

### § 3.4 D'Alembert's method (purely 1D)

Recall IDWTEQ

$$\begin{cases} u_{tt} = c^2 u_{xx} & , 0 < x < L, t > 0 \\ u(0,t) = u(L,t) = 0 & , t > 0 \\ u(x,0) = f(x) & , 0 < x < L \\ u_t(x,0) = g(x) & , 0 < x < L \end{cases}$$

D'Alembert's solution:

$$u(x,t) = \frac{1}{2} [f^*(x-ct) + f^*(x+ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} g^*(s) ds \quad (*)$$

where  $f^* \equiv 2L$  per odd ext. of  $f$

$g^* \equiv \underline{\hspace{10em}} g$

But we drop  $*$  notation below to keep notation simple.

Consistency check:

Can  $u_n(x,t) = \sin \frac{n\pi x}{L} \cos \frac{cn\pi t}{L}$  be written as  $(*)$ ?

$$\left. \begin{aligned} \sin(a+b) &= \sin a \cos b + \cos a \sin b \\ \sin(a-b) &= \sin a \cos b - \cos a \sin b \end{aligned} \right\} \Rightarrow \sin a \cos b = \frac{1}{2} [\sin(a+b) + \sin(a-b)]$$

$$\Rightarrow u_n(x,t) = \frac{1}{2} \left[ \sin \frac{n\pi}{L} (x+ct) + \sin \frac{n\pi}{L} (x-ct) \right]$$

Does  $(*)$  solve IDWTEQ?

To compute derivatives of  $(*)$  we need the following formula from calc III:

$$\frac{d}{dt} \left[ \int_{a(t)}^{b(t)} f(s;t) ds \right] = b'(t) f(b(t), t) - a'(t) f(a(t), t) + \int_{a(t)}^{b(t)} \frac{\partial f}{\partial t}(s,t) ds$$

Hence:

$$u_t = \frac{c}{2} [-f'(x-ct) + f'(x+ct)] + \frac{1}{2} [g(x+ct) + g(x-ct)]$$

$$u_{tt} = \frac{c^2}{2} [f''(x-ct) + f''(x+ct)] + \frac{c}{2} [g'(x+ct) - g'(x-ct)]$$

$$u_x = \frac{1}{2} [f'(x-ct) + f'(x+ct)] + \frac{1}{2c} [g(x+ct) - g(x-ct)]$$

$$u_{xx} = \frac{1}{2} [f''(x-ct) + f''(x+ct)] + \frac{1}{2c} [g'(x+ct) - g'(x-ct)]$$

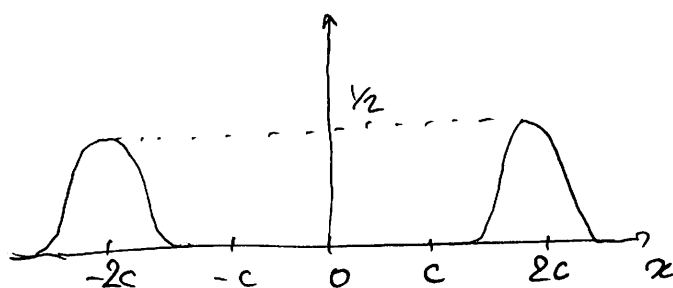
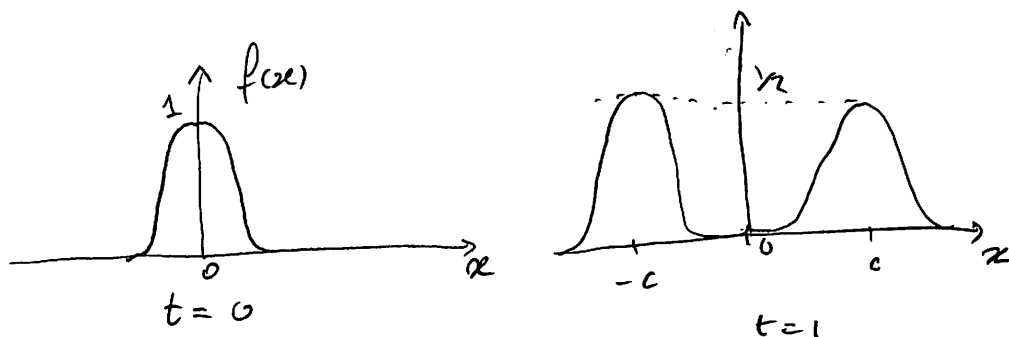
therefore we do have  $u_{tt} = c^2 u_{xx}$ .

Are we done? No we need to check B.C.

$$u(0,t) = \frac{1}{2} [f^*(-ct) + \underbrace{f^*(ct)}_{\text{odd ext.}}] + \frac{1}{2c} \int_{-ct}^{ct} \underbrace{g^*(s)}_{\text{odd}} ds = 0$$

Physical interpretation

$$\text{Assume } g(x) = 0 \Rightarrow u(x,t) = \frac{1}{2} [f^*(x-ct) + f^*(x+ct)]$$



What happens in the general case (i.e.  $g(x) \neq 0$ )

let

$$G(x) = \int_a^x g^*(z) dz$$

$g^*$  is 2L-per

$g^*$  is odd

$$G(x+2L) - G(x) = \int_x^{x+2L} g^*(z) dz \stackrel{\downarrow}{=} \int_{-L}^L g^*(z) dz = 0$$

$\Rightarrow G$  is 2L-periodic

$$\begin{aligned} \Rightarrow u(x,t) &= \frac{1}{2} [f^*(x-ct) + f^*(x+ct)] + \frac{1}{2c} [G(x+ct) - G(x-ct)] \\ &= \underbrace{\frac{1}{2} [f^*(x-ct) - \frac{1}{c} G(x-ct)]}_{\text{right propa. term}} + \underbrace{\frac{1}{2} [f^*(x+ct) + \frac{1}{c} G(x+ct)]}_{\text{left propa term}} \end{aligned}$$

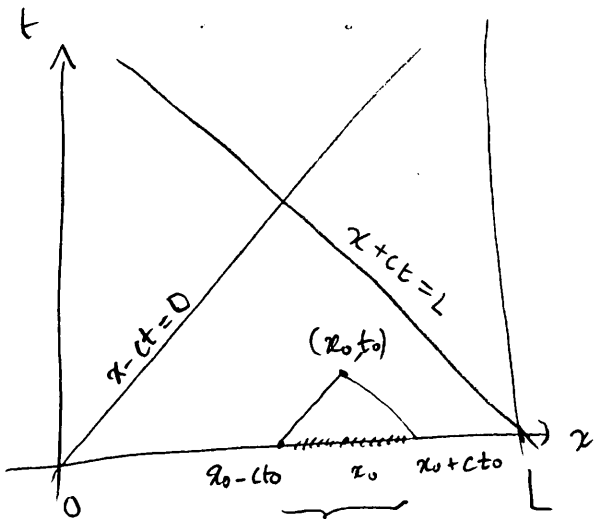
Characteristics

Characteristics = curves where sol is const.

$$x - ct = x_0 - ct_0$$

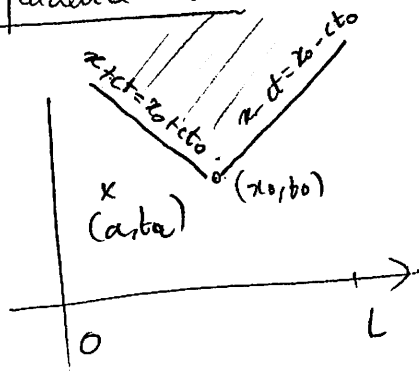
$$x + ct = x_0 + ct_0$$

Characteristics have slope  $\pm \frac{1}{c}$  in  $(x,t)$  plane



interval of dependence: value of  $u(x_0, t_0)$  can only be influenced by sol at this interval

Dependence cone:



because waves propagate at a finite speed, only observers in cone in  $xt$  plane can see disturbance at  $(x_0, t_0)$ .

Observer at  $(a_1, t_1)$  will not see disturbance!