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Fluctuations of Lévy processes from Wiener-Hopf to the Scattering

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Examples

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Other important example : The **Gamma process** (γ_t) , γ_1 is a standard exponential r.v. and $\phi(\lambda) = \text{Ln}(1 + \lambda)$.

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- The **killing formalism**:
 $\phi(0) + c$ ($c > 0$) is the exponent of a (still called) Lévy process killed at an exponential independent time with rate c .

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Followers (non-exhaustive list) : .L.Alili, J. Bertoin, L. Chaumont, R. Doney, T.Duquesne, A.Kyprianou, A. Kuznetsov, J.C Prado, V. Rivero, V.Vigon, and S.Asmussen, M.Pistorius, E. Eberlein (with financial applications).

Undoubtly a central result for fluctuations of Lévy processes.

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Concisely,

1) Every exponent of a (possibly killed) Lévy process is the **product** of the exponent of a subordinator and of the opposite of a subordinator: there exists two exponents κ and $\hat{\kappa}$ of subordinators such that

$$\phi(iu) = \kappa(iu)\hat{\kappa}(iu)$$

$[iu \in i\mathbb{R}, \text{ at least }]$. This decomposition is unique up to a positive multiplicative constant.

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3) If X is a killed process with death time ζ , or if X_t goes to $-\infty$, for $t \rightarrow +\infty$, then the r.v. $H_\zeta = \sup\{X_s; s \leq \zeta\}$ has Laplace transform $\frac{\kappa(0)}{\kappa(\lambda)}$.

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$$(-iu + h_-(a, c))(iu + h_+(a, c)) = \frac{u^2}{2} + aiu + c$$

$h_-(a, c) \in \mathbb{R}^+$ and $-h_+(a, c) \in \mathbb{R}^+$ are the 2 solutions of the equation : $-\frac{z^2}{2} + az + c = 0$

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- If ϕ has a meromorphic continuation on the right half plan (a fortiori in the whole complex plane), we obtain closed formulas.

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a) f is holomorphic on the two half-planes

$\{\Re(\lambda) > 0\} \cup \{\Re(\lambda) < 0\}$,

$$\lim_{\varepsilon \rightarrow 0^+} \frac{f(iu + \varepsilon)}{f(iu - \varepsilon)} = \kappa(iu)\check{\kappa}(iu) = \phi(iu)$$

b) $(|f(\lambda)| + |\frac{1}{\check{f}(\lambda)}|) \cdot \inf(|\lambda|, \frac{1}{|\lambda|})$ is bounded on \mathbb{C} .

\implies A complete characterization of f (thus of κ and $\check{\kappa}$), up to a multiplicative constant.

Connection with the additive decomposition

Put $\phi(0) = 1$, then $\text{Ln } \phi$ is the exponent of the subordinate process X_{γ_t} , where:

- ▶ (γ_t) a gamma-Process and X is the non killed Lévy process with exponent $\phi(iu) - \phi(0)$
- ▶ (X_{γ_t}) is a Lévy process which does not die, has no drift and is with bounded variations,

$$\text{Ln} \left(\frac{\kappa(iu)}{\kappa(0)} \right) + \text{Ln} \left(\frac{\hat{\kappa}(iu)}{\hat{\kappa}(0)} \right) = \text{Ln } \phi(iu)$$

is the additive decomposition of ϕ .

The second step : the bilateral problem

The problem : Assuming that the Lévy process dies, what is the joint distribution of the maximum value (M) and minimum value (m) ? What is the distribution of the amplitude ($M - m$) ? What is the distribution of the first exit from a bounded interval ?

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Define for $x > 0$, $\lambda \in \mathbb{C}$

$$M_{1,2}^-(x, \lambda) = M_{1,1}^+(x, \lambda) = \mathbf{P}(e^{-\lambda M}; M - m \leq x)$$

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then

$$\frac{d}{dx} \begin{pmatrix} M_{1,1}^+ \\ M_{2,1}^+ \end{pmatrix} = \begin{pmatrix} e^{-\lambda x} & v(x) \\ \hat{v}(x) & e^{\lambda x} \end{pmatrix} \begin{pmatrix} M_{1,1}^+ \\ M_{2,1}^+ \end{pmatrix} \quad (1)$$

$(v(x), \hat{v}(x))$ is the **potential**.

Auxilliary solutions

Let , for $\Re(\lambda) > 0$,

$$\begin{pmatrix} M_{12}^+(x, \lambda) \\ M_{22}^+(x, \lambda) \end{pmatrix}$$

and for $\Re(\lambda) < 0$,

$$\begin{pmatrix} M_{11}^-(x, \lambda) \\ M_{21}^-(x, \lambda) \end{pmatrix}$$

be solutions of the equation (1), holomorphic in λ , on their own half-planes.

Finally, the 2x2 matrices :

$$M^+(x, \lambda) \quad \text{for} \quad \Re(\lambda) > 0$$

$$M^-(x, \lambda) \quad \text{for} \quad \Re(\lambda) < 0$$

satisfy the differential equation :

$$\frac{d}{dx} M^{+/-}(x, \lambda) = \begin{pmatrix} e^{-\lambda x} & v(x) \\ \hat{v}(x) & e^{\lambda x} \end{pmatrix} M^{+/-}(x, \lambda)$$

To switch from the exponent ϕ to the solution of the bilateral problem is again a Riemann-Hilbert problem.

Theorem

For all $x > 0$, the matrices $M^+(x, \cdot)$ and $M^-(x, \cdot)$ satisfy for all $iu \in i\mathbb{R}$,

$$[M^-(x, iu)]^{-1} \cdot M^+(x, iu) = \begin{pmatrix} 0 & -e^{-iu x} \\ e^{iu x} & \phi(iu) \end{pmatrix}$$

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and the additional boundary properties :

a) $\lambda \rightarrow M^+(x, \lambda)$ (resp. $\lambda \rightarrow M^-(x, \lambda)$) is holomorphic on $\{\Re(\lambda) > 0\}$ (resp. on $\{\Re(\lambda) < 0\}$), continuous on the closed half-plane $\{\Re(\lambda) \geq 0\}$ (resp. on $\{\Re(\lambda) \leq 0\}$).

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b) $M^+(x, \lambda) \sim M^+(y, \lambda)$ for $\lambda \rightarrow +\infty$, $M^-(x, \lambda) \sim M^-(y, \lambda)$ for $\lambda \rightarrow -\infty$ and $\det M^-(x, \lambda) = \det M^+(x, \lambda) = 1$.

Starting from the **Scattering matrix**

$$\begin{pmatrix} 0 & -1 \\ 1 & \phi(iu) \end{pmatrix},$$

the preceding properties entirely characterize the two holomorphic functions $\lambda \rightarrow M^+(x, \lambda)$ and $\lambda \rightarrow M^-(x, \lambda)$ for all $x \in]0, +\infty[$ (up to a constant of x and λ) and the potential $(v(x), \hat{v}(x))$.

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Explicit Solutions

- Levitan et Marchenko in the 50's or 60's have solved the case that corresponds to the Brownian motion.
- Explicit solutions can be obtained when ϕ is a meromorphic function, on one half-plane and a fortiori on the two half-planes, and this corresponds to the so called "Bargmann equations" in Scattering Theory.

Some words on Scattering Theory

Scattering Theory : To a **potential** [the two functions $(v(x), \hat{v}(x))$], we associate two parametrized differential equations :

$$M'(x, \lambda) = \begin{pmatrix} e^{-\lambda x} & v(x) \\ \hat{v}(x) & e^{\lambda x} \end{pmatrix} M(x, \lambda)$$

The wronskian identity gives

$$[M^-(x, iu)]^{-1} M^+(x, iu) = \begin{pmatrix} 0 & -e^{-iu x} \\ e^{iu x} & \phi(iu) \end{pmatrix}$$

The matrix $\Phi = \begin{pmatrix} 0 & -1 \\ 1 & \phi(iu) \end{pmatrix}$ is called the Scattering Matrix.

The mapping $(v(x), \hat{v}(x)) \rightarrow \Phi$ is injective. To start from ϕ in order to determine the potential is called the **inverse scattering problem**.

Thus computing the distributions related to the bilateral problem when starting from the Lévy exponent is part of that physics problem.

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