# **Forced Undamped Oscillations**

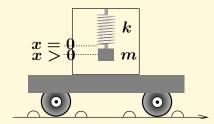
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### **Forced Undamped Motion**

The equation for study is a forced spring-mass system

$$mx''(t) + kx(t) = f(t).$$

The model originates by equating the Newton's second law force mx''(t) to the sum of the Hooke's force -kx(t) and the external force f(t). The physical model is a laboratory box containing an undamped spring-mass system, transported on a truck as in Figure 1, with external force  $f(t) = F_0 \cos \omega t$  induced by the speed bumps.



**Figure** 1. An undamped spring-mass system in a box is transported on a truck. Speed bumps on the shoulder of the road induce periodic vertical oscillations to the box.

# **Undamped Spring-Mass System**

The forced spring-mass equation without damping is

$$x''(t)+\omega_0^2\,x(t)=rac{F_0}{m}\cos\omega t, \quad \omega_0=\sqrt{k/m}.$$

The general solution x(t) always presents itself in two pieces, as the sum of the homogeneous solution  $x_h$  and a particular solution  $x_p$ . For  $\omega \neq \omega_0$ , the general solution is

$$x(t)=x_h(t)+x_p(t), \ x_h(t)=c_1\cos\omega_0t+c_2\sin\omega_0t, \ c_1,c_2 ext{ constants}, \ x_p(t)=A_1\cos\omega t, \quad A_1=rac{F_0/m}{\omega_0^2-\omega^2}.$$

A general statement can be made about the solution decomposition:

The solution is a sum of two harmonic oscillations, one of natural frequency  $\omega_0$  due to the spring and the other of natural frequency  $\omega$  due to the external force  $F_0 \cos \omega t$ .

# Rapidly and slowly varying functions

The superposition x(t) in (1) will exhibit the phenomenon of **beats** for certain choices of  $\omega_0$ ,  $\omega$ , x(0) and x'(0). For example, consider  $x(t) = \cos \omega_0 t - \cos \omega t$ . Use the trigonometric identity  $2\sin a\sin b = \cos(a-b) - \cos(a+b)$  to write  $x(t) = A(t)\sin\frac{1}{2}(\omega_0+\omega)t$  where  $A(t)=2\sin\frac{1}{2}(\omega_0-\omega)t$ . If  $\omega\approx\omega_0$ , then A(t) has natural frequency  $\alpha=\frac{1}{2}(\omega_0-\omega)$  near zero. The natural frequency  $\beta=\frac{1}{2}(\omega_0+\omega)$  can be relatively large and therefore x(t) is a product of a **slowly varying** amplitude  $A(t)=2\sin\alpha t$  and a **rapidly varying** oscillation  $\sin\beta t$ .

The physical phenomenon of **beats** refers to the periodic cancelation of sound at a slow frequency. An illustration of the graphical meaning of *beats* appears in Figure 2.

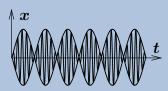
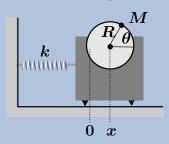


Figure 2. The phenomenon of beats. Shown is a rapidly-varying periodic oscillation  $x(t) = 2\sin 4t \sin 40t$  and the two slowly-varying envelope curves  $x_1(t) = 2\sin 4t, x_2(t) = -2\sin 4t.$ 

### Rotating drum on a cart

Figure 3 shows a model for a rotating machine, like a front-loading clothes dryer. For modeling purposes, the rotating drum with load is replaced by an idealized model: a mass M on a string of radius R rotating with angular speed  $\omega$ . The center of rotation is located along the center-line of the cart. The total mass m of the cart includes the rotating mass M, which we imagine to be an off-center lump of wet laundry inside the dryer drum.



**Figure** 3. A rotating vertical drum installed on a cart with skids.

Vibrations cause the cart to skid left or right. A spring of Hooke's constant k restores the cart to its equilibrium position x = 0. The cart has position x > 0 corresponding to skidding distance x to the right of the equilibrium position, due to the off-center load. Similarly, x < 0 means the cart skidded distance |x| to the left.

The undamped oscillator model is

(2) 
$$mx''(t) + kx(t) = RM\omega^2 \cos \omega t.$$

#### **Model Derivation**

Friction ignored, Newton's second law gives force  $F = m\overline{x}''(t)$ , where  $\overline{x}$  locates the cart's center of mass. Hooke's law gives force F = -kx(t). The centroid  $\overline{x}$  can be expanded in terms of x(t) by using calculus moment of inertia formulas. Let  $m_1 = m - M$  be the cart mass,  $m_2 = M$  the drum mass,  $x_1 = x(t)$  the moment arm for  $m_1$  and  $x_2 = x(t) + R\cos\theta$  the moment arm for  $m_2$ . Then  $\theta = \omega t$  in Figure 3 gives

(3) 
$$\overline{x}(t) = \frac{m_1 x_1 + m_2 x_2}{m_1 + m_2}$$

$$= \frac{(m - M)x(t) + M(x(t) + R\cos\theta)}{m}$$

$$= x(t) + \frac{RM}{m}\cos\omega t.$$

Force competition  $m\overline{x}'' = -kx$  and derivative expansion results in the forced harmonic oscillator

$$mx''(t) + kx(t) = RM\omega^2 \cos \omega t.$$