Math 2280-001 Week 12, April 3-7

Mon Apr 3 Finish section 5.7 notes from Friday. We will also discuss questions you may have about the section 5.6 homework due on Tuesday. A lab problem for next week will be to work the input-output problem from the second midterm using (1) matrix exponentials, and (2) the diagonalization method. We may begin that problem in lab format, if time permits.

finish up day handout today for lab problem homorrows

$$\begin{bmatrix} -t e^{2t} \\ (1-t) e^{2t} \end{bmatrix}$$

$$\begin{bmatrix} e^{2t} (1+t) & -t e^{2t} \\ t e^{2t} & (1-t) e^{2t} \end{bmatrix}$$

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$$\begin{bmatrix} e^{2t} (1+t) & -t e^{2t} \\ t e^{2t} & -e^{2t} (-1+t) \end{bmatrix}$$

$$\begin{bmatrix} e^{2t} (1+t) & -t e^{2t} \\ t e^{2t} & -e^{2t} (-1+t) \end{bmatrix}$$

$$(4)$$

Last wed we computed et A for an A not diagonalizable didn't finish rest of notes }.

recall

$$e^{At} = I + tA + \frac{t^2}{2!}A^2 + \dots$$

$$= \Phi(t)\Phi(\eta^{-1})$$

$$\uparrow$$

$$F.M., \omega |_{umns} basis η solhs $\forall x' = Ax$$$

Start here Monday 4/2

Variation of parameters: This is what fundamental matrices and matrix exponentials are especially good for....they let you solve non-homogeneous systems without guessing. Consider the non-homogeneous first order system

$$\mathbf{x}'(t) = P(t)\mathbf{x} + \mathbf{f}(t)$$

$$\underline{\boldsymbol{x}}'(t) = P(t)\underline{\boldsymbol{x}}.$$

$$\Phi'(t) = P(t) \Phi(t)$$

first order system $\underline{x}'(t) = P(t)\underline{x} + \underline{f}(t) \quad *$ Let $\Phi(t)$ be an FM for the homogeneous system $\underline{x}'(t) = P(t)\underline{x}.$ Since $\Phi(t)$ is invertible for all t we may do a change of functions for the non-homogeneous system: $\underline{x}(t) = \Phi(t)\underline{u}(t) \qquad \qquad \text{(in fad, QL) is invertible}.$

$$\underline{\boldsymbol{x}}(t) = \boldsymbol{\Phi}(t)\underline{\boldsymbol{u}}(t)$$

plug into the non-homogeneous system (*):

$$\Phi'(t)\underline{\boldsymbol{u}}(t) + \Phi(t)\underline{\boldsymbol{u}}'(t) = P(t)\Phi(t)\underline{\boldsymbol{u}}(t) + \boldsymbol{f}(t).$$

Since $\Phi' = P \Phi$ the first terms on each side cancel each other and we are left with

$$\Phi(t)\underline{u}'(t) = f(t)$$

$$\underline{u}' = \Phi^{-1}f \leftarrow RHS \text{ is just a fin of } t$$

which we can integrate to find a $\underline{u}(t)$, hence an $\underline{x}(t) = \Phi(t)\underline{u}(t)$.

Remark: This is where the (mysterious at the time) formula for variation of parameters in n^{th} order linear DE's came from....

"Recall" (February 24 notes):

Variation of Parameters: The advantage of this method is that is always provides a particular solution, even for non-homogeneous problems in which the right-hand side doesn't fit into a nice finite dimensional subspace preserved by L, and even if the linear operator L is not constant-coefficient. The formula for the particular solutions can be somewhat messy to work with, however, once you start computing.

Here's the formula: Let $y_1(x), y_2(x), ..., y_n(x)$ be a basis of solutions to the homogeneous DE

$$L(y) := y^{(n)} + p_{n-1}(x)y^{(n-1)} + \dots + p_1(x)y' + p_0(x)y = 0.$$

 $L(y) := y^{(n)} + p_{n-1}(x)y^{(n-1)} + ... + p_1(x)y' + p_0(x)y = 0.$ Then $y_p(x) = u_1(x)y_1(x) + u_2(x)y_2(x) + ... + u_n(x)y_n(x)$ is a particular solution to L(y) = f provided the coefficient functions (aka "varying parameters") $u_1(x), u_2(x), ... u_n(x)$ have derivatives satisfying the Wronskian matrix equation

$$\begin{bmatrix} y_1 & y_2 & \cdots & y_n \\ y_1' & y_2' & \cdots & y_n' \\ \cdots & \cdots & \cdots & \cdots \\ y_1^{(n-1)} & y_2^{(n-1)} & \cdots & y_n^{(n-1)} \end{bmatrix} \begin{bmatrix} u_1' \\ u_2' \\ \vdots \\ u_n' \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ \vdots \\ f \end{bmatrix}$$

$$W \left(y_1, y_2, \cdots, y_n \right)$$

But if we convert the n^{th} order DE into a first order system for $x_1 = y$, $x_2 = y'$ etc. we have

$$\begin{array}{c} x_1 & (=y) \\ x_1' = x_2 & (=y') \\ x_2' = x_3 & (=y'') \\ x_{n-1}' = x_n & (=y^{(n-1)}) \\ x_n' = \left(= y^{(n)} \right) = -p_0(x) \mathbf{x}_1 - p_1(x) \mathbf{x}_2 - \dots - p_{n-1}(x) \mathbf{x}_{n-1} + f. \end{array}$$

And each basis solution y(t) for L(y) = 0 gives a solution $[y, y', y'', ...y^{(n-1)}]^T$ to the homogeneous system

$$\begin{bmatrix} x_1' \\ x_2' \\ x_3' \\ \vdots \\ x_n' \end{bmatrix} = \begin{bmatrix} 0 & 1 & 0 & \dots & 0 \\ 0 & 0 & 1 & \dots & 0 \\ 0 & 0 & 0 & 1 \dots & 0 \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ -p_0 & -p_1 & -p_2 & \dots & -p_{n-1} \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \\ \vdots \\ x_n \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ 0 \\ \vdots \\ f \end{bmatrix}.$$

So the original Wronskian matrix for the n^{th} order linear homogeneous DE is a FM for the system above, so the formula we learned in Chapter 3 is a special case of the easier to understand one for first order systems that we just derived, namely

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Returning to first order systems, if we want to solve an IVP for a first order system rather than find the complete general solution, then the following two ways are appropriate:

1) If you want to solve the IVP

$$\underline{\boldsymbol{x}}'(t) = P(t)\underline{\boldsymbol{x}} + \boldsymbol{f}(t)$$
$$\underline{\boldsymbol{x}}(0) = \underline{\boldsymbol{x}}_0$$

The the solution will be of the form $\underline{x} = \Phi \underline{u}$ (where $\underline{u}' = \Phi^{-1} f$ as above). Thus $\underline{\boldsymbol{x}}_0 = \boldsymbol{\Phi}(0)\underline{\boldsymbol{u}}_0$

so

Thus

$$\underline{\underline{u}_{0}} = \Phi(0) \underline{\underline{u}_{0}}$$

$$\underline{\underline{u}_{0}} = \Phi(0)^{-1} \underline{\underline{x}_{0}}.$$

$$\underline{\underline{u}_{0}} = \Phi(0)^{-1} \underline{\underline{x}_{0}}.$$

$$\underline{\underline{t}} = \underline{\underline{u}_{0}} + \int_{0}^{t} \underline{\underline{u}'(s)} ds$$
F7C.

$$\underline{\boldsymbol{u}}(t) = \underline{\boldsymbol{u}}_0 + \int_0^t \underline{\boldsymbol{u}}'(s) \, \mathrm{d}s$$

$$\underline{\boldsymbol{u}}(t) = \underline{\boldsymbol{u}}_0 + \int_0^t \underline{\boldsymbol{\Phi}}^{-1}(s) \boldsymbol{f}(s) \, \mathrm{d}s.$$

Then

$$\underline{\boldsymbol{x}}(t) = \boldsymbol{\Phi}(t)\underline{\boldsymbol{u}}(t)$$

$$\underline{\boldsymbol{x}}(t) = \boldsymbol{\Phi}(t)\left(\underline{\boldsymbol{u}}_0 + \int_0^t \underline{\boldsymbol{\Phi}}^{-1}(s)\boldsymbol{f}(s) \, \mathrm{d}s\right)$$

$$\begin{cases} \underline{\boldsymbol{x}}'(t) = A \, \underline{\boldsymbol{x}} + \boldsymbol{f}(t) \\ \underline{\boldsymbol{x}}(0) = \underline{\boldsymbol{x}}_0 \end{cases}$$

 $\underline{x}(t) = \Phi(t) \left(\underline{u}_0 + \int_0^t \underline{\Phi}^{-1}(s) f(s) \, ds \right).$ 2) If you want to solve the special case IVP $\underbrace{x'(t) = Ax + f(t)}_{x(0) = \underline{x}_0}$ where A is a constant matrix, you may derive a special case of the solution formula above just as we did m

Chapter 1. This is part of empirical. Chapter 1. This is sort of amazing!

$$\begin{array}{c}
\mathbf{L}\mathbf{F}, & \underline{\mathbf{x}}'(t) = A\underline{\mathbf{x}} + \mathbf{f}(t) \\
\underline{\mathbf{x}}'(t) - A\underline{\mathbf{x}} = \mathbf{f}(t) \\
e^{-tA}(\underline{\mathbf{x}}'(t) - A\underline{\mathbf{x}}) = e^{-tA}\mathbf{f}(t) \\
\frac{d}{dt} \left(e^{-tA}\underline{\mathbf{x}}(t) \right) = e^{-tA}\mathbf{f}(t).
\end{array}$$

Integrate from 0 to *t*:

$$e^{-tA}\underline{x}(t) - \underline{x}_0 = \int_0^t e^{-sA}f(s) ds$$

Move the \underline{x}_0 over and multiply both sides by e^{tA} :

$$\underline{\boldsymbol{x}}(t) = e^{tA} \left(\underline{\boldsymbol{x}}_0 + \int_0^t e^{-sA} \boldsymbol{f}(s) ds\right).$$

$$\begin{cases}
\frac{d}{dt} \left(e^{-tA} \right) \\
= e^{-tA} \left(-Ax + x' \right)
\end{cases}$$

$$= e^{-tA} \left(-Ax + x' \right)$$

 $\underline{x}(t) = e^{tA} \left(\underline{x}_0 + \int_0^t e^{-sA} f(s) \, ds \right).$ in theory this formula tells us how to solve any constant or system of DE's!

Exercise 2 Consider the non-homogeneous problem related to the homogeneous system in Exercise 1:

$$\begin{bmatrix} x_1' \\ x_2' \end{bmatrix} = \begin{bmatrix} 3 & -1 \\ 1 & 1 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} + \begin{bmatrix} t \\ 0 \end{bmatrix}$$
$$\begin{bmatrix} x_1(0) \\ x_2(0) \end{bmatrix} = \begin{bmatrix} 0 \\ 1 \end{bmatrix}.$$

Solve this system using the formula

$$\underline{\boldsymbol{x}}(t) = e^{tA} \left(\underline{\boldsymbol{x}}_0 + \int_0^t e^{-sA} \boldsymbol{f}(s) \, ds \right)$$

(One could also try undetermined coefficients, but variation of parameters requires no "guessing.")

Tech check: (The commands are sort of strange, but might help in your homework.)

$$A := \begin{bmatrix} 3 & -1 \\ 1 & 1 \end{bmatrix} :$$

with(LinearAlgebra):
$$A := \begin{bmatrix} 3 & -1 \\ 1 & 1 \end{bmatrix}$$
:
$$MatrixExponential(t \cdot A);$$

$$f := t \rightarrow \begin{bmatrix} t \\ 0 \end{bmatrix}$$
:
$$x0 := \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$
:

$$\begin{bmatrix} e^{2t} (1+t) & -t e^{2t} \\ t e^{2t} & -e^{2t} (-1+t) \end{bmatrix}$$
 (5)

integrand := $s \rightarrow simplify(MatrixExponential(-s \cdot A).f(s))$: #integrand in formula above integrand(t); #checking

$$\begin{bmatrix} -e^{-2t} (-1+t) t \\ -t^2 e^{-2t} \end{bmatrix}$$
 (6)

integrated := unapply(map(int, integrand(s), s = 0..t), t): #"map" applies a function to each entry of an array... # "unapply" makes a function out of output

integrated(t); #checking

$$\begin{bmatrix} \frac{1}{2} t^2 e^{-2t} \\ -\frac{1}{4} + \frac{1}{4} e^{-2t} + \frac{1}{2} t e^{-2t} + \frac{1}{2} t^2 e^{-2t} \end{bmatrix}$$
 (7)

 $x := unapply(simplify(MatrixExponential(t \cdot A).(x0 + integrated(t))), t)$: x(t); #checking answer

$$\begin{bmatrix} \frac{1}{4} t (e^{2t} - 1) \\ \frac{1}{4} e^{2t} (-1 + t) + \frac{1}{4} t + \frac{1}{4} \end{bmatrix}$$
(8)

 $dsolve(\{xl'(t) = 3 \cdot xl(t) - x2(t) + t, x2'(t) = xl(t) + x2(t), xl(0) = 0, x2(0) = 0\});$

$$\left\{ xI(t) = \frac{1}{4} t e^{2t} - \frac{1}{4} t, x2(t) = -\frac{1}{4} e^{2t} + \frac{1}{4} t + \frac{1}{4} t + \frac{1}{4} t e^{2t} \right\}$$
 (9)

Monday
$$4/3$$

5.6.39
$$A = \begin{bmatrix} 1 & 3 & 3 & 3 \\ 0 & 1 & 3 & 3 \\ 0 & 0 & 2 & 3 \\ 0 & 0 & 0 & 2 \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 3 & 3 \\ 0 & 0 & 0 & 3 \\ 0 & 0 & 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 & 338 \\ 0 & 0 & 338 \\ 0 & 0 & 0 & 3 \\ 0 & 0 & 0 & 0 \end{bmatrix}$$

I told you this in class, but it's wrong. N&D do not commune so the rule of exponents does not work. Have to resort to generalized ergenspaces

Math 2280-001 Monday April 3

In some sense the examples/methods in this problem contain pretty much all of the course so 1a1.

w12.1 (from next week's homework) Consider the input-output IVP from the second midtern

Transday.

$$\begin{bmatrix} x_1'(t) \\ x_2'(t) \end{bmatrix} = \begin{bmatrix} -4 & 1 \\ 2 & -3 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} + \begin{bmatrix} 20 \\ 40 \end{bmatrix}$$
$$\begin{bmatrix} x_1(0) \\ x_2(0) \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}.$$

 $E_{\lambda=-2} = \operatorname{Span}\left\{\begin{bmatrix} 1 \\ 2 \end{bmatrix}\right\}$ $E_{\lambda=-5} = \operatorname{Span}\left\{\begin{bmatrix} 1 \\ -1 \end{bmatrix}\right\}$

w12.1a) Using matrix exponentials, the solution to

$$\underline{\boldsymbol{x}}'(t) = A\,\underline{\boldsymbol{x}} + \boldsymbol{f}(t)$$
$$\underline{\boldsymbol{x}}(0) = \underline{\boldsymbol{x}}_0$$

is given by

$$\underline{\boldsymbol{x}}(t) = e^{tA} \left(\underline{\boldsymbol{x}}_0 + \int_0^t e^{-sA} \boldsymbol{f}(s) ds\right).$$

Verify that this recovers the correct solution

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} 10 \\ 20 \end{bmatrix} - 10 e^{-2t} \begin{bmatrix} 1 \\ 2 \end{bmatrix}$$

 $\underline{\mathbf{w12.1b}}$) Alternately, when A is diagonalizable, as it is in this case, we can change output variables to reduce to a simpler system: Let

$$AS = S\Lambda$$

where the columns of S are an eigenbasis for \mathbb{R}^n or \mathbb{C}^n and Λ is the diagonal matrix of corresponding eigenvalues. Then

$$\underline{\boldsymbol{x}}'(t) = A\,\underline{\boldsymbol{x}} + \boldsymbol{f}(t)$$
$$\underline{\boldsymbol{x}}(0) = \underline{\boldsymbol{x}}_0$$

is equivalent to

$$S^{-1}\underline{\boldsymbol{x}}' = S^{-1}A\,\underline{\boldsymbol{x}} + S^{-1}\boldsymbol{f}(t).$$

For

$$\underline{\underline{\boldsymbol{u}}}(t) := S^{-1}\underline{\boldsymbol{x}}(t)$$
$$\underline{\boldsymbol{u}}'(t) = S^{-1}\underline{\boldsymbol{x}}'(t)$$

this yields the diagonal system

$$\underline{\boldsymbol{u}}'(t) = S^{-1}AS\,\underline{\boldsymbol{u}}(t) + S^{-1}\boldsymbol{f}(t)$$

i.e. the IVP

$$\underline{\boldsymbol{u}}'(t) = \Lambda \, \underline{\boldsymbol{u}}(t) + S^{-1} \boldsymbol{f}(t)$$
$$\underline{\boldsymbol{u}}(0) = S^{-1} \underline{\boldsymbol{x}}_0$$

After finding $\underline{\boldsymbol{u}}(t)$, the original solution $\underline{\boldsymbol{x}}(t) = S \underline{\boldsymbol{u}}(t)$. Rework the tank IVP using this method.

Wed Apr 5

7.1-7.2 Laplace transform, and application to DE IVPs

- The Laplace transform is a linear transformation " \mathcal{L} " that converts piecewise continuous functions f(t), defined for $t \ge 0$ and with at most exponential growth $(|f(t)| \le Ce^{Mt})$ for some values of C and M, into functions F(s) defined by the transformation formula f(s) is f(s). inputs $F(s) = \mathcal{L}\{f(t)\}(s) := \int_0^\infty f(t)e^{-st} dt$.
- Notice that the integral formula for F(s) is only defined for sufficiently large s, and certainly for s > M, because as soon as s > M the integrand is decaying exponentially, so the improper integral from t = 0 to ∞ converges.
- The convention is to use lower case letters for the input functions and (the same) capital letters for their Laplace transforms, as we did for f(t) and F(s) above. Thus if we called the input function x(t) then we would denote the Laplace transform by X(s).

Taking Laplace transforms seems like a strange thing to do. And yet, the Laplace transform \mathcal{L} is just one example of a collection of useful "integral transforms". \mathcal{L} is especially good for solving IVPs for linear DEs, as we shall see starting today. Other famous transforms - e.g. Fourier series and Fourier transform are extremely important in studying linear differential and partial differential equations. We will discuss Fourier series in about a week. These transforms are also studied in Math 3140, 3150, and in various 5000-level pure and applied math classes.

Exercise 1) Use the definition of Laplace transform

$$F(s) = \mathcal{L}\{f(t)\}(s) := \int_0^\infty f(t)e^{-st} dt$$

to check the following facts, which you will also find inside the front cover of your text book.

$$\underline{\mathbf{a}} \ \mathcal{L}\{I\}(s) = \frac{1}{s} \ (s > 0)$$

$$\underline{b} \mathcal{L}\left\{e^{\alpha t}\right\}(s) = \frac{1}{s - \alpha} \ (s > \alpha \text{ if } \alpha \in \mathbb{R}, s > a \text{ if } \alpha = a + k i \in \mathbb{C})$$

c) Laplace transform is linear, i.e.

$$\begin{split} \mathcal{L}\big\{f_1(t) + f_2(t)\big\}(s) &= F_1(s) + F_2(s) \;. \\ \mathcal{L}\big\{cf(t)\big\}(s) &= c\,F(s). \end{split}$$

$$\mathcal{L}\{f_{1}(t)+f_{2}(t)\}(s)=F_{1}(s)+F_{2}(s).$$

$$\mathcal{L}\{cf(t)\}(s)=cF(s).$$
d) Use linearity and your work above to compute $\mathcal{L}\{3-4e^{-2t}\}(s).$
a) $\mathcal{L}\{1\}=\int_{0}^{\infty} e^{-st} dt = e^{-st} = 0 - (\frac{1}{-s}) = \frac{1}{s}$ (5) 0)
b) $\mathcal{L}\{e^{st}\}(s)=\int_{0}^{\infty} e^{st} e^{-st} dt = \int_{0}^{\infty} e^{(\alpha-s)t} dt = \frac{(\alpha-s)t}{\alpha-s} = 0 - \frac{1}{\alpha-s} = \frac{1}{s-\alpha}$
c) $\mathcal{L}\{c_{1}f_{1}(t)+c_{2}f_{2}(t)\}(s)=\int_{0}^{\infty} (c_{1}f_{1}(t)+c_{2}f_{2}(t))e^{-st} dt = c_{1}F_{1}(s)+c_{2}F_{2}(s)$ because in frequency of some superscript is time.

For the linear differential equations and systems of differential equations that we've been studying, the following Laplace transforms are very important:

Exercise 2) Use complex number algebra, including Euler's formula, linearity, and the result from 1b that

$$\mathcal{L}\left\{e^{(a+ki)t}\right\}(s) = \frac{1}{s - (a+ki)}$$

to verify that

a)
$$\mathcal{L}\{\cos(kt)\}(s) = \frac{s}{s^2 + k^2}$$

b) $\mathcal{L}\{\sin(kt)\}(s) = \frac{k}{s^2 + k^2}$
c) $\mathcal{L}\{e^{at}\cos(kt)\}(s) = \frac{s - a}{(s - a)^2 + k^2}$
d) $\mathcal{L}\{e^{at}\sin(kt)\}(s) = \frac{k}{(s - a)^2 + k^2}$

(Notice that if we tried doing these Laplace transforms directly from the definition, the integrals would be messy but we could attack them via integration by parts or integral tables.)

It's a theorem (hard to prove but true) that a given Laplace transform F(s) can arise from at most one piecewise continuous function f(t). (Well, except that the values of f at the points of discontinuity can be arbitrary, as they don't affect the integral used to compute F(s).) Therefore you can read Laplace transform tables in either direction, i.e. not only to deduce Laplace transforms, but inverse Laplace transforms $\mathcal{L}^{-1}\{F(s)\}(t) = f(t)$ as well.

Exercise 3) Use the Laplace transforms we've computed and linearity to compute

$$\mathcal{L}^{-1} \left\{ \frac{7}{s} + \left(\frac{1}{s^2 + 16} \right) - \frac{10 \, s}{s^2 + 16} \right\} (t) \, .$$

$$= 7 + \frac{1}{4} \sin 4t - 10 \cos 4t$$

L is 1-1.
only for -> 0 is f=0
proof is fun"
"Stone-Weirestrasce
The"

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$ f(t) \le Ce^{Mt}$	$F(s) := \int_0^\infty f(t)e^{-st} dt$ for $s > M$			
$c_1 f_1(t) + c_2 f_2(t)$	$c_1 F_1(s) + c_2 F_2(s)$	/		
1	$\frac{1}{s} (s > 0)$	V		
e ^{a.t}	$\frac{1}{s-\alpha} \qquad (s > \Re(a))$	V		
$\cos(k t)$	$\frac{s}{s^2 + k^2} (s > 0)$	V		
sinkt sin(kt)	$\frac{1}{s^2 + k^2} \qquad \frac{k}{s^2 + k^2} (s > 0)$	\ <u>\</u>	4 52+16	k= 4
$e^{a t} \cos(k t)$	$\frac{(s-a)}{(s-a)^2 + k^2} (s > a)$	V		
$e^{at}\sin(kt)$	$\frac{k}{(s-a)^2 + k^2} (s > a)$	\rightarrow		
f'(t)	s F(s) - f(0)	V		
f''(t)	$s^2F(s) - sf(0) - f'(0)$	~		

Laplace transform table
$$\int_{0}^{\infty} f''(t) dt = \int_{0}^{\infty} f'(s) - \int_{0}^{\infty} f'(s) -$$

The integral transforms of DE's and PDE's were designed to have the property that they convert the corresponding linear DE and PDE problems into algebra problems. For the Laplace transform it's because of these facts:

Exercise 4a) Use integration by parts and the definition of Laplace transform to show that

$$\mathcal{L}\lbrace g'(t)\rbrace(s) = s \,\mathcal{L}\lbrace g(t)\rbrace(s) - g(0) = s \,G(s) - g(0) . \quad \blacksquare$$

<u>4b)</u> Use the result of \underline{a} , applied to the function f'(t) to show that

$$\mathcal{L}\{f''(t)\}(s) = s^{2}F(s) - sf(0) - f'(0).$$
4c) What would you guess is the Laplace transform of $f'''(t)$? Could you check this?

$$\mathbf{L}\{g'(t)\}(s) = \int_{0}^{\infty} g'(t) e^{-st} dt \qquad = e^{-st} g(t) \int_{0}^{\infty} - \int_{0}^{\infty} g(t) e^{-st} dt \qquad = e^{-st} dt \qquad = \int_{0}^{\infty} g(t) e^{-st}$$

Here's an example of using Laplace transforms to solve DE IVPs, in the context of Chapter 3 and the mechanical (and electrical) application problems we considered there.

Exercise 5) Consider the undamped forced oscillation IVP

$$x''(t) + 4x(t) = 10\cos(3t)$$

 $x(0) = 2$
 $x'(0) = 1$

If x(t) is the solution, then both sides of the DE are equal. Thus the Laplace transforms are equal as well... .so, equate the Laplace transforms of each side and use algebra to find $\mathcal{L}\{x(t)\}(s) = X(s)$. Notice you've computed X(s) without actually knowing x(t)! If you were happy to stay in "Laplace land" you'd be done. In any case, at this point you can use our table entries to find $x(t) = \mathcal{L}^{-1}\{x(t)\}(s)$.

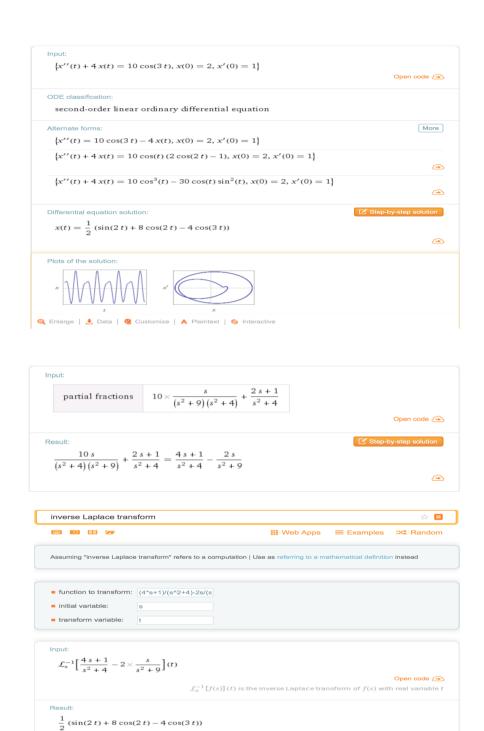
(Notice that if your algebra skills are good you've avoided having to use the Chapter 3 algorithm of (i) find x_H (ii) find an x_P (iii) $x = x_P + x_H$ (iv) solve IVP.) Magic! Or, would you have prefered to convert to a

first order system and to have used variation of parameters with an FM or e^{At} , like in Chapter 5:-)?)

respectively. Solve 171.) Magnetic of, would you have prefered to convert to a system and to have used variation of parameters with an FM or
$$e^{At}$$
, like in Chapter 5:-)?)

for soll $x(t)$, $X(LHS) = X(RHS)$ because $LHS = RHS$

$$\begin{array}{c}
f(t) & f(t) \\
\hline
x'' & f$$



Exercise 6) Use Laplace transform as above, to solve the IVP for the following underdamped, unforced oscillator DE:

The content is above, to solve the type for the following underdamped, unforced is:

$$x''(t) + 6x'(t) + 34x(t) = 0$$

$$x(0) = 3$$

$$x''(0) = 1$$

$$x'' + 5x'(5) - x(0)$$

$$x'' + 5x'(5$$

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7.1-7.4 Laplace transform, and application to DE IVPs, especially those in Chapter 3. Today we'll continue to fill in the Laplace transform table (at the end of the notes). Along the way we'll revisit some of the mechanical oscillation differential equations from Chapter 3.

Exercise 1) (to review) Use the table to compute

$$\begin{array}{ccc} \underline{1a} & \mathcal{L}\left\{4 + 2e^{-4t}\right\}(s) \\ \underline{1b} & C^{-1} & 2 & 6 \end{array}$$

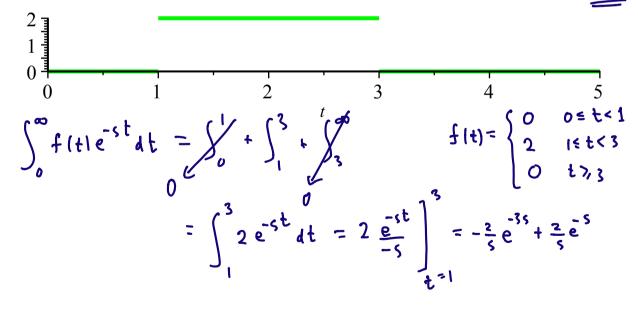
$$\underline{1b}) \quad \mathcal{L}^{-1}\left\{\frac{2}{s-2}+\frac{6}{s}\right\}(t) \ .$$

$$\frac{4}{5} + 2 \frac{1}{5+4}$$

Exercise 2) (to review the definition) Use the definition of Laplace transform,

$$F(s) = \mathcal{L}\{f(t)\}(s) := \int_0^\infty f(t)e^{-st} dt$$

to find the Laplace transform of the step function graphed below. (The function is equal to zero for $t \ge 3$.)



Exercise 3a) Use the Table entry we proved on Wednesday for derivatives (via integration by parts), namely

¥
$$\mathcal{L}\{g'(t)\}(s) = s \mathcal{L}\{g(t)\}(s) - g(0) = s G(s) - g(0)$$

and math induction to show that for $n \in \mathbb{N}$

$$\mathcal{L}\left\{f^{(n)}(t)\right\}(s) = s^{n} F(s) - s^{n-1} f(0) - s^{n-2} f'(0) - \dots - s f^{(n-2)}(0) - f^{(n-1)}(0) .$$

$$h = 2 : \mathcal{L}\left\{f''(t)\right\}(s) = s^{2} F(s) - s f(0) - f'(0)$$

• know for n=1 * on Wed. • Show if it's true for n=k than it's true for n=k+1 true $\forall n=1,2,...$ • Show if it's true for n=k than it's true for n=k+1 $n \in \mathbb{N}$. Assume $\sum_{k=1}^{n} f(k) = \sum_{k=1}^{n} f(k) =$

$$= s^{k+1} F(s) - s^{k} f(0) - s^{k-1} f'(0) - \cdots - s^{k} f^{(k-1)} - f^{(k)}(0)$$

<u>3b</u>) (Integrals are "negative" derivatives): Use the Laplace transform first-derivative formula above to show that

state degree derivatives). Ose the Laplace transform instructivative formula
$$\mathcal{L}\left\{\int_{0}^{t} f(\tau) d\tau\right\}(s) = \frac{1}{s} \mathcal{L}\left\{f(t)\right\}(s) = \frac{F(s)}{s} \dots$$

$$\mathcal{L}\left\{\int_{0}^{t} \left(\int_{0}^{r} f(\tau) d\tau\right) dr\right\}(s) = \frac{F(s)}{s^{2}} \dots$$

$$\mathcal{L}\left\{\int_{0}^{t} \left(\int_{0}^{r} f(\tau) d\tau\right) dr\right\}(s) = \frac{F(s)}{s} \dots$$

$$\mathcal{L}\left\{\int_{0}^{t} \left(\int_{0}^{r} f(\tau) d\tau\right) dr\right\}(s) = \frac{F(s)}{s} \dots$$

$$\mathcal{L}\left\{\int_{0}^{t} \left(\int_{0}^{r} f(\tau) d\tau\right) dr\right\}(s) = \frac{F(s)}{s} \dots$$

$$\mathcal{L}\left\{\int_{0}^{t} \left(\int_{0}^{r} f(\tau) d\tau\right) d\tau\right\}(s) = \frac{F(s)}{$$

Exercise 4) Use the result of <u>3a</u> to show that for $n \in \mathbb{N}$,

$$\mathcal{L}\{t\}(s) = \frac{1}{s^2}$$

$$\mathcal{L}\{t^2\}(s) = \frac{2}{s^3}$$

$$\mathcal{L}\{t^n\}(s) = \frac{n!}{s^{n+1}}.$$

$$f(t) = t^{n}$$

$$f'(t) = n t^{n-1}$$

$$f''(t) = n (n-1) t^{n-2}$$

$$\vdots$$

$$f^{(n)} = n (n-1) - 3 \cdot 2 \cdot 1$$

$$f(t) = t^{n}$$

$$f^{(n)}(t) = n! \quad \exists$$

$$so \quad \chi \left\{ n! \right\} = s^{n} F(s) - s^{n-1} f(s) - s^{n-2} f'(s) - \dots - f^{(n-n)}(s)$$

$$\frac{n!}{s} = s^{n} F(s) \qquad \frac{f(t)}{s} F(s)$$

$$F(s) = \frac{n!}{s^{n+1}}$$

$$sinkt \quad \frac{k}{s^{2}+1}$$

$$\frac{f(t) F(s)}{1 \frac{1}{5}}$$

$$\sinh t \frac{h}{5^2 + h^2}$$

$$\cosh t \frac{s}{5^2 + h^2}$$

Exercise 5) Find $\mathcal{L}^{-1}\left\{\frac{1}{s(s^2+4)}\right\}(t)$

- a) using the result of <u>3b</u>.
- b) using partial fractions.

a)
$$\mathcal{L}\left\{\int_{s}^{t}f(t)dt\right\}(s)=\frac{F(s)}{s}$$

$$F(s) = \frac{1}{s^{2}+4} = \frac{1}{2} \frac{2}{s^{2}+4}$$

$$f(t) = \frac{1}{2} \sin 2t$$

$$\int_{0}^{t} \frac{1}{2} \sin 2t \, dt = -\frac{1}{4} \cos 2t + \frac{1}{4}$$

$$\int_{0}^{t} \frac{1}{4} \cos 2t + \frac{1}{4} \int_{0}^{t} (s) = \frac{1}{5} \frac{1}{(s^{2}+4)}$$

b)
$$\frac{1}{s(s^2+4)} = \frac{A}{s} + \frac{Bs+c}{s^2+4}$$

$$| = A(s^{2}+4)+(Bs+C)s = 4A + os + os^{2} (A+B)$$

$$\frac{1}{s(s^{2}+4)} = \frac{1}{4} \frac{1}{s} - \frac{1}{4} \frac{s}{s^{2}+4}$$

$$4A = 1 \quad A = \frac{1}{4}$$

$$B = -\frac{1}{4}$$

$$C = 0$$

Exercise 6) (first translation theorem). Use the definition of Laplace transform to show that

$$\mathcal{L}\left\{e^{at}f(t)\right\}(s) = \mathcal{L}\left\{f(t)\right\}(s-a) = F(s-a)$$

$$\int_{0}^{\infty} e^{at}f(t) e^{-st} dt$$

$$\int_{0}^{\infty} f(t) e^{-(s-a)t} dt$$

$$\mathcal{L}\left\{f(t)\right\}(s-a)$$

Exercise 7) As a special case of Exercise 6, show

$$\mathcal{L}\left\{t e^{a t}\right\}(s) = \frac{1}{\left(s - a\right)^{2}}.$$

$$\mathcal{L}\left\{t^{n} e^{a t}\right\}(s) = \frac{n!}{\left(s - a\right)^{n+1}}$$

A harder table entry to understand is the following one - go through this computation and see why it seems reasonable, even though there's one step that we don't completely justify. The table entry is

Here's how we get it:
$$F(s) = \mathcal{L}\{f(t)\}(s) := \int_{0}^{\infty} f(t)e^{-st} dt = \int_{0}^{\infty} \frac{d}{ds}f(t)e^{-st} dt.$$

$$\Rightarrow \frac{d}{ds}F(s) = \frac{d}{ds}\int_{0}^{\infty} f(t)e^{-st} dt = \int_{0}^{\infty} \frac{d}{ds}f(t)e^{-st} dt.$$

It's this last step which is true, but needs more justification. We know that the derivative of a sum is the sum of the derivatives, and the integral is a limit of Riemann sums, so this step does at least seem reasonable. The rest is straightforward:

$$\int_0^\infty \frac{d}{ds} f(t) e^{-st} dt = \int_0^\infty f(t) (-t) e^{-st} dt = -\mathcal{L} \{t f(t)\}(s) \qquad \Box.$$

For resonance and other applications ...

Exercise 8) Use $\mathcal{L}\{tf(t)\}(s) = -F'(s)$ directly, or Euler's formula and $\mathcal{L}\{te^{\alpha t}\}(s) = \frac{1}{(s-\alpha)^2}$ to

show

a)
$$\mathcal{L}\{t\cos(kt)\}(s) = \frac{s^2 - k^2}{(s^2 + k^2)^2}$$

$$\underline{b} \mathcal{L}\left\{\frac{1}{2k} t \sin(kt)\right\} (s) = \frac{s}{\left(s^2 + k^2\right)^2}$$

c) Then use a and the identity

$$\frac{1}{\left(s^2 + k^2\right)^2} = \frac{1}{2k^2} \left(\frac{s^2 + k^2}{\left(s^2 + k^2\right)^2} - \frac{s^2 - k^2}{\left(s^2 + k^2\right)^2} \right)$$

to verify the table entry

$$\mathcal{L}^{-1}\left\{\frac{1}{\left(s^2 + k^2\right)^2}\right\}(t) = \frac{1}{2k^2} \left(\frac{1}{k}\sin(kt) - t\cos(kt)\right).$$

a)
$$\chi \{ t \omega s h t \} (s) = -\frac{d}{ds} \chi \{ \omega s h t \} (s) = -\frac{d}{ds} \frac{s}{s^2 + k^2} = -\frac{1 \cdot (s^2 + k^2)}{(s^2 + k^2)^2} = \frac{5^2 - k^2}{(s^2 + k^2)^2}$$

Exercise 9) Use Laplace transforms to write down the solution to

$$x''(t) + \omega_0^2 x(t) = \underline{F_0} \sin(\omega_0 t)$$
$$x(0) = x_0$$
$$x'(0) = v_0.$$

what phenomena do solutions to this DE illustrate (even though we're forcing with $\sin(\omega_0 t)$ rather than $\cos(\omega_0 t)$)? How would you have tried to solve this problem in Chapter 3?

$$\mathcal{Y}: \zeta^{2} \chi(\varsigma) - \varsigma \chi_{\bullet} - v_{o} + \omega_{\bullet}^{2} \chi(\varsigma) = \frac{F_{o}}{m} \frac{\omega_{o}}{\varsigma^{2} + \omega_{\bullet}^{2}}$$

$$\chi(\varsigma) \left(\varsigma^{2} + \omega_{\bullet}^{2}\right) = \frac{F_{o}}{m} \frac{\omega_{o}}{\varsigma^{2} + \omega_{\bullet}^{2}} + \chi_{\bullet} \varsigma + v_{o}$$

$$\chi(\varsigma) = \frac{F_{o} \omega_{o}}{m} \frac{1}{(\varsigma^{2} + \omega_{\bullet}^{2})^{2}} + \chi_{o} \frac{\varsigma}{\varsigma^{2} + \omega_{\bullet}^{2}} + v_{o} \frac{1}{\varsigma^{2} + \omega_{\bullet}^{2}}$$

$$\chi(t) = \frac{F_{o} \omega_{o}}{m} \frac{1}{2\omega_{\bullet}^{2}} \left(\frac{1}{\omega_{o}} \sin \omega_{o}^{2} t - t \omega_{o} \omega_{o}^{2} t\right) + \chi_{o} \cos \omega_{o}^{2} t + \frac{v_{o}}{\omega_{o}^{2}} \sin \omega_{o}^{2} t$$

$$\chi(t) = \frac{F_{o} \omega_{o}}{m} \frac{1}{2\omega_{o}^{2}} \left(\frac{1}{\omega_{o}} \sin \omega_{o}^{2} t - t \omega_{o} \omega_{o}^{2} t\right) + \chi_{o} \cos \omega_{o}^{2} t + \frac{v_{o}}{\omega_{o}^{2}} \sin \omega_{o}^{2} t$$

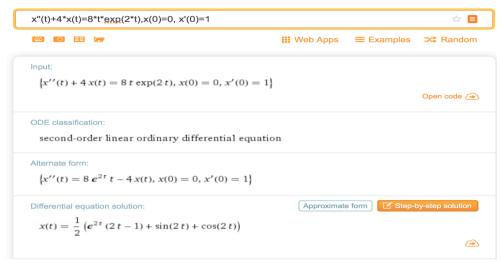
Exercise 10) Solve the following IVP. Use this example to recall the general partial fractions algorithm.

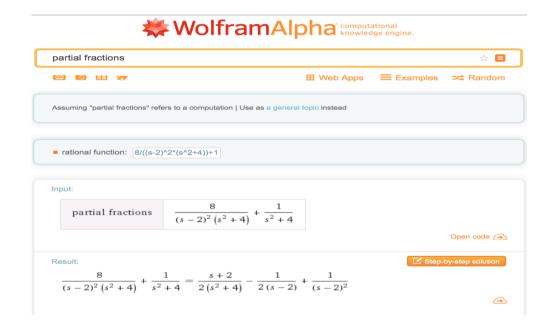
$$x''(t) + 4x(t) = 8 t e^{2t}$$

 $x(0) = 0$
 $x'(0) = 1$

Wolfram checks:









			₩eb Apps	E xamples	⊃⊄ Rando
Assuming "in	verse laplace transform	m" refers to a computation Use as ref	erring to a mathe	matical definition inst	tead
■ function to	o transform: (1/2)s	/(s^2+4)+1/(s^2			
■ initial vari	able: s				
■ transform	variable: t				
Input: $\mathcal{L}_s^{-1} \left[\frac{1}{2} \right]$	$\left(\frac{s}{s^2+4}+\frac{1}{s^2+4}\right)$	$\left[1 - \frac{1}{2} \times \frac{1}{s-2} + \frac{1}{(s-2)^2}\right](t)$			
- 2	3 + 4 3 + 4	2 3 - 2 (3 - 2)-3			Open code
		$\mathcal{L}_s^{-1}[f(s)](t)$ is the inver	rse Laplace trai	nsform of $f(s)$ with	h real variab
Result:					

$f(t)$, with $ f(t) \le Ce^{Mt}$	$F(s) := \int_0^\infty f(t)e^{-st} dt \text{ for } s > M$	↓ verified
$c_{1}f_{1}(t) + c_{2}f_{2}(t)$	$c_1 F_1(s) + c_2 F_2(s)$	Verified
1 t	$\frac{1}{s} \qquad (s > 0)$	₽
t^2	$\frac{1}{s^2}$ $\frac{2}{s^3}$	
$t^n, n \in \mathbb{N}$	$\frac{s^3}{n!}$ $\frac{n!}{s^{n+1}}$	
e ^{a.t}	$\frac{1}{s-\alpha} \qquad (s>\Re(\alpha))$	k
$\cos(k t)$	$\frac{s}{s^2 + k^2} (s > 0)$	>
$\sin(k t)$	$\frac{\frac{k}{k}}{s^2 + k^2} (s > 0)$	×
$\cosh(k t)$	$\frac{s}{s^2 - k^2} (s > k)$	
$\sinh(k t)$	$\frac{k}{s^2 - k^2} (s > k)$	
$e^{a t} \cos(k t)$	$\frac{\frac{(s-a)}{(s-a)^2 + k^2}}{\frac{k}{(s-a)^2 + k^2}} (s > a)$	M
$e^{at}\sin(kt)$	$\frac{k}{(s-a)^2 + k^2} (s > a)$	≥
$f'(t)$ $f''(t)$ $f^{(n)}(t), n \in \mathbb{N}$ $\int_0^t f(\tau) d\tau$	$s F(s) - f(0)$ $s^{2}F(s) - s f(0) - f'(0)$ $s^{n} F(s) - s^{n-1}f(0) - \dots - f^{(n-1)}(0)$ $\frac{F(s)}{s}$	X X Q
$t f(t)$ $t^{2} f(t)$ $t^{n} f(t), n \in \mathbb{Z}$	$ \begin{array}{c} -F'(s) \\ F''(s) \\ (-1)^n F^{(n)}(s) \end{array} $	

$\frac{f(t)}{t}$	$\int_s^\infty F(\sigma)d\sigma$	
$t\cos(k t)$ $\frac{1}{2 k} t \sin(k t)$ $\frac{1}{2 k^3} (\sin(k t) - k t \cos(k t))$	$\frac{s^2 - k^2}{(s^2 + k^2)^2}$ $\frac{s}{(s^2 + k^2)^2}$ $\frac{1}{(s^2 + k^2)^2}$	
$e^{at}f(t)$ $t e^{at}$	F(s-a)	
$t^n e^{at}, n \in \mathbb{Z}$	$\frac{\frac{1}{(s-a)^2}}{\frac{n!}{(s-a)^{n+1}}}$	

Laplace transform table