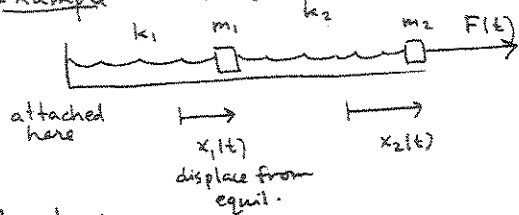


Math 2280-1  
Tuesday March 1

§ 4.1, 4.3

- finish damped oscillator re-interpretation as 1<sup>st</sup> order system, page 4 <sup>Monday</sup>
- spring system (and conversion below), then general discussion.

Example: mass-spring systems (§5.3)



Newton:

$$m_1 x_1''(t) = \underbrace{k_2 (x_2 - x_1)}_{\substack{\text{2nd spring is stretched} \\ \text{this much}}} - k_1 x_1$$

$$m_2 x_2''(t) = -k_2 (x_2 - x_1) + F(t)$$

e.g.  $\left. \begin{array}{l} m_1 = 2 \\ m_2 = 1 \\ k_1 = 4 \\ k_2 = 2 \\ f(t) = 40 \sin 3t \end{array} \right\} \Rightarrow \left\{ \begin{array}{l} 2x_1'' = 2(x_2 - x_1) - 4x_1 = -6x_1 + 2x_2 \\ x_2'' = -2(x_2 - x_1) + 40 \sin 3t = 2x_1 - 2x_2 + 40 \sin 3t \end{array} \right.$

or  $\begin{bmatrix} x_1'' \\ x_2'' \end{bmatrix} = \begin{bmatrix} -3 & 1 \\ 2 & -2 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} + \begin{bmatrix} 0 \\ 40 \sin 3t \end{bmatrix}$

- What's a natural IVP for this 2<sup>nd</sup> order system of 2 linear DE's?  
(We won't solve it today - see §7.4)

note: Any system of DE's can be converted into an equivalent system of 1<sup>st</sup> order DE's, by introducing extra unknown functions.

(1)  $\begin{cases} x_1'' = -3x_1 + x_2 \\ x_2'' = 2x_1 - 2x_2 + 40 \sin 3t \end{cases}$

Let  $y_1 = x_1'$   
 $y_2 = x_2'$

(2)  $\begin{cases} x_1' = y_1 \\ y_1' = -3x_1 + x_2 \\ x_2' = y_2 \\ y_2' = 2x_1 - 2x_2 + 40 \sin 3t \end{cases}$

solutions of (1) yield solutions of (2) and vice versa. CHECK!!

This correspondence is sometimes useful.

We'll solve systems like this in §5.3

Def A 1<sup>st</sup> order system of DE's of the form

$$* \quad \frac{d\vec{x}}{dt} - A(t)\vec{x} = \vec{f}(t) \quad A(t) \text{ an } n \times n \text{ matrix}$$

is called a linear 1<sup>st</sup> order system.

If  $\vec{f}(t) \equiv 0$  it is called homogeneous

Notice (check!)

$$L(\vec{x}(t)) := \frac{d\vec{x}}{dt} - A(t)\vec{x}(t) \quad \text{is } \underline{\text{linear}} \quad \left( \begin{array}{l} \text{domain} = \text{continuously diffble vector valued fns } \vec{x}(t) \\ \text{codomain} = \text{cont. vector valued fns} \end{array} \right)$$

Corollary: The general sol'n to \* is

$$\vec{x}(t) = \vec{x}_H(t) + \vec{x}_p(t)$$

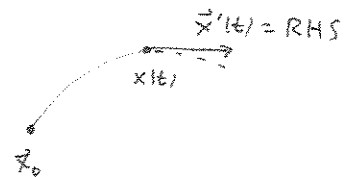
where  $\vec{x}_p(t)$  is a particular sol'n to \* &  $\vec{x}_H(t)$  is the general sol'n to the homogeneous DE

$$\frac{d\vec{x}}{dt} = A(t)\vec{x}.$$

Theorem If  $A(t)$  and  $\vec{f}(t)$  are defined and continuous on an interval  $I$ , and  $t_0 \in I$ , then the IVP

$$\begin{cases} \frac{d\vec{x}}{dt} = A(t)\vec{x} + \vec{f}(t) \\ \vec{x}(t_0) = \vec{x}_0 \end{cases}$$

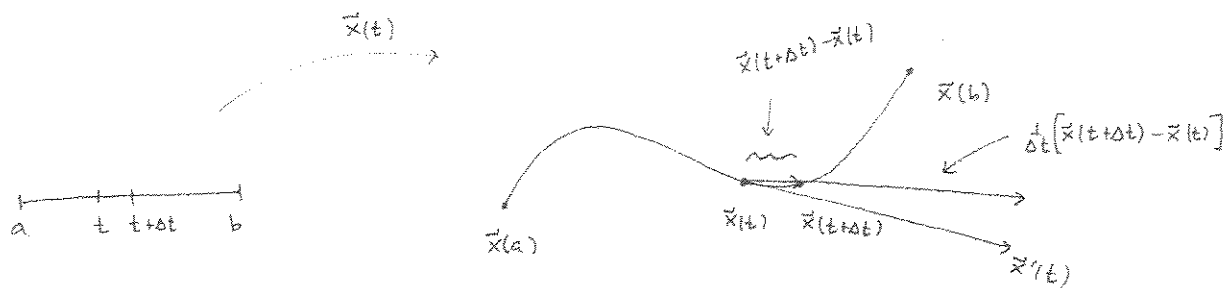
always has a unique sol'n on  $I$



proof: omitted, but the idea is you know where to start at time  $t_0$ , and you know your tangent vector  $\vec{x}'(t)$  if you know where you are and what time it is, so you know where to go.

The linear nature of this DE turns out to guarantee local (in time)  $\exists!$ , and that the solution can't blow up in finite time, so that it exists  $\forall t \in I$

Digression/review (2210 or 1260), about the meaning of  $\vec{x}'(t)$



$\vec{x}: [a, b] \rightarrow \mathbb{R}^n$   
domain codomain

Thus the first order IVP solution to

$$\begin{cases} \frac{d\vec{x}}{dt} = \vec{F}(t, \vec{x}) \\ \vec{x}(t_0) = \vec{x}_0 \end{cases}$$

can be interpreted as ~~searching~~ the particle motion starting at  $\vec{x}_0$ , for which the "velocity" vector is known, in terms of current position and time, via  $\vec{F}(t, \vec{x})$ .

$$\begin{aligned} \vec{x}'(t) &:= \lim_{\Delta t \rightarrow 0} \frac{1}{\Delta t} (\vec{x}(t+\Delta t) - \vec{x}(t)) \\ &= \lim_{\Delta t \rightarrow 0} \frac{1}{\Delta t} \begin{bmatrix} x_1(t+\Delta t) - x_1(t) \\ x_2(t+\Delta t) - x_2(t) \\ \vdots \\ x_n(t+\Delta t) - x_n(t) \end{bmatrix} \\ &= \lim_{\Delta t \rightarrow 0} \begin{bmatrix} \frac{x_1(t+\Delta t) - x_1(t)}{\Delta t} \\ \vdots \\ \frac{x_n(t+\Delta t) - x_n(t)}{\Delta t} \end{bmatrix} \\ &= \begin{bmatrix} x_1'(t) \\ x_2'(t) \\ \vdots \\ x_n'(t) \end{bmatrix} \end{aligned}$$

We called  $\vec{x}'(t)$  the tangent vector or velocity vector (if  $\vec{x}(t)$  was particle motion).

Cor 1 The solution space to the homogeneous system

(4)

$$\frac{d\vec{x}}{dt} - A(t)\vec{x} = \vec{0}$$

is  $n$ -dim'l

Proof: Let  $V$  be the solution space to this system,  
which is a subspace of differentiable vector-valued functions  $\vec{r}(t)$   
defined for  $t \in I$  (because  $V$  is closed under addition & scalar multiplication).

Let  $t_0 \in I$ .

Define  $T: V \rightarrow \mathbb{R}^n$  by  $T(\vec{x}(t)) := \vec{x}(t_0) \in \mathbb{R}^n$ .

$T$  is linear, and  
is an isomorphism

(1-1 & onto) by the  $\exists!$  theorem. Thus  $\dim(V) = \dim(\mathbb{R}^n) = n$  ■

Def If  $\vec{y}_1(t), \dots, \vec{y}_n(t)$  are vector-valued functions with image in  $\mathbb{R}^n$

then  $\begin{bmatrix} \vec{y}_1 & \vec{y}_2 & \dots & \vec{y}_n \end{bmatrix}$  is the Wronskian matrix  
and its det, called  $W(\vec{y}_1, \dots, \vec{y}_n)$  is  
the Wronskian,

Theorem: If  $\{\vec{y}_1, \dots, \vec{y}_n\}$  each solve  $\frac{d\vec{x}}{dt} - A(t)\vec{x} = \vec{0}$  on  $I$

then either  $W(\vec{y}_1, \dots, \vec{y}_n) \equiv 0$  on  $I$

or  $W(\vec{y}_1, \dots, \vec{y}_n) \neq 0 \quad \forall t \in I$ .

The first case means  $\{\vec{y}_1, \dots, \vec{y}_n\}$  are dependent, the second case means they're independent.

proof: If  $\det \begin{bmatrix} \vec{y}_1 & \vec{y}_2 & \dots & \vec{y}_n \end{bmatrix} \neq 0$  at  $t_0$  then every

IVP at  $t_0$  has a unique soltn  $\vec{x}(t) = c_1 \vec{y}_1 + c_2 \vec{y}_2 + \dots + c_n \vec{y}_n$

( $\vec{c}$  solves  $\begin{bmatrix} \vec{y}_1(t_0) & \dots & \vec{y}_n(t_0) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \\ \vdots \\ c_n \end{bmatrix} = \vec{b}$  where  $\vec{b}$  is the given  
initial value vector  $\vec{x}(t_0)$ )

In this case  $\{\vec{y}_1, \vec{y}_2, \dots, \vec{y}_n\}$

are a basis for the soltn space  $V$ , so for any (other)  $t_1 \in I$ ,

$\text{span}\{\vec{y}_1(t_1), \vec{y}_2(t_1), \dots, \vec{y}_n(t_1)\} = \mathbb{R}^n$  because every IVP at  $t_1$  has  
a soltn  $\vec{x}(t) = c_1 \vec{y}_1 + c_2 \vec{y}_2 + \dots + c_n \vec{y}_n$ .

thus  $\det \begin{bmatrix} \vec{y}_1 & \vec{y}_2 & \dots & \vec{y}_n \end{bmatrix} \neq 0$  at  $t_1$  as  
well. ■

Cor 2 (the real reason  $\exists!$  theorem holds for  $n^{\text{th}}$  order linear DE IVP.)

(5)

$$\text{Let } L(y) := y^{(n)} + a_{n-1}(t)y^{(n-1)} + \dots + a_1(t)y' + a_0(t)y, \quad a_j \in C(I)$$

$$f \in C(I)$$

$$\text{Let } t_0 \in I, \quad \vec{b} = [b_0, b_1, \dots, b_{n-1}] \in \mathbb{R}^n$$

Then  $\exists!$  soltn  $y(t)$  to IVP:  $y \in C^n(I)$ .

$$\text{IVP1} \begin{cases} L(y) = f & \text{in } I \\ y(t_0) = b_0 \\ y'(t_0) = b_1 \\ \vdots \\ y^{(n-1)}(t_0) = b_{n-1} \end{cases}$$

proof: This IVP is equivalent to the first order system IVP for  $\vec{x}(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \\ \vdots \\ x_n(t) \end{bmatrix}$ :

$$\text{IVP2} \begin{cases} x_1' = x_2 \\ x_2' = x_3 \\ \vdots \\ x_n' = f - a_{n-1}x_{n-1} - a_{n-2}x_{n-2} - \dots - a_0x_1 \\ x_1(t_0) = b_0 \\ x_2(t_0) = b_1 \\ \vdots \\ x_n(t_0) = b_{n-1} \end{cases}$$

$$\vec{x}' = \begin{bmatrix} 0 & 1 & 0 & \dots & 0 \\ 0 & 0 & 1 & \dots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \dots & 1 & 0 \\ -a_0 & -a_1 & \dots & -a_{n-1} & 0 \end{bmatrix} \vec{x} + \begin{bmatrix} 0 \\ 0 \\ \vdots \\ 0 \\ f \end{bmatrix}$$

i.e.  $\exists!$   $y$  soltn to IVP 1

iff

$$\begin{matrix} y = x_1 \\ y' = x_2 \\ \vdots \\ y^{(n-1)} = x_n \end{matrix} \quad \text{is ! soltn to IVP2}$$

Thus Cor 2, the Chapter 3  $\exists!$  theorem,  
follows from the linear systems  $\exists!$  theorem in Chapter 4 (page 1 theorem).