

Math 2210-3
 Wednesday Feb. 3
 9 11.4-11.5

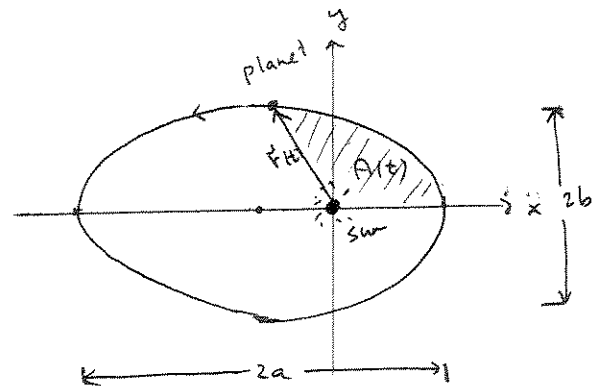
HW for 2/10
 TB/A

①

Kepler's Laws and Newton's amazing deduction

with data from Tycho Brahe,
 Johannes Kepler formulated the
 following empirical laws of planetary
 motion

- ① Planets orbit the sun in ellipses,
 with the sun at a focus of the
 ellipse
- ② A line segment from the focus sun
 to the planet (i.e. $\vec{r}(t)$ if sun is origin)
 sweeps out area at a constant rate
 ("equal areas in equal times")
- ③ The square of an orbit's period is
 proportional to the cube of the ellipse's
 major axis (same proportionality constant
 for all planetary orbits)



↑
 Kepler's Laws

Newton's
 Laws

Newton was able to explain Kepler's Laws with
 and

$$\begin{aligned}\vec{F} &= m \vec{F}''(t) \\ \vec{F} &= \frac{GmM}{r^2} \left(-\frac{\vec{r}}{r} \right)\end{aligned}$$

m = planet mass
 M = sun mass
 G = universal const
 $r = \|\vec{r}\|$

which reduces to

$$\vec{F}''(t) = -\frac{c}{r^2} \left(\frac{\vec{r}}{r} \right)$$

↑
 unit
 vector from
 planet to sun
 (sun feels opposite force,
 but since $M \gg m$ it
 hardly accelerates)

The text explains how $NL \Rightarrow KL$.

But the profound science is that Kepler's
 empirical laws imply that planetary acceleration
 must be towards the sun, and that the strength
 of the acceleration satisfies the inverse square law

In other words, Newton deduced the inverse square law of
 gravitational attraction as a mathematical consequence
 of Kepler's Laws. ~ he never actually directly measured planetary
 to sun gravitational attraction !!

Newton's deductions

(2)

Step 1: (II) Equal areas in equal time with planar orbit)
holds if and only if $\vec{r}''(t)$ is parallel to $\vec{r}$, i.e.
planet accelerates towards sun:

in time Δt ,
the growth in area ΔA is

$$\Delta A \approx \underbrace{\frac{1}{2} \|\vec{r}(t) \times \Delta \vec{r}\|}_{\text{triangle area}}$$

$$\frac{\Delta A}{\Delta t} \approx \frac{1}{2 \Delta t} \|\vec{r}(t) \times \Delta \vec{r}\| = \frac{1}{2} \|\vec{r}(t) \times \frac{\Delta \vec{r}}{\Delta t}\|$$

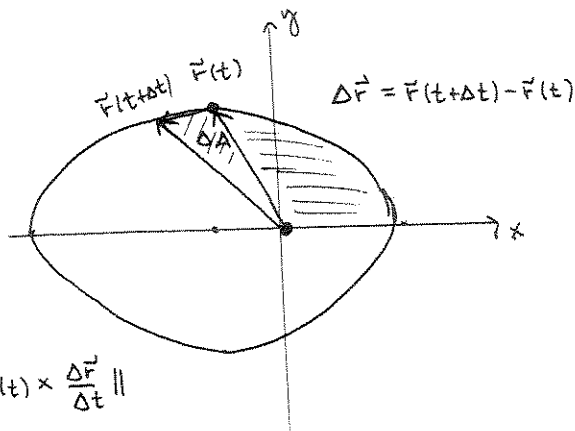
$$\lim_{\Delta t \rightarrow 0} : \text{const} = \frac{dA}{dt} = \frac{1}{2} \|\vec{r} \times \vec{r}'\|$$

(II)

but $\vec{r} \times \vec{r}'$ is a positive multiple
of $\hat{k}$, so in fact

$$\boxed{\vec{r} \times \vec{r}' = \vec{C}} \quad (= \text{const } \hat{k})$$

$$\text{take } \frac{d}{dt} : \underbrace{\vec{r}' \times \vec{r}' + \vec{r} \times \vec{r}''}_{\vec{0}} = \vec{0} \Rightarrow \boxed{\vec{r}'' \parallel \vec{r}} \quad \text{(all steps reversible)}$$



1b) For later, compute $\frac{dA}{dt} = C_1$ in terms of orbital data; i.e. the period T & the ellipse:

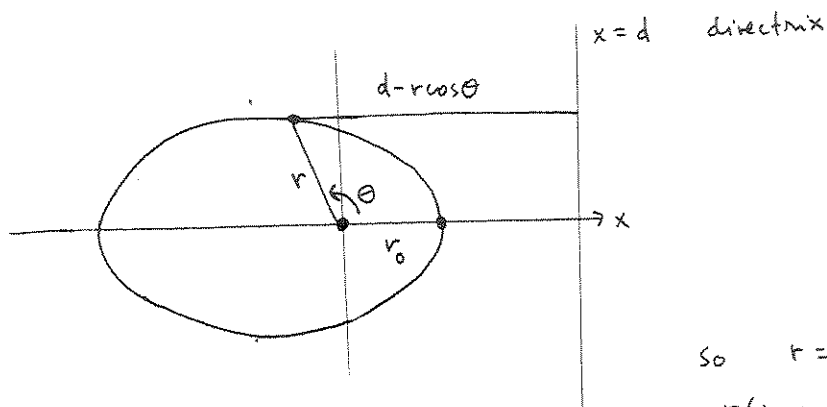
$$\int_0^T \underbrace{\frac{dA}{dt}}_{C_1} dt = \text{ellipse area} = \pi ab$$

$$C_1 T = \pi ab$$

$$\boxed{\frac{dA}{dt} = C_1 = \frac{\pi ab}{T}}$$

$$\boxed{\|\vec{r} \times \vec{r}'\| = 2C_1 = \frac{2\pi ab}{T}}$$

Step 2 polar coord. form of an ellipse (or other conic) with focus at origin



$$\frac{r}{d - r \cos \theta} = e \quad (\text{eccentricity}) \quad \text{const.}$$

so $r = e(d - r \cos \theta)$

$$r(1 + e \cos \theta) = ed$$

$$r = \frac{ed}{1 + e \cos \theta}$$

$e < 1$ ellipse
 $e > 1$ hyperbola

Step 3 rectangular (x-y) eqn for same ellipse ~ details shown for your convenience.

$$\frac{\sqrt{x^2 + y^2}}{d - x} = e$$

$$\sqrt{x^2 + y^2} = e(d - x)$$

$$x^2 + y^2 = e^2(d^2 - 2dx + x^2)$$

$$x^2(1 - e^2) + 2e^2dx + y^2 = e^2d^2$$

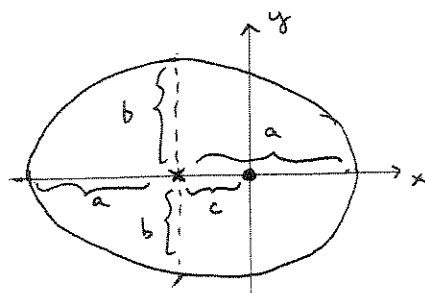
$$(1 - e^2)\left(x^2 + \frac{2e^2}{1 - e^2}dx\right) + y^2 = e^2d^2$$

$$(1 - e^2)\left(x + \frac{de^2}{1 - e^2}\right)^2 + y^2 = e^2d^2 + \frac{d^2e^4}{1 - e^2}$$

$$= \frac{e^2d^2}{1 - e^2}$$

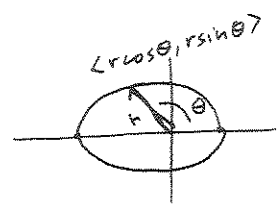
$$\frac{\left(x + \frac{de^2}{1 - e^2}\right)^2}{\left(\frac{e^2d^2}{(1 - e^2)}\right)} + \frac{y^2}{\left(\frac{e^2d^2}{1 - e^2}\right)} = 1$$

$$a = \frac{e^2d^2}{1 - e^2}, \quad b = \frac{ed}{\sqrt{1 - e^2}}, \quad c = \frac{de^2}{1 - e^2}$$



④

Step 4 : $\vec{F}'$, $\vec{F}''$ computations using polar coordinates



$$\begin{aligned} \vec{F}(t) &= r \langle \cos \theta, \sin \theta, 0 \rangle \text{ in } \mathbb{R}^3 \\ r \langle \cos \theta, \sin \theta \rangle &\text{ in } \mathbb{R}^2 \end{aligned} \quad \begin{aligned} r &= r(t) \\ \theta &= \theta(t) \end{aligned}$$

so $\vec{F}'(t) = r' \langle \cos \theta, \sin \theta, 0 \rangle + r \theta' \langle -\sin \theta, \cos \theta, 0 \rangle$

check!

thus $\vec{F} \times \vec{F}' = \vec{F} \times \left(\frac{r'}{r} \vec{F} + r \theta' \langle -\sin \theta, \cos \theta, 0 \rangle \right) = r^2 \theta' \hat{k}$

0

deduce from Step (b)

$r^2 \theta' = 2C_1 = \frac{2\pi ab}{T}$

continue to $\vec{F}''$:

$$\begin{aligned} \vec{F}''(t) &= r'' \langle \cos \theta, \sin \theta, 0 \rangle + r' \theta' \langle -\sin \theta, \cos \theta, 0 \rangle \\ &\quad + r' \theta' \langle -\sin \theta, \cos \theta, 0 \rangle + r \theta'' \langle -\sin \theta, \cos \theta, 0 \rangle \\ &\quad + r (\theta')^2 \langle -\cos \theta, -\sin \theta, 0 \rangle \end{aligned}$$

$$\vec{F}''(t) = (r'' - r(\theta')^2) \underbrace{\langle \cos \theta, \sin \theta \rangle}_{\text{radial unit vector}} + (2r'\theta' + r\theta'') \underbrace{\langle -\sin \theta, \cos \theta \rangle}_{\perp \text{ to radial}}$$

since acceleration is radial (step 1), this is its strength, and this is zero.

So, write $f = r'' - r(\theta')^2$. We (Newton) hope $f = \frac{c}{r^2}$, i.e. $r^2 f$ is constant?

Compute!

$$\begin{aligned} r^2 f &= r^2 r'' - r^3 (\theta')^2 \\ &= r^2 r'' - \frac{1}{r} (r^2 \theta')^2 \\ r^2 f &= r^2 r'' - \frac{1}{r} (2C_1)^2 \end{aligned} \quad \begin{array}{l} \text{from box above} \end{array}$$

almost done!

$$r^2 f = r^2 r'' - \frac{1}{r} (2C_1)^2 \quad \text{from page 4} \quad \text{want this to be const!!}$$



$$r = \frac{ed}{1 + e \cos \theta} \quad \text{step 2}$$

$$r(1 + e \cos \theta) = ed$$

$$r'(1 + e \cos \theta) - re \theta' \sin \theta = 0$$

$$r' \left(\frac{ed}{r} \right) - re \theta' \sin \theta = 0$$

$$r' \cancel{ed} - \underbrace{r^2 \theta'}_{2C_1} \cancel{\sin \theta} = 0$$

$$r' = \frac{2C_1}{d} \sin \theta$$

$$r'' = \frac{2C_1}{d} \theta' \cos \theta$$

So

$$\begin{aligned} r^2 f &= r^2 \left(\frac{2C_1}{d} \theta' \right) \cos \theta - \frac{1}{r} (2C_1)^2 \\ &= (2C_1)^2 \left[\frac{\cos \theta}{d} - \frac{1}{r} \right] \\ &= (2C_1)^2 \left[\frac{\cos \theta}{d} - \frac{1 + e \cos \theta}{ed} \right] \\ &= (2C_1)^2 \left(-\frac{1}{ed} \right) \quad \text{a constant!} \end{aligned}$$

$$= 4\pi^2 \frac{a^2 b^2}{T^2} \left(-\frac{a}{b^2} \right) \quad \swarrow \text{Step 3!!}$$

$$r^2 f = -\frac{a^3}{T^2} 4\pi^2$$

$$= -k 4\pi^2 \quad k \text{ from Kepler's 3rd law III!}$$

$$f = -\frac{4\pi^2 k}{r^2} \quad \text{independent of planet!}$$

$$\boxed{\vec{r}'' = -\frac{4\pi^2 k}{r^2} \frac{\vec{r}}{\|\vec{r}\|}}$$

Kepler's observations $\Rightarrow$ Newton's theory of gravitation!