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Amazing Structures of Optimal Multicomponent Composites

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Optimal structures

- Structures, that minimize the sum of energies in several homogeneous and orthogonal external fields, are called "optimal".
 - ✓ P1: Minimal energy in a fixed current minimal total resistivity,
 - ✓ P2: Minimal energy in a fixed electrical field minimal conductivity.

Optimal structures = minimizing sequences in the quasiconvex envelope problem.

$$J = \inf_{\chi} W; \quad W = \inf_{v} \iint_{\Omega} \left[v^{T} K_{i} v + \gamma_{i} \right] \chi_{i} dx, \quad v \text{ is } \Omega \text{ - periodic}, \quad \iint_{\Omega} v dx = v_{0}$$
$$\sum_{j,k} A_{ijk} \frac{\partial v_{j}}{\partial x_{k}} = 0, i = 1, ...p \quad (\nabla \cdot v = 0, \nabla \times v = 0)$$
$$\chi_{i} = \begin{cases} 1 \text{ if } x \in \Omega_{i} \\ 0 \text{ if } x \notin \Omega_{i} \end{cases}, \quad \Omega = \bigcup \Omega_{i}.$$

Instead of the cost γ_i we fix volume fractions $m_i = \int_{\Omega} \chi_i$

Motivations

- Technological capabilities, such as 3d printing and microfabrication, allow a huge variety of structures to be manufactured for roughly the same price, and one wants to know ``the best'' structure.
- Mathematically, the best composite structure is a solution to the quasiconvex envelope problem for a multiwell Lagrangian. Few examples for two-well Lagrangians are known (starting with Lurie &Ch 82, 84, see also Tartar 85) but the technique for multiwell quasiconvex envelopes is in its infancy.
- Aesthetically, the optimal structures are delightful, and intellectually, they present an elegant riddle that links geometry and PDE



Comment on minimizers

- Contrary to homogenization, the problem of optimal composite asks for a partition (may be a fractal-type one) of the domain, that corresponds to extremal average properties.
- Materials (wells) are identified by the range fields (gradients).

Fields in materials are ordered and bounded







Features

- **Simple fields complex structures** In optimal structures, the fields v(x) or ۲ some their components are constant in Omega_i. To provide the constancy, the structures typically become fractals.
- The optimality is established by **geometrically independent bounds** and ٠ corresponding sufficient optimality conditions on v(x). The bounds account for corollaries of differential properties of v(x). They might be exact or rough. Exact bounds are realized by optimal structures (or a sequence of structures)
- The geometry of two-material optimal composites is steady and intuitive, ۲



but geometry of three-material optimal composites is complex, typically non–unique, and cannot be simply guessed.

First optimal three-phase structures Require sufficient large volume fractions of the "best" material k 1. k1-k3 coated circles mimic k2 and become "invisible" mimics Multicoated spheres Lurie, Ch 86 Parallel matrix laminates Milton Kohn, 88



Sigmund, Gibiansky, 2000.

Parallel coated spheres Milton 82

$$\mathbf{k}_* - k_2 \mathbf{I} \le 0$$

Attainability condition

Beyond translation method

- Translation method accounts for differential constraints through an analog of Green's theorem and provides translation bounds.
 - Translation bound is exact for two-well problem and for the case when the volume fraction m1 of the "best material" is larger than a threshold.
 - Translation bound is definitely loose when m1 -> 0.
- One assumes that the fields in an optimal geometry are subject to some additional inequalities that become exact for small m1.
- Question: how to reveal these inequalities?

How the inequalities were revealed for 2d problem

- Gibiansky and Sigmund (2002) suggested a new class of isotropic optimal structures
- Albin, Ch, Nesi investigated the fields in these and similar structures and extended results to anisotropic composites.
- Ch (2009) investigated fields in materials in limiting structures (minimal values of m1). conjectured and proved new inequalities for optimal fields; in the asymptotical case, they coincide with previously found inequalities by Alessandrino and Nesi.
- When these inequalities are added to the Translation method procedure, the novel bounds emerge, that are achievable everywhere.
- For anisotropic bounds, additional inequalities were conjectured by Ch. and G.Dzierzanowski; the proofs are no be perfected.

Meaning of the inequalities

$$A: \quad \det(e_1 - e_3) \ge 0$$

follows from compatibility of the supporting points of the envelope:

Any supporting point is in a rank-one connection with a point in the convex envelope of other supporting points



$$B: \quad \det(e_1 - e_0) \le 0$$

is an optimization requirement (work in progress). The optimality test disproves the field in the forbidden interval

Predicted location of optimal fields in the "best" material



Eigenvalues of supports of quasiconvex envelope (fields in optimal materials), required by the bounds and matching optimal structures



Example: Isotropic Bounds, k3=infinity

Th.

(a) Effective conductivity k_{*} of every isotropic three-phase composite satisfies the bound (Special case of Nesi (1995) and Ch. (2009) bounds)



(b) The bound is exact everywhere.

Optimal structures

- 1. Repeated laminates: Use laminate formulae repeatedly in different scales.
- Coated-spheres-type wheels assembly (effective field theory)
 The assembly elements, "wheels", vary depending on the fractions of materials.





Assembly elements: the optimal topology depends on m₁











m₁>m₁₁ Hashin-Shtrikman bound. K₁ is connected

m₁₂<m₁<m₁₁ No connected phases

m₁<m₁₂ k₂ (gray) is connected

3D problem: search for optimal structures

Problem: so far, the only known bound for 3d is the translation (Hashin– Shtrikman) bound.

The optimal structures should hint for new inequalities and new bounds. We want to extend the domain of applicability for the known bounds, then investigate the relations on the fields in materials in the limiting structures, conjunct and (hopefully) prove new inequalities and derive the bounds.

Difficulties: Laminates need a large number of iterations, choice of substructures, and other not obvious decisions.

P1. Optimal structures of minimal resistance

Given three materials with conductivities k1, k2, and k3: k1>k2 > k3=0 (k3 is a perfect insulator) Find a range of isotropic structures of maximum effective conductivity.

$$v_0^T k_0 v_0 + \gamma_0 = \inf \int_{\Omega} F(v) dx, \quad \nabla \times v = 0,$$

$$F = \min \left(v^T k_1 v + \gamma_1, v^T k_2 v + \gamma_2, v^T k_3 v \right)$$

Known structures: Range of realizability of H-S structures by Coated spheres structures:

$$= \qquad m_1 \ge m_1^{cr} = \frac{3k_2}{k_2 + 2k_1}(1 - m_2)$$





Columnar structure: Condition of piece-wise constancy of fields in bulk structures

After Zhikov, Sigmund.. – separation of variables

$$k_1 e_1 = k_3 e_3, \quad k_2 e_2 = k_4 e_4$$

if $\frac{k_1}{k_3} = \frac{k_2}{k_4}$ then e_i are constant $i = 1..4$

Effective properties:

If
$$\frac{k_1}{k_3} = \frac{k_2}{k_4}$$
 then
 $k_0 = S_A \frac{k_1 k_3}{\gamma k_3 + (1 - \gamma) k_1} + S_B \frac{k_2 k_4}{\gamma k_4 + (1 - \gamma) k_2}$



Bulk blocks assemblage - 1



Inspired by Sigmund-Gibiansky structures,



studied by N.Albin for his PhD.



Parameters

• The volume fractions *cL* and *c* of *k1* in the laminates and cylinders, respectivally, are connected to satisfy the columnar condition. $c_L = \frac{k_1 c^2}{k_2 (2-c)}$,

• The volume fractions and effective conductivity are expressed through structural parameters parameters:

$$\begin{split} m_2 &= \gamma^3, \\ m_1 &= 3\gamma(1-\gamma)c\,\frac{ck_1(1-\gamma)+\gamma k_2(2-c)}{(2-c)k_2}, \\ k_0 &= \gamma ck_1\frac{2ck_1(1-\gamma)+\gamma k_2(2-c)}{(2-c)(c\gamma k_1+(1-\gamma)k_2)} \end{split}$$

Optimality



Larger *m1*: "Enveloping" theorem

Th: If a structure S is translation–optimal, than a laminate form S and k1 is also translation–optimal.

$$\frac{1}{k_0 + (d-1)k_1} \ge \sum_{i=1}^N \frac{m_i}{k_i + (d-1)k_1} - \text{H-S bound}$$
$$\operatorname{Tr}(K_0 - k_1 I)^{-1} \ge \sum_{i=2}^N \frac{m_i d}{k_i - k_1} + \frac{m_1}{k_1}, \ K_0 \le \langle \mathbf{k} \rangle \text{I-Translation bound}$$

Therefore, it is sufficient to look for a *critical optimal* structure with minimal amount of k1.



k1

Range of optimality



Additional relation on fields (dual variables) in the limiting structures

$$\det(\mathbf{e}_1 - \mathbf{e}_3) = 0 \ \forall x \in \Omega_1$$

Conjectured inequality:

$$\det(\mathbf{e}_1 - \mathbf{e}_3) \ge 0 \ \forall x \in \Omega_1$$

Anisotropic optimal structures (similar to 2d consideration by Ch, Dzierzanowski, 2014)



- The consideration is similar but more messy.
- We show that this anisotropic structure satisfies translation bound at the same value of m1 for all values of gamma_i, if

 $m_2 = \gamma_1 \gamma_2 \gamma_3,$ $m_1 = g(\gamma_1, \gamma_2, \gamma_3)$

 Optimality of degenerative anisotropic structures is conjuncted, but not proven

P2. Dual problem of minimal resistance. Controllable Differential scheme

• Problem: Given three materials with resistances k1, k2, k3

$$k_1 < k_2 < k_3 = \infty$$

find a structure of minimal conductance (max resistance)

$$v_0^T k_0 v_0 + \gamma_0 = \inf \int_{\Omega} F(v) dx, \quad \nabla \cdot v = 0,$$

$$F = \min \left(v^T k_1 v + \gamma_1, v^T k_2 v + \gamma_2, v^T k_3 v \right)$$

The bulk block structures are not optimal.

Controllable differential scheme: Differential equation of structural evolution

 Three infinitesimal layers of volume *dm* of transversally isotropic composite is added in three orthogonal directions to an isotropic material "seed" with conductivity *k(m)*:



Added composite of variable conductivity



Assume that the added layer is a 2d-optimal composite with perfectly conducting cylindrical inclusions

$$k_a = \infty, \quad k_b = k_1 \left(\frac{2}{c} - 1\right).$$

- The potentials at the upper and lower surface are equal, the current flows only in the plane.
- Notice, that the element can "turn" the current and transmit the potential

$$\left|J_{N}^{+}-J_{N}^{-}\right|=\left|J_{plane}\right|, \quad u^{+}=u^{-}$$

c=c(m) – the volume fraction of k1 in the composition – Is a control (a variable parameter, that is adjusted to minimize the conductivity of the whole structure

Optimization problem

Choose c=c(m) to minimize k(1) if 1. $m\frac{dk}{dm} = \frac{2k_b - k}{3}$, $k_b = k_1\left(\frac{2}{c} - 1\right)$ - eq of struct. evolution 2. $k(m_2) = k_2$ - nucleous from k_2 3. $\int c dm = m_1, c(m) \in (0,1].$ Technique: $H = \lambda(m) \frac{2k_1\left(\frac{2}{c}-1\right)-k}{3m} + \beta c$ - Hamiltonian $\frac{d\lambda}{dm} = - \frac{\partial H}{\partial k}, \, \lambda(1) = 1,$ Optimality $\frac{\partial H}{\partial c} = 0 \Longrightarrow c = c_{opt}(\beta) \Longrightarrow \int_{0}^{1} c_{opt} dm = m_1 \Longrightarrow \beta = \beta_0.$ conditions

Optimization problem

Choose *c*=*c*(*m*) to minimize *k*(1)

Solution:

$$c_{opt}(m) = \frac{2m_1}{1 - m_2^{\frac{2}{3}}} m^{-\frac{1}{3}}$$

$$k_{opt} = m_2^{\frac{1}{3}} k_1 - 2k_1 \left(1 - m_2^{\frac{1}{3}}\right) = \frac{2k_1}{m_1} (m_2^{\frac{2}{3}} - 1)^2$$

$$k_{opt} \text{ satisfies H-S bound, when}$$

$$m_1 = m_1^{cr} = m_2^{\frac{1}{3}} \left(1 - m_2^{\frac{1}{3}}\right) \frac{3k_1}{k_2 + 2k_1}$$

Range of optimality



Spherical layer and fields.



At the perforated layer:

u = const

$$k_1 \frac{\partial u}{\partial r}\Big|_{r=r_1} - k_2 \frac{\partial u}{\partial r}\Big|_{r=r_2} = G(u)$$

- Using spherical coordinates and the hint from diff.
 scheme, we also can
 separate variables and
 compute properties of
 coated anisotropic spheres
 geometry, using effective
 field theory. The result is
 the same.
- Larger m1- add new outer layer from m1.

Next evolution: Conjuncted structures

- In the limiting structure, the current in k1 is two–dimensional, the normal current is zero.
- Therefore, the next evolution of optimal structures, when m1 decreases, will likely follow the evolution of 2d optimal structures:







k2 enters the enveloping composite, moving from the central sphere. k1 moves into radial layers

k2 also moves into radial layers, replaces k1

One can check, that the limiting k2-k3 structure is optimal and equivalent to coated spheres Metamaterials - structures with exotic properties

- Bi-products of optimal structures study:
- Using differential scheme and separation of variables in curved coordinates, one can create various geometries with unusual overall properties, inner degrees of freedom, abnormal and nonlinear etc.
- **Example:** A homogeneous external field applied to the spiral assemblies (Ch, Pruss, 13) with inclusions is transformed into a rotated homogeneous field inside inclusions.

Sun rises from the North





Optimal design problem

(in collaboration with G. Dzierzanowski and N. Briggs)

- **Optimal design problem:** determine optimal structures and distribute them in the design domain, keeping the total cost.
- Three material design: Given two materials the strong expensive one, the cheap weak one, and the void, find their layout in a fixed loaded from the boundary domain, that minimizes the stress energy, i.e. maximizes the stiffness.

$$\begin{split} &\inf_{\tau} \left(\int F(\tau, g) dx + \oint f(u^{T} \tau n), \right) \quad \tau : \ \tau = \tau^{T}, \nabla \cdot \tau = 0, \\ &F = \min(W_{i}(\tau) + g_{i}), \quad W_{i} = k_{i} Tr(\tau^{2})/2 \\ &k_{1} < k_{2} < k_{3} = \infty, \quad g_{1} = 1, \quad g_{3} = 0, \quad g_{2} = \gamma, \end{split}$$

Multimaterial design



Multimaterial designs are to two-material ones as color TV is to black–and–white TV.

Multimaterial design

Color TV

Two-material design

Black-and-white TV



Optimal structures and Quasiconvex envelope



Optimal structures depend on intensity orientation and degree of anisotropy of stress tensor. The degenerate into two material composites or pure materials.

> Quasiconvex envelope is a multifaced surface; we obtain analytic expressions for all the faces.

Example of three-material design Bridge #1







Bridge #2, smaller total weight





Cantilever, loaded at the middle





Cantilever, loaded at the the edge





Conclusion (List of unsolved problems)

- The technique for finding and proving inequalities on optimal fields needs to be developed.
- More examples of multimaterial optimal structures are needed.
- Techniques for several materials optimal design is to be developed and applied to elastic and electromagnetic materials.
- Quasiconvex envelope can be analytically described, 3d optimal structures are doable and cool.

Thank you!

Conjectured Range of the fields in an optimal 3d 3material structure



Thank you!



Next evolution: Conjuncted structures

- In the limiting structure, the current in k1 is two–dimensional, the normal current is zero.
- Therefore, the next evolution of optimal structures, when m1 decreases, likely will follow the evolution of 2d optimal structures:



enveloping composite, moving from the central part k1 moves into radial layers

k2 also moves into radial layers, replaces k1 One can check, that the resulting k2-k3 structure is optimal and equivalent to coated spheres

Bulk blocks assemblage - 2

$$\begin{aligned} K2 &= \begin{pmatrix} k_2 & 0 & 0 \\ 0 & k_2 & 0 \\ 0 & 0 & k_2 \end{pmatrix}, \\ Kc &= \begin{pmatrix} kca & 0 & 0 \\ 0 & kcb & 0 \\ 0 & 0 & kcb \end{pmatrix}, \quad kca &= \frac{k_1}{c_c}, \quad kcb &= k_1 \frac{2 - c_c}{c_c}, \end{aligned}$$

$$K3 = 0,$$

$$Kl = \begin{pmatrix} 0 & 0 & 0 \\ 0 & kl & 0 \\ 0 & 0 & kl \end{pmatrix}, \quad kl = k_1 c_1,$$



The structures realize sufficient optimality conditions: the variety is due to active/inactive point wise inequalities on the fields in opt. composite.

Smaller m1: conjecture.

 Notice, that k1 is located in 2d cylindrical or 1d laminate substructures. When m1 decreases, k1 can be replaced by k2, as in 2d structures.



• One can check, that the limiting k2-k3 bulk structure realizes HS bounds for two materials $m_1 \rightarrow 0$ (is equivalent to the coated spheres)