

THE FINITE ELEMENT METHOD

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This is the plan for this part of the class

§ 11.5 Rayleigh-Ritz method (FEM for BVP) } Brenner & Scott

§ 12.4 The FEM for elliptic problems

§ 12.1 - 12.3 Elliptic, parabolic and hyperbolic problems } Barden & Fairweather
Kincaid & Cheney

What is the finite element method? It is a systematic way of generating discrete algorithms for approximating the solution of differential equations.

We start first with an overview of FEM by looking at BVP:

$$(BVP) \begin{cases} -\frac{d^2 u}{dx^2} = f & \text{in } (0,1) \\ u(0) = 0 \\ u'(1) = 0 \end{cases}$$

If u is the solution and v is any suff. smooth function with $v(0)=0$:

$$(f, v) := \int_0^1 f(x) v(x) dx = - \int_0^1 u'' v dx = \underbrace{-u' v \Big|_0^1}_{=0} + \underbrace{\int_0^1 u' v' dx}_{:= a(u, v)}$$

Thus we reformulate (BVP) as follows:

Let $V = \{ v \in L^2(0,1) \mid a(v, v) < \infty \text{ and } v(0)=0 \}$

If u solves (BVP) then $u \in V$ and

$$\boxed{\forall v \in V \quad a(u, v) = (f, v)} \quad (WF)$$

This is the "variational" or "weak" formulation of BVP.

Can we show that if u satisfies (WF) then u solves (BVP)?

This holds if we assume additional regularity of f and u :

Theorem: let $f \in C^0([0,1])$ and $u \in C^2([0,1])$ satisfy (WF). (45)

Then u solves (BVP).

Proof: let $v \in V \cap C^1([0,1])$.

$$(*) \quad (f, v) = a(u, v) = \int_0^1 u' v' dx \stackrel{\text{IBP}}{=} \underbrace{u' v \Big|_0^1}_{= u'(1)v(1)} - \int_0^1 u'' v dx$$

Choosing $v \in V$ s.t. $v(1) = 0$:

$$(f, v) = - \int_0^1 u'' v dx$$

$$(\Rightarrow) \quad \left(\underbrace{f + u''}_{w \in C^0}, v \right) = 0.$$

Assume for contradiction that $w \neq 0$. Then w must assume one sign on some interval $[a, b] \subset [0, 1]$ (by continuity). Choose

$$v(x) = \begin{cases} (x-a)^2(x-b)^2 & \text{in } [a, b] \\ 0 & \text{otherwise} \end{cases}$$

(or any other v with one sign for that matter)

then $(w, v) \neq 0$ which is a contradiction. $\Rightarrow -u'' = f$.

Now for B.C.:

use $v(x) = x$ in $(*)$ to get:

$$u'(1) = (f + u'', v) = 0$$

The other B.C. comes from $u \in V$:

$$u(0) = 0$$

QED.

The following nomenclature is given to boundary conditions.

$u(x) = 0 \Rightarrow$ Dirichlet B.C. $\equiv$ essential B.C. because it appears explicitly in definition of V .

$u'(x) = 0 \Rightarrow$ Neumann B.C. $\equiv$ natural B.C. because it appears implicitly in variational formulation

In light of previous theorem it is clear that the weak formulation is a way of interpreting (BVP) with more loose assumption on f . (76)

Ritz Galerkin Approximation

Let $S \subset V$ be any finite dimensional subspace of V .

Replacing V by S in (WF) we get:

$$\boxed{\text{Find } u_S \in S \text{ such that } a(u_S, v) = (f, v) \quad \forall v \in S} \quad (RG)$$

We will see that by this simple device we have created a linear system from BVP!

Theorem Given $f \in L^2(\Omega, 1)$, (RG) has a unique solution.

Proof: Let $\{\phi_i\}_{i=1}^n$ be a basis of S . We rewrite (RG) in terms of this basis:

$$\text{Let } u_S = \sum_{i=1}^n U_i \phi_i \text{ then:}$$

$$a(u_S, v) = \sum_{i=1}^n U_i a(\phi_i, v) = (f, v) \quad \forall v \in S$$

$$\Leftrightarrow a(u_S, \phi_j) = \sum_{i=1}^n U_i a(\phi_i, \phi_j) = (f, \phi_j) \quad j=1 \dots n$$

$$\Leftrightarrow KU = F$$

where $K \in \mathbb{R}^{n \times n}$ and $K_{ij} = a(\phi_j, \phi_i)$

$$F \in \mathbb{R}^n, \quad F_i = (f, \phi_i)$$

$$U \in \mathbb{R}^n \quad U = \begin{pmatrix} U_1 \\ \vdots \\ U_n \end{pmatrix}$$

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This linear system has a unique solution: if it did not then $\exists V \neq 0$ s.t. $KV = 0$

Let $v = \sum_{j=1}^n V_j \phi_j$ then

$$a(v, \phi_j) = 0 \quad j=1 \dots n$$

$$\Rightarrow \sum_{j=1}^n V_j a(v, \phi_j) = 0$$

$$\Rightarrow a(v, v) = 0 = \int_0^1 (v')^2 dx$$

$$\Rightarrow v' = 0 \Rightarrow v = \text{constant} = v(0) = 0.$$

$$\Rightarrow V = 0. \quad \underline{QED.}$$

note: some technical points have been swept under the rug here.

$K \equiv$ stiffness matrix

$K = K^T$ because the "energy inner prod" $a(\cdot, \cdot)$ is symm (symm)

K is also positive definite since

$$V^T K V = \sum_{i,j=1}^n k_{ij} v_i v_j = a(v, v) \geq 0 \quad \text{for } v = \sum_{j=1}^n v_j \phi_j$$

$$\text{and } V^T K V = 0 \Rightarrow v = 0 \Rightarrow V = 0.$$

Error Estimate (is the approx any good?)

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Galerkin orthogonality:

$$a(u, w) = (f, w) \quad \forall w \in S$$

$$a(u_S, w) = (f, w)$$

$$\Rightarrow \boxed{a(u - u_S, w) = 0 \quad \forall w \in S}$$

Now introduce the so called energy norm:

$$\|u\|_E = (a(u, u))^{\frac{1}{2}}$$

which satisfies Schwartz inequality:

$$|a(v, w)| \leq \|v\|_E \|w\|_E \quad \forall v, w \in V$$

Then $\forall v \in S$:

$$\|u - u_S\|_E^2 = a(u - u_S, u - u_S)$$

$\rightarrow$ by Galerkin $\perp$

$$= a(u - u_S, u - v) + \overbrace{a(u - u_S, v - u_S)}$$

$$\leq \|u - u_S\|_E \|u - v\|_E$$

If $\|u - u_S\|_E \neq 0$ then

$$\|u - u_S\|_E \leq \|u - v\|_E \quad \forall v \in S$$

$$\Rightarrow \|u - u_S\|_E \leq \inf_{v \in S} \|u - v\|_E$$

Since $u_S \in S$:

$$\inf_{v \in S} \|u - v\|_E \leq \|u - u_S\|_E$$

Therefore:

$$\|u - u_S\|_E = \inf_{v \in S} \|u - v\|_E$$

Moreover there is an element (namely u_S) in S for which the infimum is attained. Therefore we get: (79)

$$\text{Theorem} \quad \|u - u_S\|_E = \min_{v \in S} \|u - v\|_E$$

What this says is that the RG approx gives automatically the best approx. to u in S , as measured by the energy norm.

Now let's see what happens with the L^2 norm:

$$\|v\|_2 = (v, v)^{\frac{1}{2}} = \left(\int_0^1 v^2 dx \right)^{\frac{1}{2}}$$

To estimate $\|u - u_S\|$ we use a duality argument (Aubin-Nitsche trick)

Let w be the solution of:

$$\begin{cases} -w'' = u - u_S & \text{on } [0, 1] \\ w(0) = w'(1) = 0 \end{cases}$$

Then:

$$\begin{aligned} \|u - u_S\|^2 &= (u - u_S, u - u_S) \\ &= (u - u_S, -w'') \\ &\stackrel{\text{IBP}}{=} \alpha(u - u_S, w) - \cancel{(u - u_S)w'|_0^1} \\ &= \alpha(u - u_S, w - v) \quad \forall v \in S \\ &\stackrel{\text{Galerkin}}{\perp} \end{aligned}$$

Thus using Schwarz inequality

$$\begin{aligned} \|u - u_S\|^2 &\leq \|u - u_S\|_E \|w - v\|_E \|w''\| \\ \|u - u_S\| &\leq \|u - u_S\|_E \|w - v\|_E / \sqrt{\|u - u_S\|} \end{aligned}$$

Taking the infimum over $v \in S$ we get:

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$$\|u - u_S\| \leq \|u - u_S\|_{\mathcal{E}} \inf_{v \in S} \frac{\|w - v\|_{\mathcal{E}}}{\|w\|}$$

Thus the L^2 error can be made much smaller than the energy norm error provided we make the assumption that for some $\varepsilon > 0$:

$$\inf_{v \in S} \|w - v\|_{\mathcal{E}} \leq \varepsilon \|w\|$$

This is an "approximation assumption" and we get:

$$\|u - u_S\| \leq \varepsilon \|u - u_S\|_{\mathcal{E}} \leq \varepsilon \left(\varepsilon \|u\| \right) = \varepsilon^2 \|u\| = \varepsilon^2 \|f\|.$$

↑
approx. assum.
again with w
replaced by u

To summarize:

$$\|u - u_S\|_{\mathcal{E}} \leq \varepsilon \|f\|$$

$$\|u - u_S\| \leq \varepsilon^2 \|f\| \rightarrow \text{smaller error}$$

We shall construct subspace (finite elements) where ε can be made arbitrarily small.