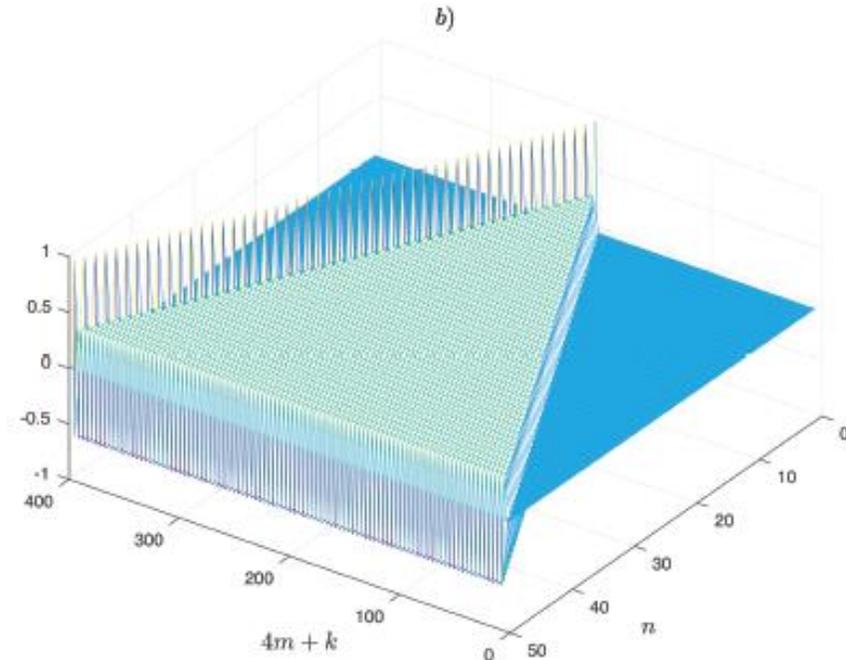
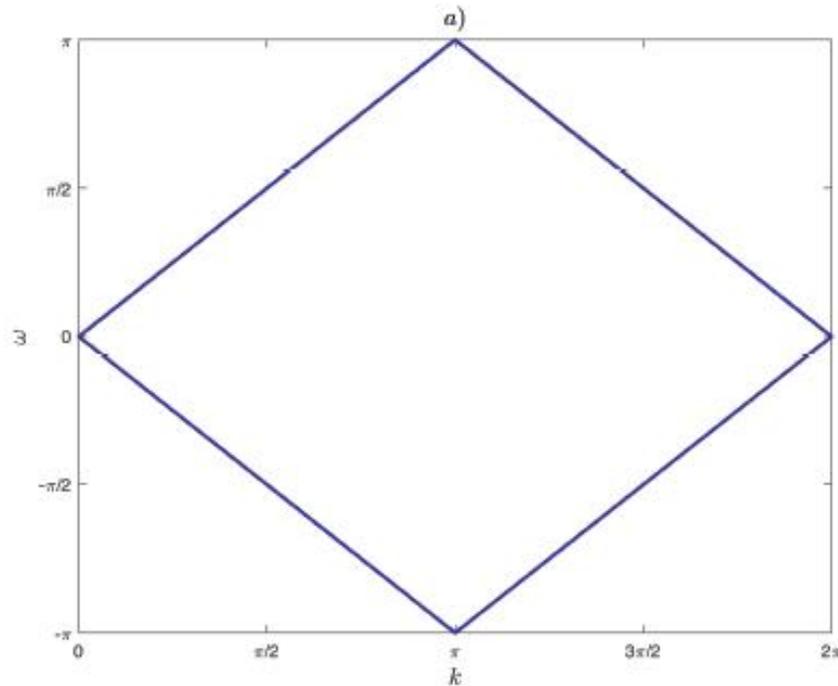
Dispersion relation :  $\lambda(k) = \exp(i\omega)$

# Dispersion diagram for the two-phase checkerboard

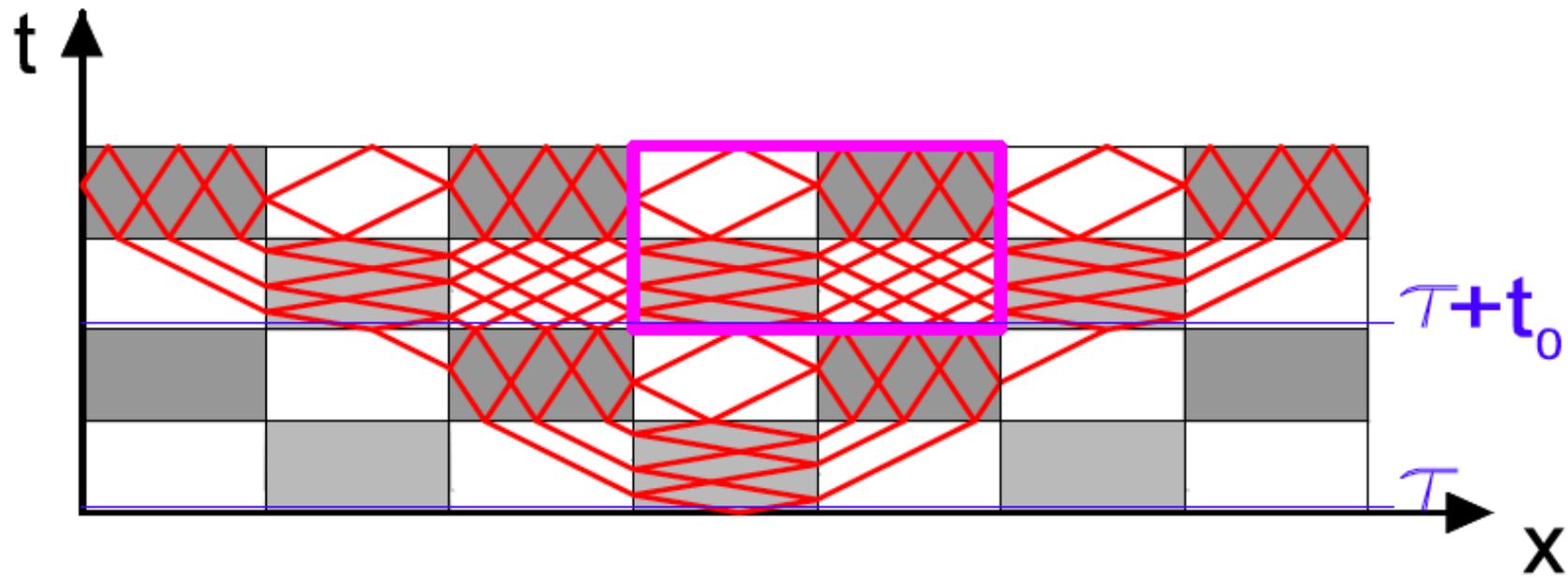


$$\lambda(k) = \exp(\pm ik)$$

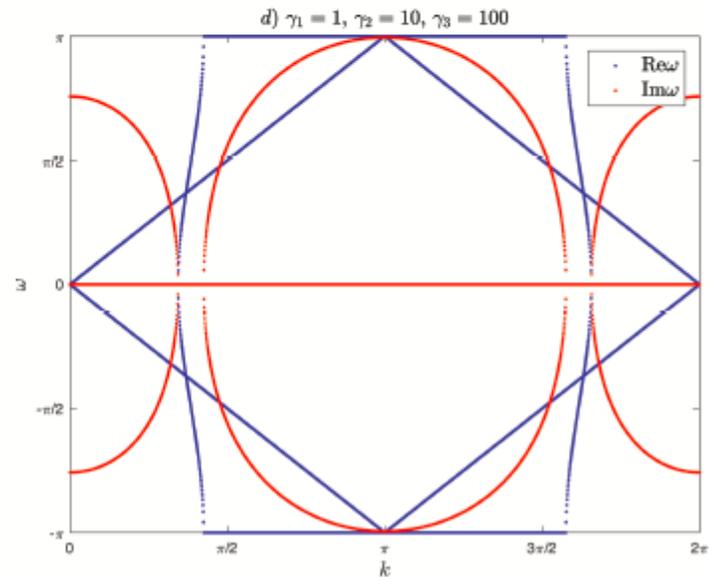
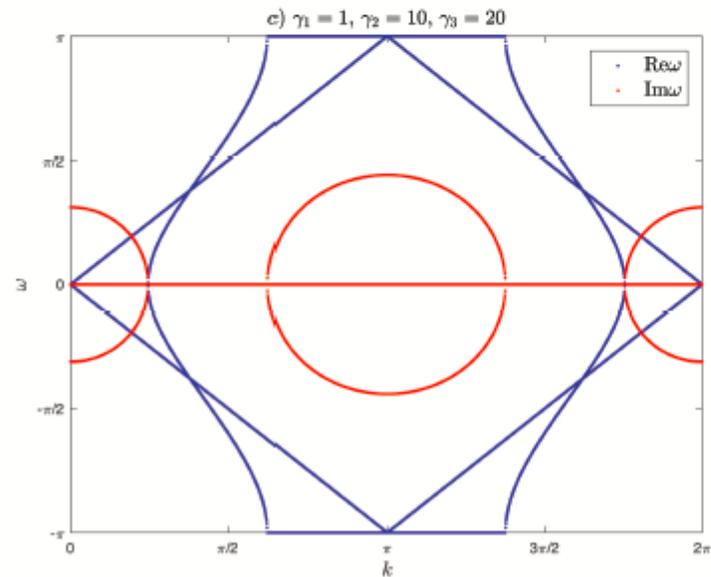
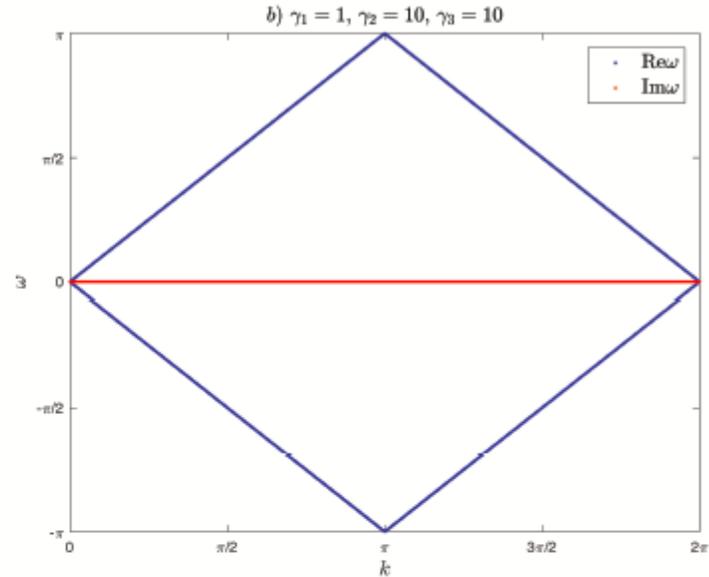
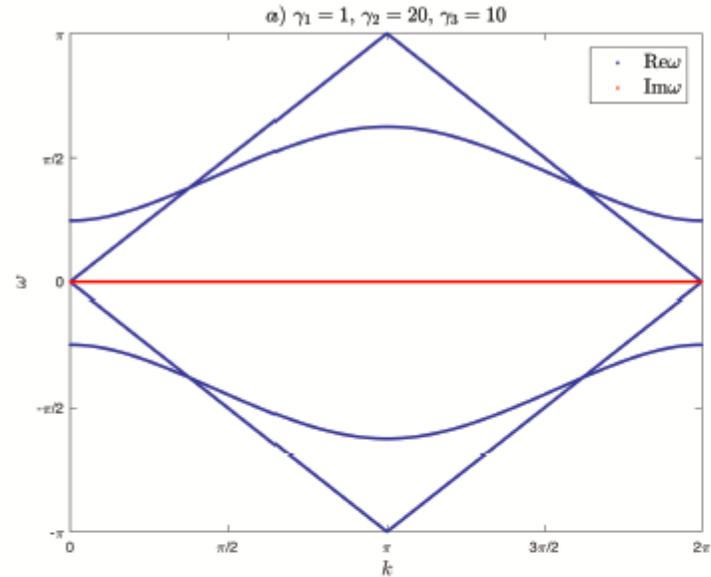
$$\text{Dispersion relation: } \omega = \pm k$$



Three-phase checkerboard with phases having speed in a certain ratio

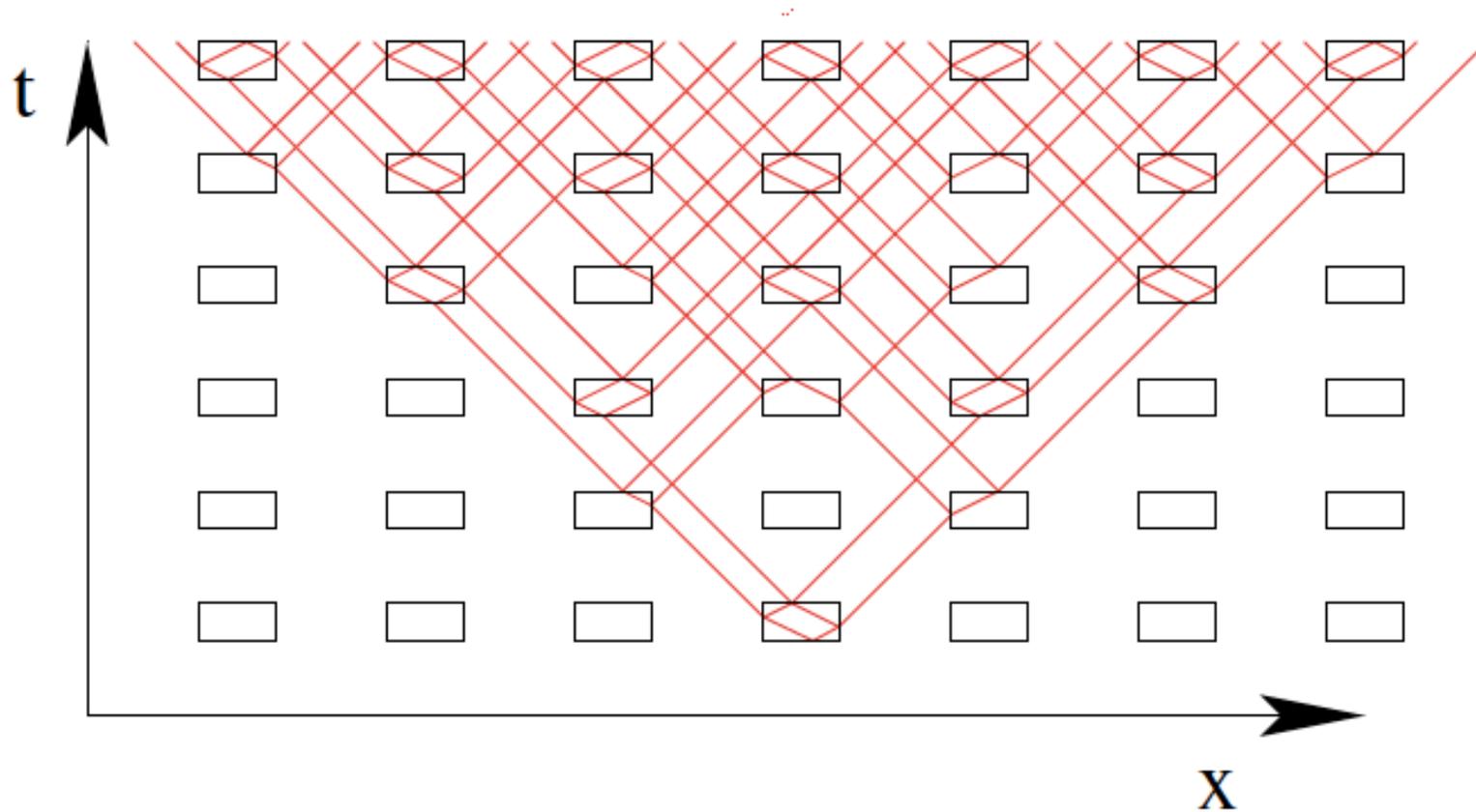


# Dispersion diagrams for the three-phase checkerboard



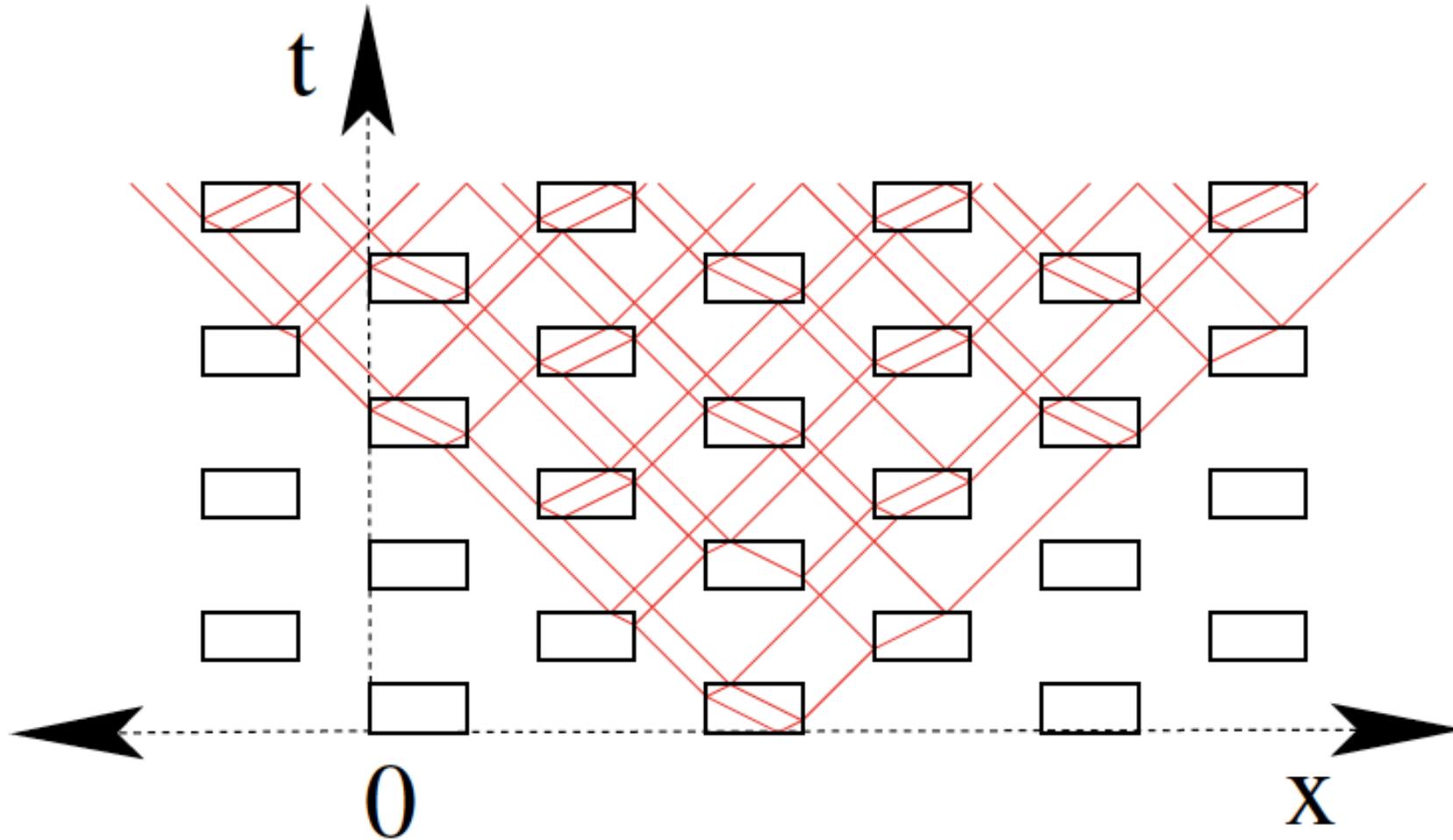
Bloch Waves are:  
**Infinitely Degenerate!**

# What about other field pattern geometries?



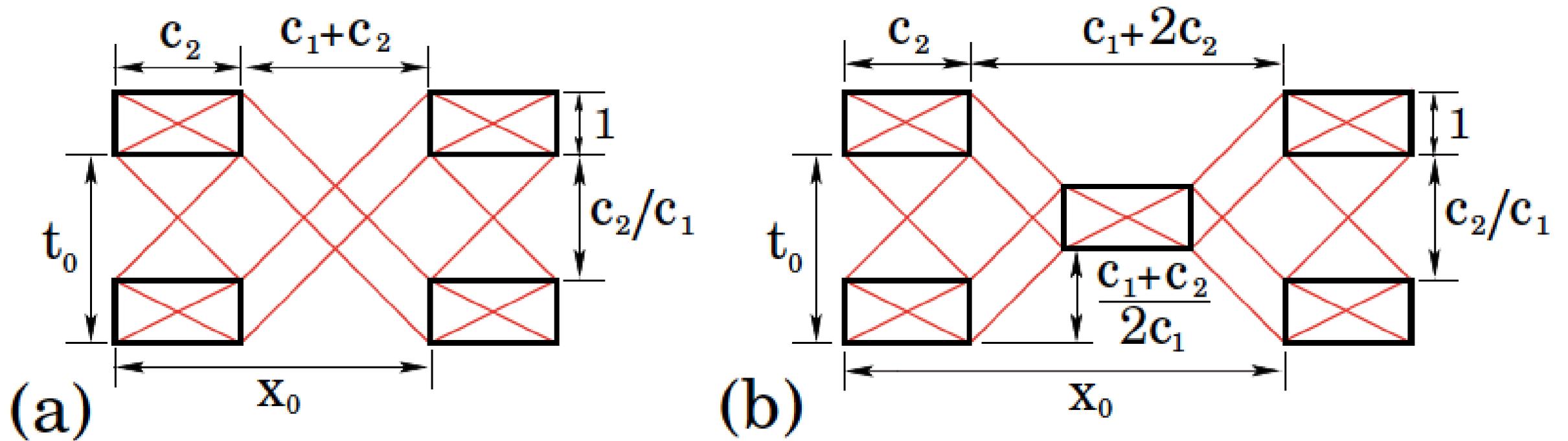
Characteristic lines lie on a pattern—the field pattern.  
Note the P-T symmetry of the microstructure.

Alternatively one can have staggered inclusions:

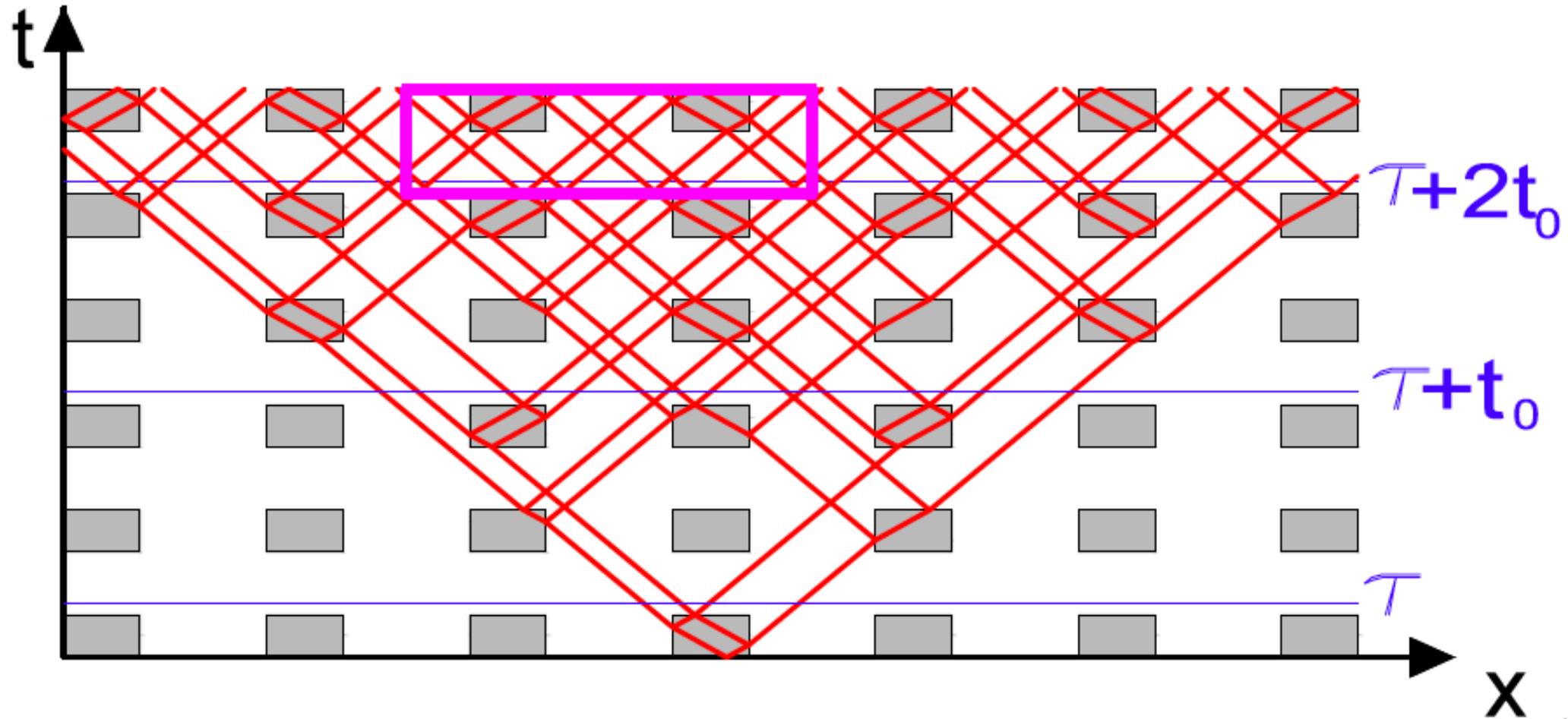


Again, note the P-T symmetry of the microstructure.

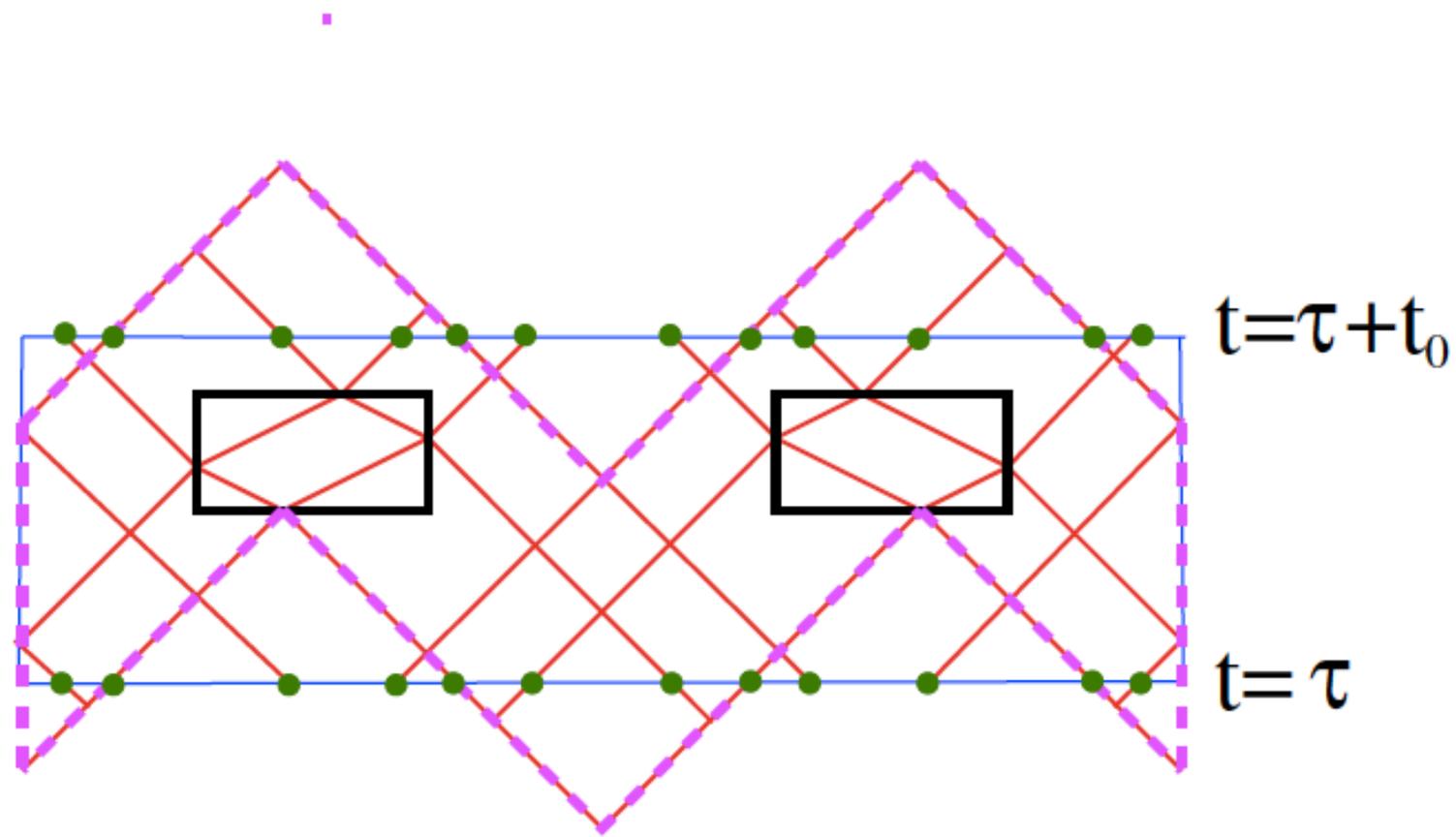
# Geometry: Relation to Characteristic Lines



# Space-time microstructure with rectangular inclusions



# Numerical results: Transfer Matrix

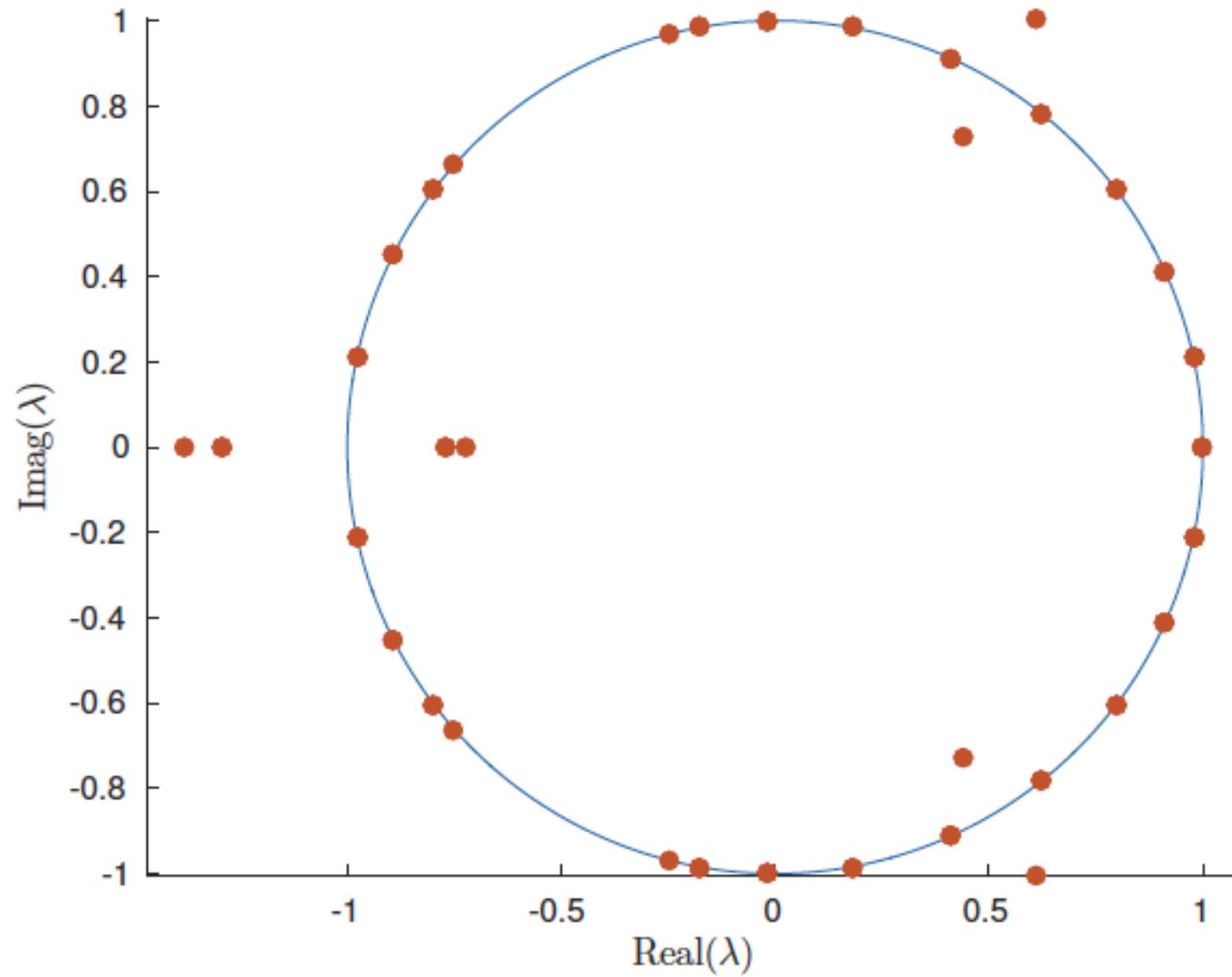


Note: The period cell of the field pattern is twice that of the microstructure!

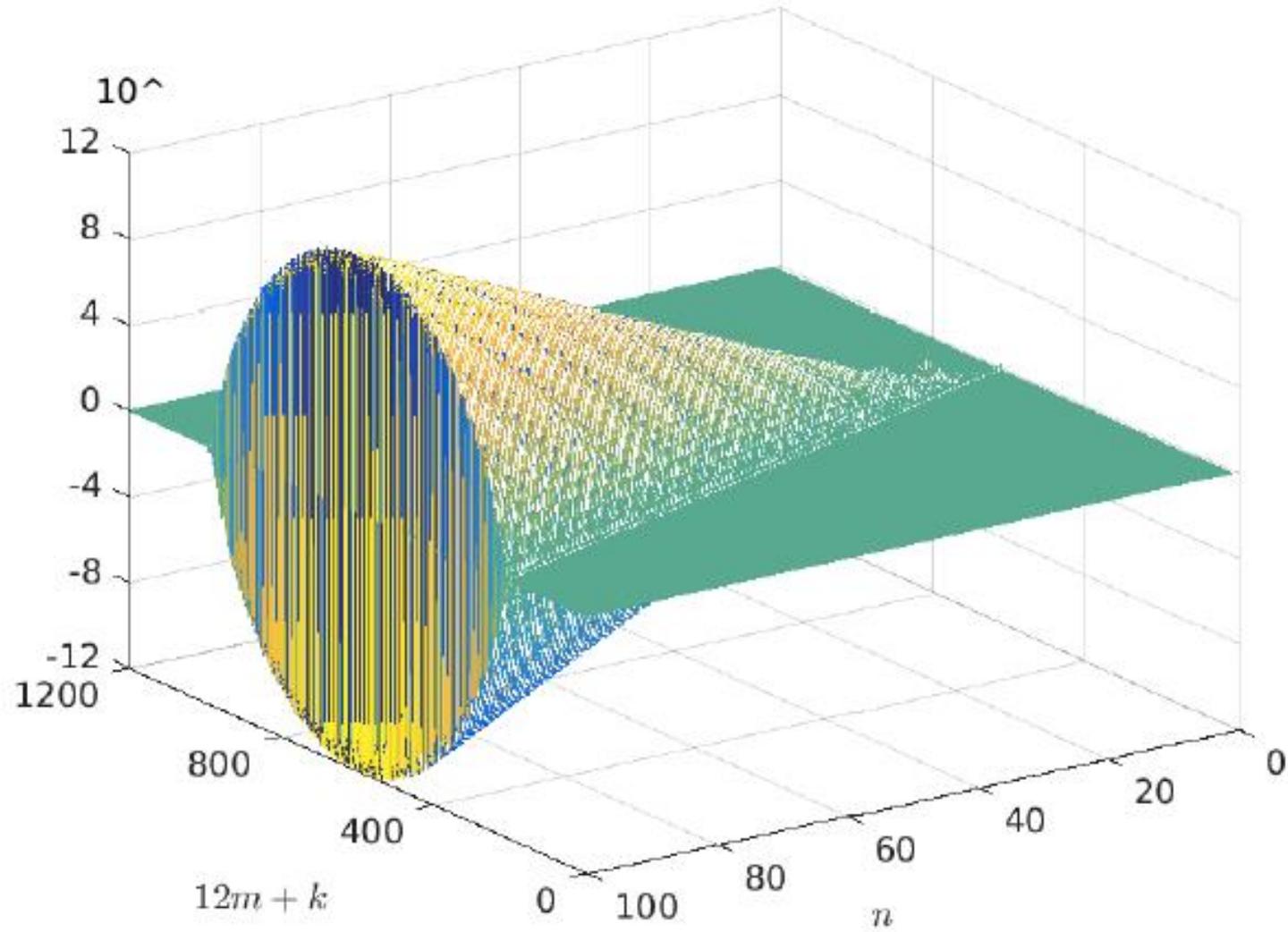
$$j(k, m, n + 1) = \sum_{k', m'} T_{(k, m), (k', m')} j(k', m', n)$$

$$T_{(k, m), (k', m')} = G_{k, k'}(m - m')$$

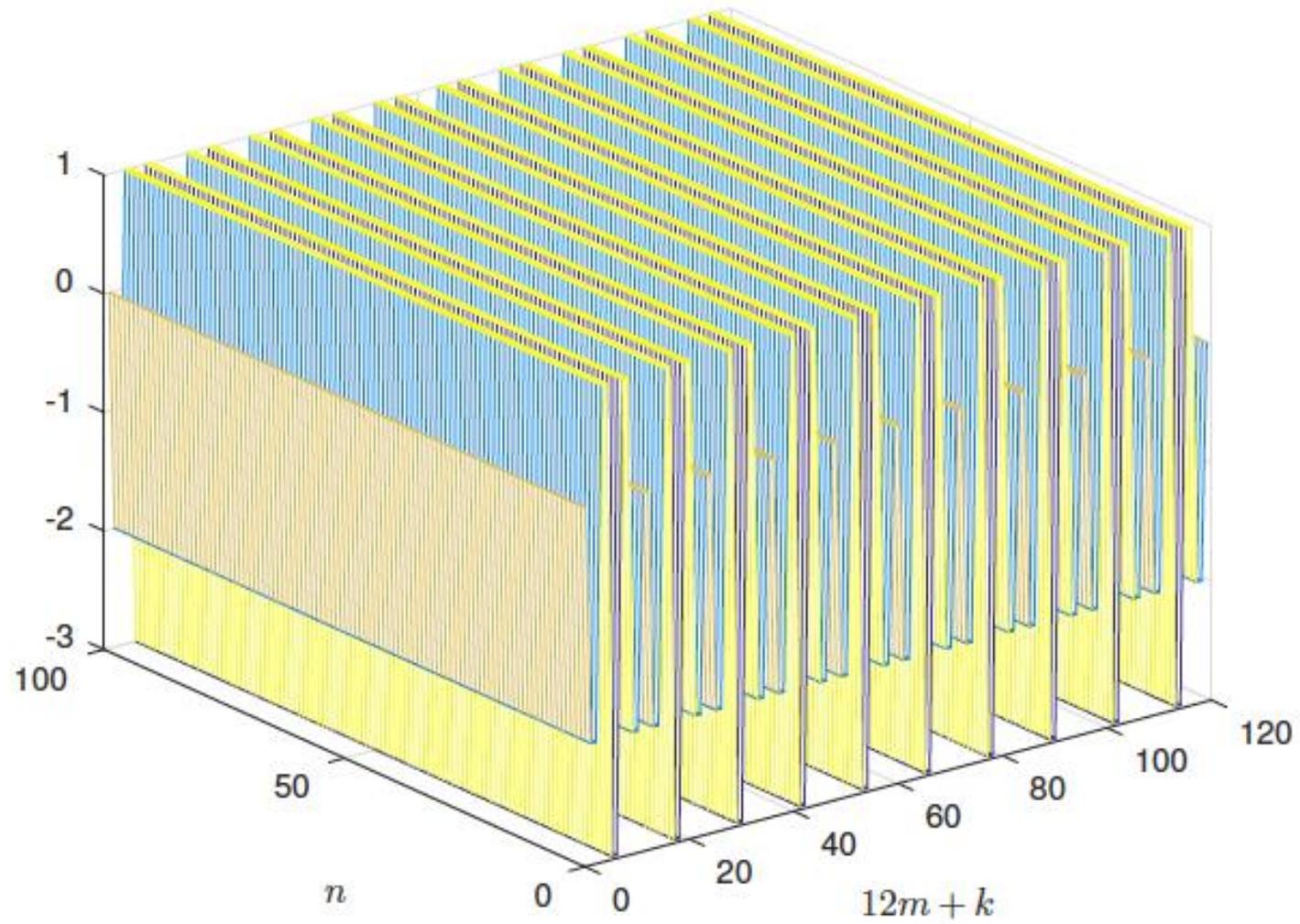
# Eigenvalues of the transfer matrix



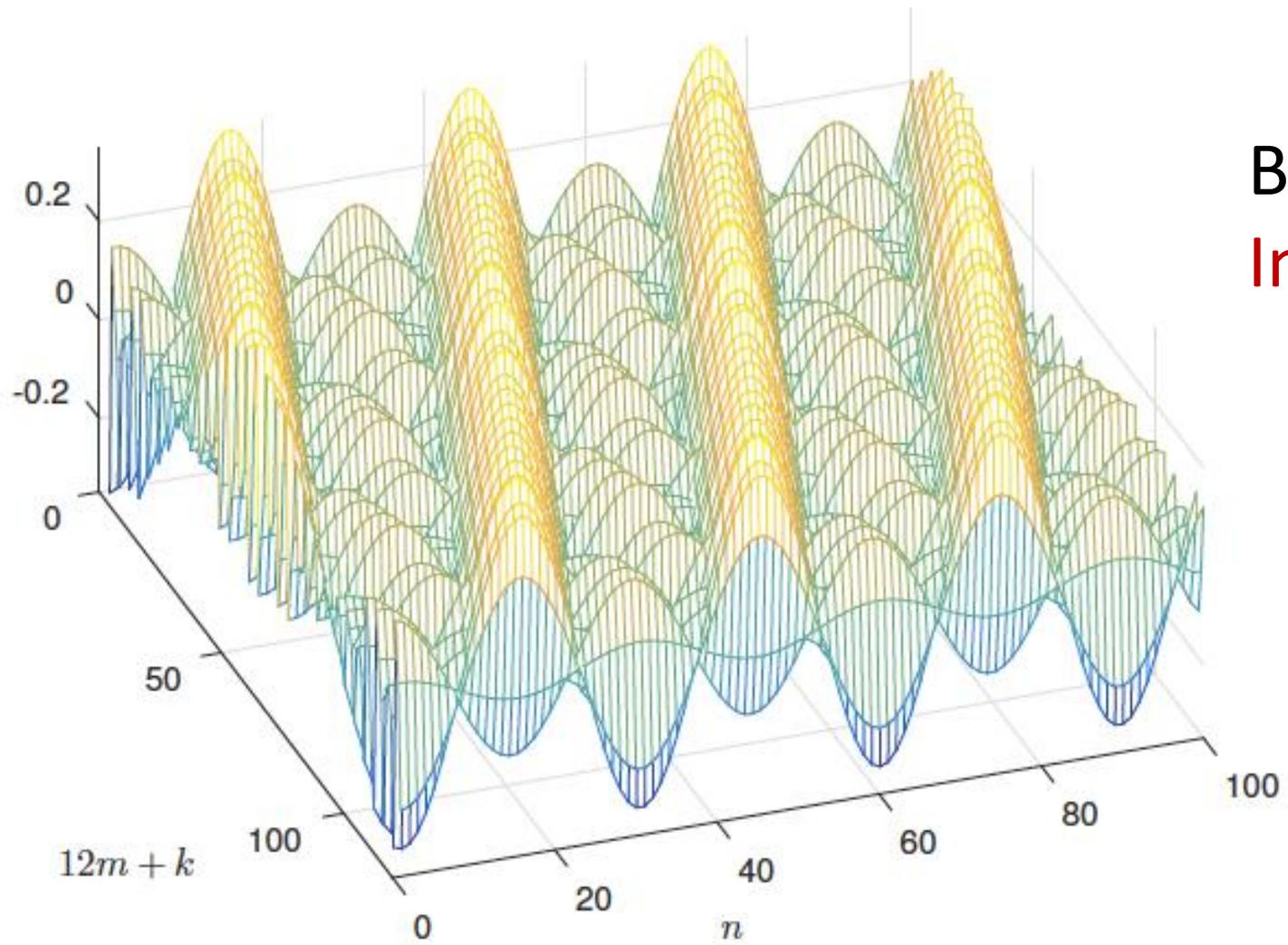
# Blow up



# Periodic solution

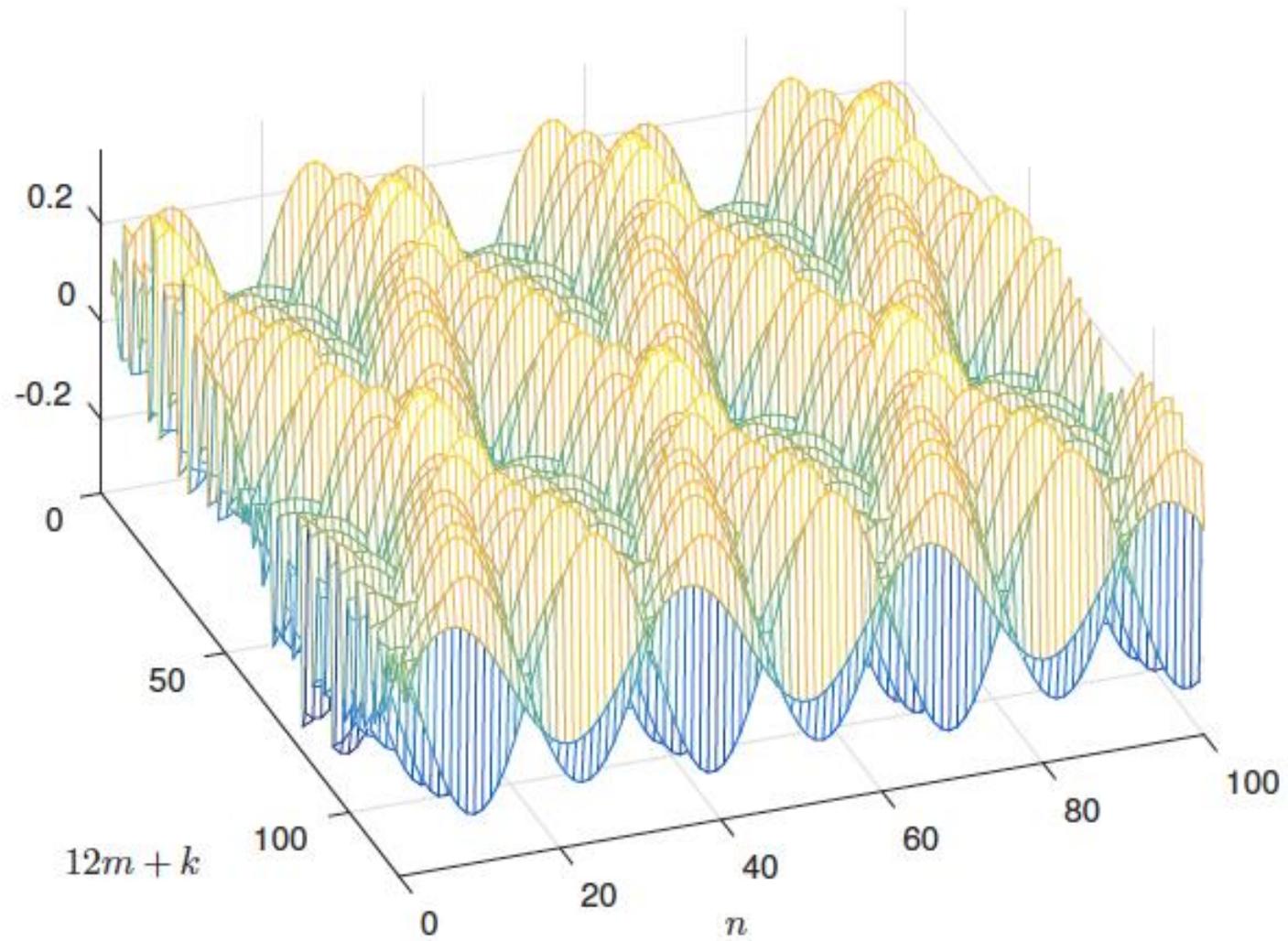


# An example of a solution that does not blow up

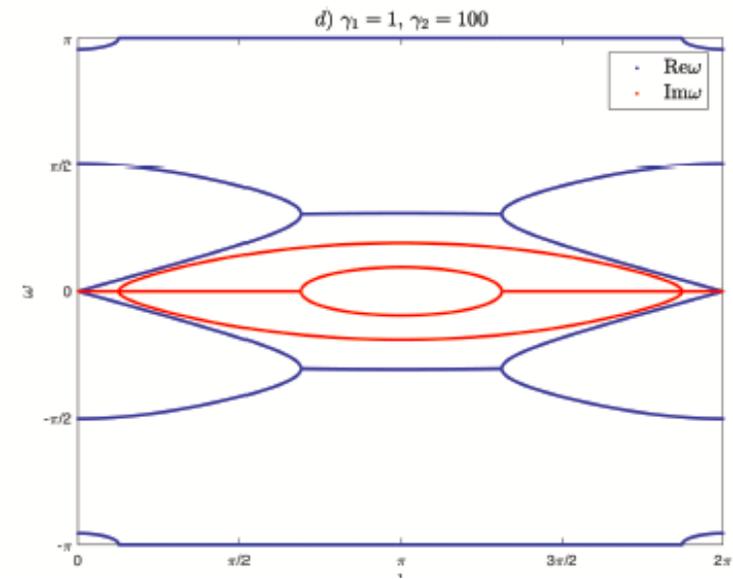
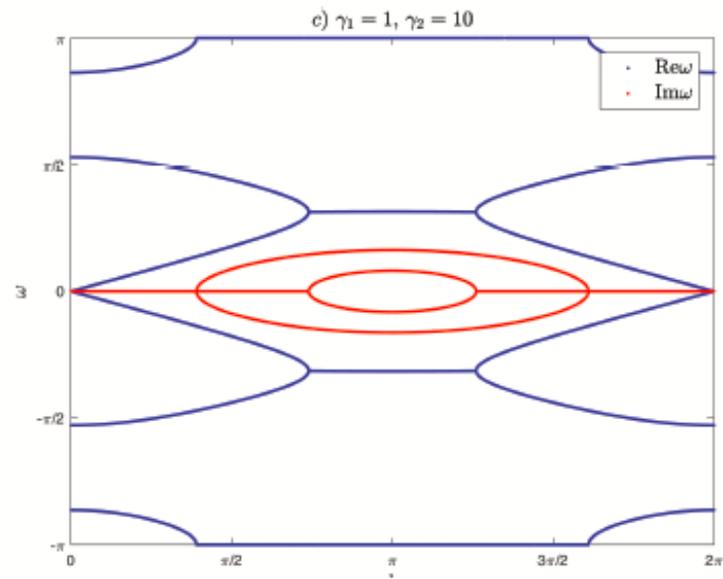
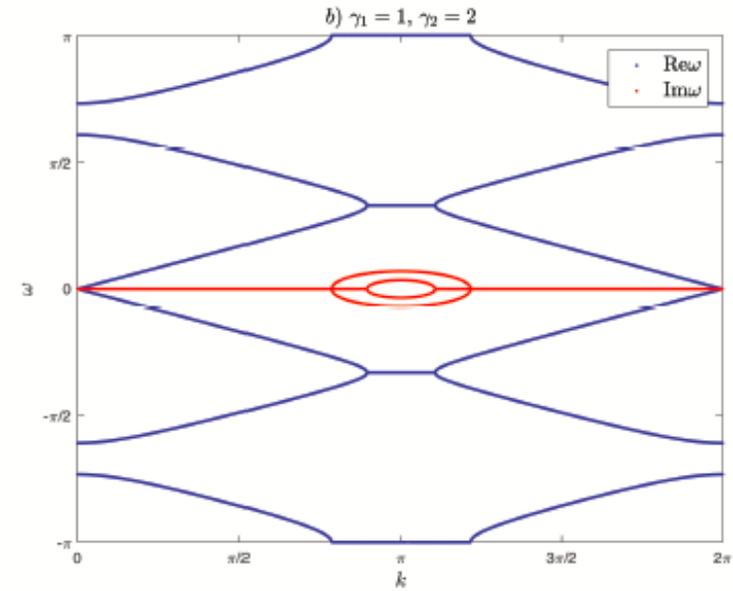
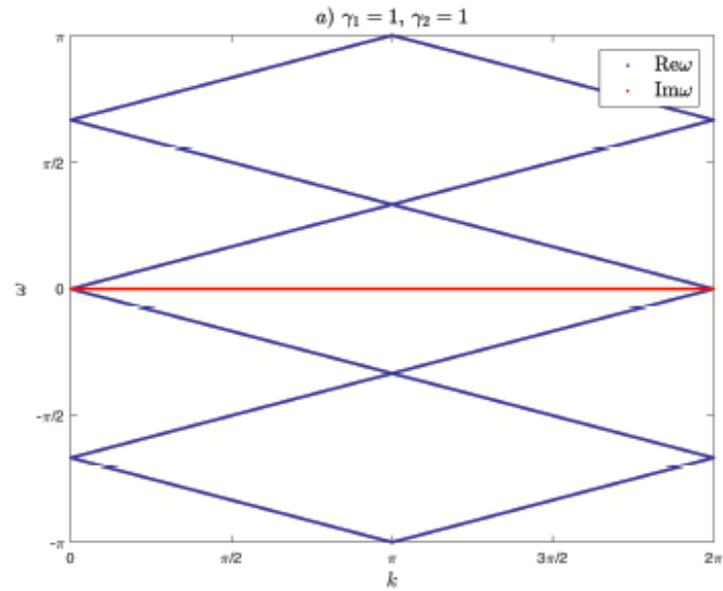


Bloch Wave:  
**Infinitely Degenerate!**

# One more solution that does not blow up



# Dispersion diagrams for the microstructure with inclusions



Thank you for your attention!!



Milton GW, Mattei O, 2017. *Field patterns: a new mathematical object*. Proc R Soc A. 473:20160819.



Mattei O, Milton GW, 2017. *Field patterns without blow up*. To appear in New J Phys.



Mattei O, Milton GW, 2017. *Field patterns: a new type of wave with infinitely degenerate band structure*. Submitted.

# Extending the Theory of Composites to Other Areas of Science

Edited By  
Graeme W. Milton



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to Other Areas of Science

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