- 1. Consider the velocity field for a Newtonian fluid between two infinite parallel disks that are separated by a gap H. The lower disk rotates about the central axis with angular velocity Ω . Assume that the other disk is stationary. The governing equations are given below.
 - Continuity:

$$\frac{1}{r} \left[\partial_r (r u_r) + \partial_\theta u_\theta \right] + \partial_z u_z = 0.$$

r-momentum:

$$\rho\left(u_r\partial_r u_r + \frac{u_\theta}{r}\partial_\theta u_r + u_z\partial_z u_r - \frac{u_\theta^2}{r}\right) = -\partial_r p$$
$$+ \mu\left[\partial_r\left(\frac{1}{r}\partial_r(ru_r)\right) + \frac{1}{r^2}\partial_{\theta\theta}u_r - \frac{2}{r^2}\partial_\theta u_\theta + \partial_{zz}u_r\right]$$

 θ -momentum:

$$\rho\left(\partial_t u_\theta + u_r \partial_r u_\theta + \frac{u_\theta}{r} \partial_\theta u_\theta + u_z \partial_z u_\theta + \frac{u_\theta u_r}{r}\right) = -\frac{1}{r} \partial_\theta p$$
$$+ \mu \left[\partial_r \left(\frac{1}{r} \partial_r (r u_\theta)\right) + \frac{1}{r^2} \partial_{\theta\theta} u_\theta + \frac{2}{r^2} \partial_\theta u_r + \partial_{zz} u_\theta\right]$$

z-momentum:

$$\rho\left(\partial_t u_z + u_r \partial_r u_z + \frac{u_\theta}{r} \partial_\theta u_z + u_z \partial_z u_z\right) = -\partial_z p + \mu \left[\partial_r \left(\frac{1}{r} \partial_r (ru_z)\right) + \frac{1}{r^2} \partial_{\theta\theta} u_z + \partial_{zz} u_z\right]$$

The boundary conditions are:

At z = 0, $u_r = u_z = 0$ and $u_\theta = \Omega r$. At z = H, $u_r = u_z = u_\theta = 0$.

- 1. Nondimensionalize the equations. Note that there are two choices for the characteristic pressure either $p_c = \Omega \mu$ or $p_c = \rho \Omega^2 H^2$. Explain how the second choice is the correct one when considering the limit $\text{Re} \to 0$, while the first one is to be used in the limit $\text{Re} \to 0$.
- 2. Solve the equation in the creeping flow limit, i.e Re = 0. Note that the solution is radially symmetric and independent of θ .
- 3. Use a regular perturbation series to obtain the first order solution, i.e. up to and including terms O(Re).
- 4. Use Matlab or another software to plot the secondary flow (the first order perturbation).
- 2. Consider a general similarity solution $\Psi = F(x)f(\eta)$ with $\eta = y/g(x)$ to the boundary layer equations

$$u\partial_x u + v\partial_y u = -\frac{1}{\rho}\partial_x p + \nu\partial_y u \quad \partial_x u + \partial_y v = 0$$

where Ψ is a the stream function.

1. Show that the condition, $u \to U(x)$ as $y/\delta \to \infty$, demands that F(x) = cU(x)g(x), where c is a constant, and then choose c to be 1.

2. Show that substitution in the boundary layer equations leads to

$$f'^{2} - \left(1 + \frac{U}{U'}\frac{g'}{g}\right)ff'' = 1 + \frac{\nu f'''}{g^{2}U'}.$$

Eliminate the pressure by matching to the outer solution (Bernoulli), i.e.

$$-\frac{1}{\rho}\partial_x p = u_e \partial_x u_e = U(x)U'(x)$$

- 3. Deduce that a similarity solution is only possible if either $U(x) \sim (x x_0)^m$ or $U(x) \sim e^{\alpha x}$, where x_0, m, α are constants.
- 4. In the case, $U(x) = Ax^m$, A > 0, show that g(x) is proportional to $x^{(1-m)/2}$, and show that choosing

$$g(x) = \sqrt{\frac{2\nu}{(m+1)Ax^{m-1}}}$$

leads to

$$f''' + ff'' + \frac{2m}{m+1}(1 - f'^2) = 0,$$

subject to f(0) = f'(0) = 0 and $f'(\infty) = 1$.

3. Consider uniform slow flow past a circular cylinder, and show that the problem reduces to

$$\left(\frac{\partial^2}{\partial r^2} + \frac{1}{r}\frac{\partial}{\partial r} + \frac{1}{r^2}\frac{\partial^2}{\partial \theta^2}\right)^2\Psi = 0,$$

with $\frac{\partial \Psi}{\partial r} = \frac{\partial \Psi}{\partial \theta} = 0$ on r = a and $\Psi \sim Ur \sin \theta$ as $r \to \infty$. Show that seeking a solution of the form $\Psi = f(r) \sin \theta$ leads to

$$\Psi = \left[Ar^3 + Br\log r + Cr + \frac{D}{r}\right]\sin\theta$$

and thus fails, in that for no choice of the arbitrary constants can all the boundary conditions be satisfied.