

Math 2250-4
Monday April 18

67.5 Defective eigen-spaces.

and how to find solution space to

$$\frac{d\vec{x}}{dt} = A\vec{x}$$

when A is not diagonalizable.

Example Find the solution space to

$$\frac{d\vec{x}}{dt} = \begin{bmatrix} 3 & 4 \\ 0 & 3 \end{bmatrix} \vec{x}$$

step 1 (as always).

$$\begin{vmatrix} 3-\lambda & 4 \\ 0 & 3-\lambda \end{vmatrix} = (3-\lambda)^2 = (\lambda-3)^2$$

$$\lambda = 3:$$

$$\begin{array}{c|c} 0 & 4 \\ 0 & 0 \end{array} \begin{array}{c} 0 \\ 0 \end{array}$$

eigenbasis $\left\{ \begin{bmatrix} 1 \\ 0 \end{bmatrix} \right\}$

$$v_2 = 0$$

$$\vec{v} = t \begin{bmatrix} 1 \\ 0 \end{bmatrix};$$

one sol'n $e^{3t} \begin{bmatrix} 1 \\ 0 \end{bmatrix}$; need another independent sol'n!

discussion: Let $A\vec{v} = \lambda_i \vec{v}$, but the λ_i -eigenspace is "defective", i.e. $\dim(\lambda_i\text{-eigenspace}) < k_i$ where charact poly has factor $(\lambda - \lambda_i)^{k_i}$.

Does multiplying by t work?

i.e. we know (for $\lambda = \lambda_i$)

$$\vec{x}(t) = e^{\lambda t} \vec{v} \text{ solves } \vec{x}' = A\vec{x}$$

does $\vec{z}(t) = t e^{\lambda t} \vec{v}$:

$$\vec{z}' = (e^{\lambda t} + t\lambda e^{\lambda t}) \vec{v}$$

$$A\vec{z} = t e^{\lambda t} A\vec{v} = t\lambda e^{\lambda t} \vec{v}$$

failure.

well, how about

$$\vec{z}(t) = e^{\lambda t} (t\vec{v} + \vec{w}) \text{ (same } \vec{v})$$

$$\vec{z}'(t) = [e^{\lambda t} + t\lambda e^{\lambda t}] \vec{v} + \lambda e^{\lambda t} \vec{w}$$

$$A\vec{z} = t\lambda e^{\lambda t} \vec{v} + e^{\lambda t} A\vec{w}$$

Success if

$$A\vec{w} = \lambda \vec{w} + \vec{v}$$

$$\begin{cases} (A - \lambda I) \vec{w} = \vec{v} \\ (A - \lambda I) \vec{v} = 0 \end{cases}$$

"a chain of length 2"

... last page ...

in our example

$$(A - \lambda I) \vec{w} = \vec{v}:$$

$$\begin{array}{c|c} 0 & 4 \\ 0 & 0 \end{array} \begin{array}{c} 1 \\ 0 \end{array}$$

$$4w_2 = 1 \text{ e.g. } \vec{w} = \begin{bmatrix} 0 \\ 1/4 \end{bmatrix}$$

$$\vec{z}(t) = e^{3t} \left(t \begin{bmatrix} 1 \\ 0 \end{bmatrix} + \begin{bmatrix} 0 \\ 1/4 \end{bmatrix} \right)$$

General sol'n:

$$\vec{x}_H(t) = c_1 e^{3t} \begin{bmatrix} 1 \\ 0 \end{bmatrix} + c_2 e^{3t} \left(t \begin{bmatrix} 1 \\ 0 \end{bmatrix} + \begin{bmatrix} 0 \\ 1/4 \end{bmatrix} \right)$$

Theorems: Let $A_{n \times n}$ as on page 1, with

$$|A - \lambda I| = (-1)^n (\lambda - \lambda_1)^{k_1} (\lambda - \lambda_2)^{k_2} \dots (\lambda - \lambda_m)^{k_m}$$

$$\begin{aligned} k_1 + \dots + k_m &= n \\ \lambda_i &\in \mathbb{C} \quad 1 \leq i \leq m \\ \lambda_i &\text{ distinct} \end{aligned}$$

(1) The generalized λ_i -eigenspace is defined to be

$$\ker (A - \lambda_i I)^{k_i}$$

it is always k_i dimensional (in $\mathbb{C}^n$)

(2) By amalgamating bases of the generalized eigenspaces you get a basis for $\mathbb{C}^n$

(3) The generalized λ_i -eigenbases can be chosen to be made of one or more "chains" $\{\vec{u}_1, \vec{u}_2, \dots, \vec{u}_k\}$ where for $\lambda = \lambda_i$

$$\begin{aligned} (A - \lambda I) \vec{u}_1 &= \vec{0} & (\vec{u}_1 \text{ is an eigenvector}) \\ (A - \lambda I) \vec{u}_2 &= \vec{u}_1 \\ &\vdots \\ (A - \lambda I) \vec{u}_k &= \vec{u}_{k-1} \end{aligned}$$

(see on page 1, $(\vec{v}, \vec{w})$ was a chain of length 2)

(4) Each chain $\{\vec{u}_1, \vec{u}_2, \dots, \vec{u}_k\}$ yields k lin. ind. solns to $\frac{d\vec{x}}{dt} = A\vec{x}$:

$$\vec{x}_1(t) = \vec{u}_1 e^{\lambda t}$$

$$\vec{x}_2(t) = (\vec{u}_1 t + \vec{u}_2) e^{\lambda t}$$

$$\vec{x}_3(t) = (\vec{u}_1 \frac{1}{2} t^2 + \vec{u}_2 t + \vec{u}_3) e^{\lambda t}$$

$\vdots$

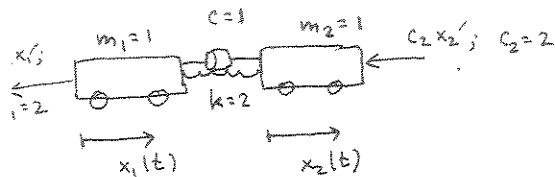
$$\vec{x}_k(t) = (\vec{u}_1 \frac{1}{k!} t^k + \dots + \vec{u}_{k-1} t + \vec{u}_k) e^{\lambda t}$$

$$\begin{aligned} \leftarrow \text{notice } \vec{x}_j' &= \lambda \vec{x}_j + \vec{x}_{j-1} \\ A \vec{x}_j &= \lambda \vec{x}_j + \vec{x}_{j-1} \text{ too!} \end{aligned}$$

amalgamating these

solutions yields a basis of sol'ns to $\frac{d\vec{x}}{dt} = A\vec{x}$!!!

Example 6 p 457



$$\begin{aligned} m_1 x_1'' &= +k(x_2 - x_1) - c_1 x_1' + c(x_2' - x_1') \\ m_2 x_2'' &= -k(x_2 - x_1) - c_2 x_2' - c(x_2' - x_1') \end{aligned}$$

$$\begin{bmatrix} x_1'' \\ x_2'' \end{bmatrix} = \begin{bmatrix} -2 & 2 \\ 2 & -2 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} + \begin{bmatrix} -3 & 1 \\ 1 & -3 \end{bmatrix} \begin{bmatrix} x_1' \\ x_2' \end{bmatrix}$$

cannot use §5.3 (no damping.)
Must convert to 1st order system

$$\begin{aligned} x_1 &= x_1 \\ x_2 &= x_2 \\ x_3 &= x_1' \\ x_4 &= x_2' \end{aligned} \quad \begin{bmatrix} x_1' \\ x_2' \\ x_3' \\ x_4' \end{bmatrix} = \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ -2 & 2 & -3 & 1 \\ 2 & -2 & 1 & -3 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \\ x_4 \end{bmatrix}$$

Maple to the rescue! (See next page).

$$t \ddot{u}_1 + \ddot{u}_2$$

$$\begin{bmatrix} x_1 \\ x_2 \\ x_1' \\ x_2' \end{bmatrix} = c_1 \begin{bmatrix} 1 \\ 1 \\ 0 \\ 0 \end{bmatrix} + e^{-2t} \left[c_2 \begin{bmatrix} 1 \\ 1 \\ -2 \\ -2 \end{bmatrix} + c_3 \begin{bmatrix} 1 \\ -1 \\ -2 \\ 2 \end{bmatrix} + c_4 \left[t \begin{bmatrix} -1 \\ 1 \\ 2 \\ -2 \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ -1 \\ 1 \end{bmatrix} \right] \right]$$

↑
explain.

↑
nice physically;
the sum of MAPLE's
2 $\lambda = -2$ eigenvects,
times -2.

↑
the difference
of Maple's
2 $\lambda = -2$
eigenvects, $\times 2$.
also nice
physical interpretation

defective eigenspace computations

```
> with(LinearAlgebra):
> A := Matrix(4, 4, [0, 0, 1, 0, 0, 0, 0, 1,
-2, 2, -3, 1, 2, -2, 1, -3]);
```

$$A := \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ -2 & 2 & -3 & 1 \\ 2 & -2 & 1 & -3 \end{bmatrix} \quad (1)$$

```
> Eigenvectors(A);
```

$$\begin{bmatrix} 0 \\ -2 \\ -2 \\ -2 \end{bmatrix}, \begin{bmatrix} 1 & 0 & -\frac{1}{2} & 0 \\ 1 & -\frac{1}{2} & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix} \quad (2)$$

```
> Iden := Matrix(4, shape = identity);
```

$$Iden := \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \quad (3)$$

```
> B := A + 2*Iden; #eigenvalue of A is -2, this is A-λ·I
```

$$B := \begin{bmatrix} 2 & 0 & 1 & 0 \\ 0 & 2 & 0 & 1 \\ -2 & 2 & -1 & 1 \\ 2 & -2 & 1 & -1 \end{bmatrix} \quad (4)$$

```
> NullSpace(B^3); # the nullspace of (A-λ·I)^k is k -- dimensional, where k is the algebraic multiplicity of λ
NullSpace(B^2);
NullSpace(B);
```

$$\left\{ \begin{bmatrix} -\frac{1}{2} \\ 0 \\ 0 \\ 1 \end{bmatrix}, \begin{bmatrix} -\frac{1}{2} \\ 0 \\ 1 \\ 0 \end{bmatrix}, \begin{bmatrix} -1 \\ 1 \\ 0 \\ 0 \end{bmatrix} \right\}$$

$$\left\{ \begin{bmatrix} -\frac{1}{2} \\ 0 \\ 1 \\ 0 \end{bmatrix}, \begin{bmatrix} -\frac{1}{2} \\ 0 \\ 0 \\ 1 \end{bmatrix}, \begin{bmatrix} -1 \\ 1 \\ 0 \\ 0 \end{bmatrix} \right\}$$

$$\left\{ \begin{bmatrix} -\frac{1}{2} \\ 0 \\ 1 \\ 0 \end{bmatrix}, \begin{bmatrix} 0 \\ -\frac{1}{2} \\ 0 \\ 1 \end{bmatrix} \right\} \quad (5)$$

>
From the above we deduce that the chain structure of the generalized $\lambda = -2$ nullspace is one chain of length two and one chain of length 1.

```
> u2 := Vector([0, 0, -1, 1]);
#for chain of length 2 don't start with an eigenvector. I like this one for its physical meaning -
#it's the sum of the first two generalized eigenvectors.
u1 := B.u2;
```

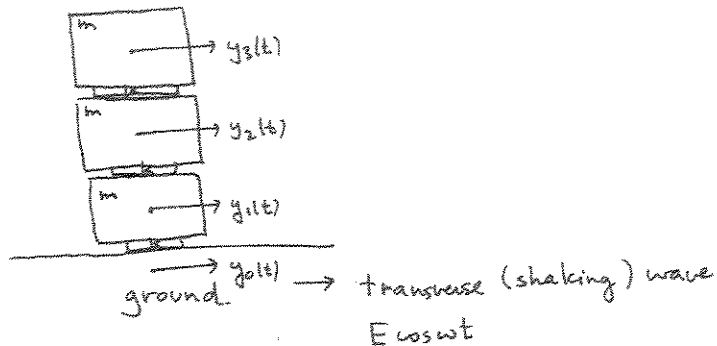
$$u2 := \begin{bmatrix} 0 \\ 0 \\ -1 \\ 1 \end{bmatrix}$$

$$u1 := \begin{bmatrix} -1 \\ 1 \\ 2 \\ -2 \end{bmatrix} \quad (6)$$

Earthquake Project Comments

(A careful explanation of the book's claim that even though the building is only being shaken at the base, in a frame of reference moving with ground, each floor experiences an equal "inertial force".)

consider a 3 story building (in project is 7 stories) modeled as a mass-spring model, with each story having mass m , and floor junctions modeled as "springs" with constants k (with respect to transverse as opposed to parallel motion)



$$y_0(t) = E \cos \omega t \quad \text{so}$$

$$\begin{aligned} y_0'' &= -E\omega^2 \cos \omega t \\ m y_1'' &= -k(y_1 - y_0) + k(y_1 - y_2) \\ m y_2'' &= -k(y_2 - y_1) - k(y_2 - y_3) \\ m y_3'' &= -k(y_3 - y_2) \end{aligned}$$

reduction
of fns

let

$$\begin{aligned} x_0 &= y_0 - E \cos \omega t \equiv 0 \\ x_1 &= y_1 - E \cos \omega t \\ x_2 &= y_2 - E \cos \omega t \\ x_3 &= y_3 - E \cos \omega t \end{aligned}$$

$$\begin{aligned} x_0 &\equiv 0 \\ m x_1'' &= m y_1'' + m E \omega^2 \cos \omega t \\ m x_2'' &= m y_2'' + m E \omega^2 \cos \omega t \\ m x_3'' &= m y_3'' + m E \omega^2 \cos \omega t \end{aligned}$$

so

$$\begin{aligned} m x_1'' &= -k(x_1 - 0) - k(x_1 - x_2) + m E \omega^2 \cos \omega t \\ m x_2'' &= -k(x_2 - x_1) - k(x_2 - x_3) + m E \omega^2 \cos \omega t \\ m x_3'' &= -k(x_3 - x_2) + m E \omega^2 \cos \omega t \end{aligned}$$

so the x 's are the displacements from equilibrium as measured by an observer on the ground

Notice $y_i - y_j = x_i - x_j$

$$\begin{bmatrix} x_1'' \\ x_2'' \\ x_3'' \end{bmatrix} = \begin{bmatrix} -2k/m & k/m & 0 \\ k/m & -2k/m & k/m \\ 0 & k/m & -k/m \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \end{bmatrix} + E \omega^2 \cos \omega t \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}$$

it's as if all floors are being forced - book calls this inertial forcing due to frame of reference