

**Lecture Notes on Multiparameter Processes:
Ecole Polytechnique Fédérale de Lausanne,
Switzerland**

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April–June 2001

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Preface

These are the notes for a one-semester course based on ten lectures given at the *Ecole Polytechnique Fédérale de Lausanne*, April–June 2001. My goal has been to illustrate, in some detail, some of the salient features of the theory of multiparameter processes and in particular, Cairoli’s theory of multiparameter martingales. In order to get to the heart of the matter, and develop a kind of intuition at the same time, I have chosen the simplest topics of random walks, Brownian motions, etc. to highlight the methods. The full theory can be found in *Multi-Parameter Processes: An Introduction to Random Fields* (henceforth, referred to as MPP) which is to be published by *Springer-Verlag*, although these lectures also contain material not covered in the mentioned book.

The first eight lectures are introductory material and everything is covered in some detail. The last two lectures are about potential theory of processes; in order to cover enough material, I have decided not to cover balayage; this makes the arguments of these two lectures somewhat heuristic, but the details are all there in MPP for those who are interested to learn more about the subject.

While I have tried to write the lectures faithfully, the heuristic “clumping” picture that I have tried to develop in the lectures is sadly missing from these notes. My hope is that, if all else is forgotten, one theme will be remembered, and that is the connections between our heuristic notion of clumping, and the rôle played by multiparameter martingales. Some of these notions are being rigorized in joint work with Yimin Xiao (under the general heading of *locally grown random sets*) as I write, but this will have to wait for now.

I am grateful to EPF-L for their hospitality, in particular, to Professor Robert C. Dalang, and for making my visit possible as well as highly enjoyable. Also, my heartfelt thanks go to all of the attendees of these lectures. Working with them has been a distinct pleasure for me. Last but not least, I wish to thank the United States’ *National Science Foundation* for the generous support of my work on the topic of these lectures during the past six to seven years.

D. Khoshnevisan
June 2001
Lausanne, Switzerland

Lecture 1

Examples from Markov chains

In these lectures, I will develop some of the foundations of a theory of multiparameter Markov processes that is motivated by a number of problems coming from

- probability (intersections of processes, Brownian sheet, percolation on trees, Markov chains, potential theory);
- mathematical physics (Schrödinger operators, the $(\frac{1}{2}\Delta)^2$ operator of elasticity, polymer measures); as well as
- recent advances in geology (use of stable sheets in modeling rock strata instead of percolation models).

We will see that even in the study of classical stochastic processes, multiparameter processes arise as natural objects, although these connections sometimes go unnoticed.

Some of the material of this course is based on my forthcoming book *Multi-Parameter Processes: An Introduction to Random Fields* (Springer). I will make the relevant portions of this book available as needed during the lectures.

A rough outline of this course is as follows:

- Introduction and motivating examples from Markov chains, percolation, and Brownian motion;
- Capacity, energy, and Hausdorff dimension;
- Cairoli's theory of multiparameter martingales;
- Multiparameter Markov processes;
- The Fitzsimmons–Salisbury and the Hirsch–Song theories of potential;
- Brownian sheet and potential theory; Kahane's problem.

Time permitting, and depending on the audience's interests, we may discuss some aspects of

- Probability on sigma-finite spaces, Lévy processes and random walks; the Kesten–Bretagnolle criterion for Lévy processes ;
- analysis of sample paths, stochastic codimension.

Many of the techniques in these notes are based on various applications of the theory of martingales. In the one-parameter setting, this is Doob’s theory, and in the multi-parameter setting, it is Cairoli’s. Many of you may know that R. Cairoli was a Professor here at *EPF–Lausanne*, and his theory evolved here, locally.

In the first few lectures, we will play with some of the fundamental concepts developed later in this course, but in a simple setting, where it is easy to see what is going on, without need for theoretical developments.

1 Recurrence of Markov Chains

Let S be denumerable, and consider a Markov chain $X = \{X_n; n \geq 0\}$ on S . Recall that this means there are probabilities $\{\mathbb{P}_x; x \in S\}$, and *transition functions* $\{p_n; n \geq 0\}$, on $S \times S$, such that

- $\mathbb{P}_x\{X_0 = x\} = 1$, for all $x \in S$;
- $\mathbb{P}_x\{X_n = a | \mathcal{F}_k\} = \mathbb{P}_x\{X_n = a | X_k\} = p_{n-k}(X_k, a)$, \mathbb{P}_x -a.s. for all $n \geq k \geq 0$ and all $a \in S$, where $\mathcal{F} = \{\mathcal{F}_j; j \geq 0\}$ is the filtration generated by X .

One thinks of $p_n(a, b)$ as the probability that, in n time steps, the Markov chain goes from a to b .

Following Poincaré, Pólya, Chung, etc., we say that $x \in S$ is *recurrent*, if starting from $x \in S$, $n \mapsto X_n$ hits x infinitely often, with probability one. More precisely, x is recurrent if

$$\mathbb{P}_x\{\forall m \geq 0, \exists n \geq m : X_n = x\} = 1.$$

When is x recurrent? The classical condition of Pólya (for simple walks), Chung–Fuchs, etc. is

Theorem 1.1 (Pólya’s Criterion) $x \in S$ is recurrent if and only if $\sum_n p_n(x, x) = +\infty$.

A rough explanation of the proof that is to come is needed. Namely, our proof uses the fact that if X hits x , it will do so several times in close proximity of one another. Thus, near the times when X hits x , one expects to observe an unusual contribution to the $\sum_j \mathbf{1}_{\{X_j=x\}}$. In even rougher terms, the random set $\{n : X_n = x\}$ is comprised of a bunch of i.i.d.–looking “clumps”. The Paley–Zygmund inequality provides us with a key method in analysing clumping situations. It states

Lemma 1.2 (Paley–Zygmund’s Inequality) Suppose $Z \geq 0$ is a nontrivial random variable in $L^2(\mathbb{P})$. Then, for all $\varepsilon \in [0, 1]$,

$$\mathbb{P}(Z > \varepsilon \mathbb{E}\{Z\}) \geq (1 - \varepsilon)^2 \frac{(\mathbb{E}\{Z\})^2}{\mathbb{E}\{Z^2\}}.$$

Proof Using Cauchy–Schwarz inequality,

$$\begin{aligned} \mathbb{E}\{Z\} &\leq \varepsilon \mathbb{E}\{Z\} + \mathbb{E}(Z; Z > \varepsilon \mathbb{E}\{Z\}) \\ &\leq \varepsilon \mathbb{E}\{Z\} + \sqrt{\mathbb{E}\{Z^2\} \cdot \mathbb{P}(Z > \varepsilon \mathbb{E}\{Z\})}. \end{aligned}$$

Just solve. □

Proof of Theorem 1.1 Let $J_n = \sum_{j=0}^n \mathbf{1}_{\{X_j=x\}}$ to see that x is recurrent if and only if $\mathbb{P}_x\{J_\infty = \infty\} = 1$. Clearly,

$$\sum_n p_n(x, x) < +\infty \iff \mathbb{E}_x\{J_\infty\} < +\infty \implies J_\infty < +\infty, \mathbb{P}_x\text{-a.s.}$$

This is one half of the result. “Clumping” says that this should be sharp. To prove it, we assume that

$$\sum_n p_n(x, x) = \mathbb{E}_x\{J_\infty\} = +\infty,$$

and first notice that

$$\begin{aligned} \mathbb{E}_x\{J_n^2\} &\leq 2 \sum_{0 \leq i \leq j \leq n} \mathbb{P}_x\{X_i = x, X_j = x\} \\ &= 2 \sum_{0 \leq i \leq j \leq n} p_i(x, x) p_{j-i}(x, x) \\ &\leq 2 [\mathbb{E}_x\{J_n\}]^2. \end{aligned}$$

Of course, one always has $\|Z\|_2 \geq \|Z\|_1$. The above states that a kind of converse to this holds for $Z = J_n$.

This, together with the Paley–Zygmund inequality, shows for all $n \geq 1$ and all $\varepsilon \in [0, 1]$,

$$\mathbb{P}_x(J_n \geq \varepsilon \mathbb{E}_x\{J_n\}) \geq \frac{(1 - \varepsilon)^2}{2}.$$

Thus, using our assumption that $\mathbb{E}_x\{J_\infty\} = +\infty$, and choosing any $\varepsilon > 0$,

$$\begin{aligned} \mathbb{P}_x\{J_\infty = +\infty\} &= \lim_{n \rightarrow \infty} \mathbb{P}_x(J_\infty \geq \varepsilon \mathbb{E}_x\{J_n\}) \\ &\geq \liminf_{n \rightarrow \infty} \mathbb{P}_x(J_n \geq \varepsilon \mathbb{E}_x\{J_n\}) \\ &\geq \frac{1}{2}(1 - \varepsilon)^2. \end{aligned}$$

This being true for all $\varepsilon \in (0, 1]$, we have shown that $\mathbb{P}_x\{J_\infty = +\infty\} \geq \frac{1}{2}$. We need to show this probability is 1. Usually, this step is done by stopping time arguments. Instead, we make an argument that has multiparameter extensions, and this is where martingales come in. Henceforth, we assume $\mathbb{P}_x\{J_\infty = \infty\}$ is positive, and strive to show that it is one.

Let us define $M_n = \mathbb{P}_x\{J_\infty = +\infty \mid \mathcal{F}_n\}$, and apply Doob's martingale convergence theorem. Indeed, \mathcal{F}_n increases to $\bigvee_n \mathcal{F}_n$ which means that $\lim_n M_n = \mathbf{1}_{\{J_\infty = \infty\}}$, \mathbb{P}_x -a.s. On the other hand, by the Markov property, $M_n = \mathbb{P}_{X_n}\{J_\infty = +\infty\}$ (why?) Therefore, on the event that X_n hits x infinitely often (i.e., on $\{J_\infty = +\infty\}$), $M_n = \mathbb{P}_x\{J_\infty = +\infty\}$, infinitely often. In particular, \mathbb{P}_x -a.s. on $\{J_\infty = \infty\}$,

$$1 = \mathbf{1}_{\{J_\infty = \infty\}} = \lim_n M_n = \mathbb{P}_x\{J_\infty = \infty\}.$$

This finishes our proof. □

SOMETHING TO TRY. Show that there are constants C_p such that for all $p > 1$ and all $n \geq 1$,

$$\|J_n\|_p \leq C_p \|J_n\|_1.$$

Thus, $\{J_n\}_{n \geq 1}$ forms what is called a uniformly hypercontractive family of operators on the underlying probability space. Hypercontractivity is known to be a powerful property in analysis.

(HINT. By Jensen's inequality, it suffices to do this for integers $p > 1$.)

2 Recurrence for Inhomogeneous Markov Chains

This subsection may not be in the lecture and can be omitted on first reading.

Recall that a process $\{X_n; n \geq 0\}$ is an *inhomogeneous Markov chain* if $\mathbb{P}_x\{X_n = a \mid \mathcal{F}_m\} = \mathbb{P}_x\{X_n = a \mid X_m\} = p_{m,n}(X_m, a)$, when $n > m$. Thus, the only difference between these and regular Markov chains is that the transition functions, $p_{m,n}$, need not satisfy $p_{m,n}(a, b) = p_{0,n-m}(a, b)$.

SOMETHING TO TRY. If X is an inhomogeneous Markov chain on S , the process $Y = \{(n, X_n); n \geq 0\}$ defines a homogeneous Markov chain on $\mathbb{Z} \times S$. The process Y is Doob's *space-time* process.

What happens when is x recurrent for an inhomogeneous Markov chain X ? Perhaps not surprisingly, the answer is more complicated, although similar ideas as our proof of Theorem 1.1 still work. Namely, recurrence still occurs by clumping. However, due to the inhomogeneity of the chain, the clumps need not be evenly distributed. In more mathematical terms, we need to consider a weighted version of the number of times that the chain hits x . Due to the absence of 0–1 laws, we also revise our definition of recurrence and say that x is *recurrent* if $\mathbb{P}_x\{J_\infty = \infty\} > 0$, where J_∞ was defined in the previous subsection.

Supposing that F is a subset of $\mathbb{Z} = \{0, 1, \dots\}$, we write $\mu \in \mathcal{P}(F)$ meaning that $\{\mu_i; i \in F\}$ are probabilities: $\mu_i \in [0, 1]$ and $\sum_{i \in F} \mu_i = 1$. By the *energy* at x of such a μ , we mean the functional

$$\mathcal{E}_x(\mu) = \sum_{\substack{0 \leq k \leq j \\ p_{0,j}(x,x) > 0}} \frac{p_{k,j}(x,x)}{p_{0,j}(x,x)} \mu_j \mu_k,$$

where, we recall, $p_{k,j}$'s are the transition functions of the inhomogeneous Markov chain X .

The *capacity* at x of F is defined by Gauss' minimum energy principle. Namely,

$$\mathcal{C}_x(F) = \left[\inf_{\mu \in \mathcal{P}(F)} \mathcal{E}_x(\mu) \right]^{-1},$$

where $\inf \emptyset = +\infty$, and $1 \div \infty = 0$. Then, we have the following quantitative estimate that is interesting even for homogeneous Markov chains:

Theorem 2.1 *For all $F \subseteq \mathbb{Z}$ and all $x \in S$,*

$$\frac{1}{2} \mathcal{C}_x(F \cap G_x) \leq \mathbb{P}_x \{ \exists n \in F : X_n = x \} \leq \mathcal{C}_x(F \cap G_x),$$

where $G_x = \{i \geq 0 : \mathbb{P}_x[X_i = x] > 0\}$.

In particular, x is recurrent if and only if $\lim_{N \rightarrow \infty} \mathcal{C}_x(\{N, N+1, \dots\}) > 0$.

Lecture 2

Examples from Percolation on Trees and Brownian Motion

1 Proof of Theorem 2.1

The characterization of recurrence follows from the asserted inequalities, upon noticing that whenever $F \subset F'$, $\mathcal{C}(F) \leq \mathcal{C}(F')$, so that $\lim_N \mathcal{C}(\{N, N+1, \dots\})$ exists. It remains to verify the two inequalities.

For any $\mu \in \mathcal{P}(F \cap G_x)$, define

$$I_\mu = \sum_{i \in G_x} \frac{\mathbf{1}_{\{X_i=x\}}}{\mathbb{P}_x\{X_i=x\}} \mu_i.$$

Clearly, $\mathbb{P}_x\{X_i=x\} = p_{0,i}(x,x)$, and

$$\mathbb{E}_x\{I_\mu\} = 1. \tag{1.1}$$

This functional I_μ is a normalized, weighted, version of J_n in our proof of Theorem 1.1.1. With this in mind, we estimate its second moment in a similar way as we did that of J_n . Namely,

$$\begin{aligned} \mathbb{E}_x\{I_\mu^2\} &\leq 2 \sum_{\substack{i,j \in G_x: \\ i \leq j}} \frac{\mathbb{P}_x\{X_i = X_j = x\}}{p_{0,i}(x,x)p_{0,j}(x,x)} \mu_i \mu_j \\ &= 2\mathcal{E}_x(\mu), \end{aligned} \tag{1.2}$$

since $\mathbb{P}_x\{X_i = X_j = x\} = p_{0,i}(x,x)p_{0,j}(x,x)$ if $i \leq j$. Eq. (1.1) and the above combine, thanks to the Paley–Zygmund inequality [Lemma 1.1.2], to show that for all $\mu \in \mathcal{P}(F \cap G_x)$,

$$\mathbb{P}_x\{\exists n \in F : X_n = x\} \geq [2\mathcal{E}_x(\mu)]^{-1}.$$

Optimize over all $\mu \in \mathcal{P}(F)$ to derive the first inequality of Theorem 1.2.1. So far, we have followed our proof of Theorem 1.1.1 but with J_n replaced by I_μ . For the other bound, we need to do more work, since not all weights μ work.

For this bound, we can assume, without loss of generality, that $\mathbb{P}_x\{\exists n \in F : X_n = x\} > 0$. Let $T = \inf\{n \in F : X_n = x\}$, where $\inf \emptyset = +\infty$. Then, T is a stopping time and our assumption is equivalent to $\mathbb{P}_x\{T < \infty\} > 0$. In particular, $\mu \in \mathcal{P}(F \cap G_x)$, where we define

$$\mu_i = \mathbb{P}_x\{T = i \mid T < \infty\}, \quad \forall i = 0, 1, 2, \dots$$

We consider the following martingale based on I_μ for *this* μ :

$$M_n = \mathbb{E}_x\{I_\mu \mid \mathcal{F}_n\}, \quad \forall n \geq 0.$$

Clearly,

$$\begin{aligned} M_n &= \sum_{i \in G_x} \frac{\mathbb{P}_x\{X_i = x \mid \mathcal{F}_n\}}{p_{0,i}(x, x)} \mu_i \\ &\geq \sum_{\substack{i \in G_x: \\ i \geq n}} \frac{\mathbb{P}_x\{X_i = x \mid \mathcal{F}_n\}}{p_{0,i}(x, x)} \mu_i \\ &= \sum_{\substack{i \in G_x: \\ i \geq n}} \frac{p_{n,i}(X_n, x)}{p_{0,i}(x, x)} \mu_i \\ &\geq \sum_{\substack{i \in G_x: \\ i \geq n}} \frac{p_{n,i}(x, x)}{p_{0,i}(x, x)} \mu_i \cdot \mathbf{1}_{\{X_n = x\}}. \end{aligned}$$

In particular,

$$M_T \mathbf{1}_{\{T < \infty\}} \geq \sum_{\substack{i \in G_x: \\ i \geq T}} \frac{p_{T,i}(x, x)}{p_{0,i}(x, x)} \mu_i \cdot \mathbf{1}_{\{T < \infty\}}.$$

Taking expectations, and using the special form of μ gives

$$\mathbb{E}_x\{M_T; T < \infty\} \geq \mathcal{E}_x(\mu) \cdot \mathbb{P}_x\{T < \infty\}.$$

It remains to show that the left hand side is 1. But,

$$\begin{aligned} \mathbb{E}_x\{M_T; T < \infty\} &\leq \liminf_{n \rightarrow \infty} \mathbb{E}_x\{M_{T \wedge n}\} \\ &= \liminf_{n \rightarrow \infty} \mathbb{E}_x\{M_0\} \\ &= 1. \end{aligned}$$

Justification: the first inequality is from Fatou's lemma, the second equality holds by Doob's optional stopping theorem, and the third is from Eq. (1.1) above. This concludes our proof. \square

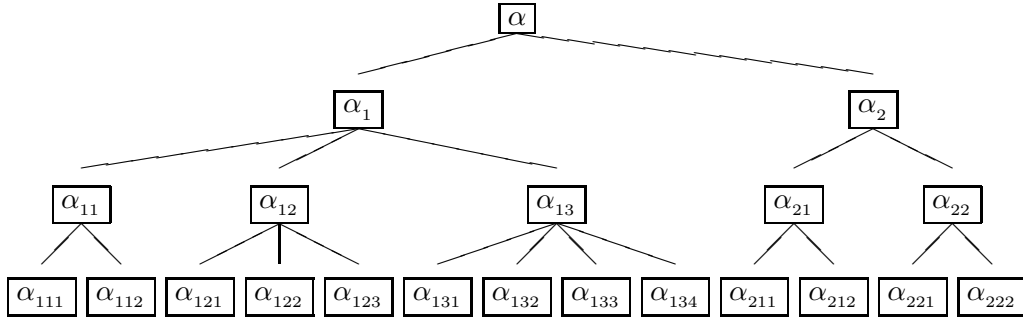


Figure 2.1: A typical ancestral tree Γ

2 Lyons' Theorem

Consider a finite rooted tree Γ , viewed as an ancestral tree, and with root α . An example of such a tree can be found in Figure 2.1, where α gives birth to two children, α_1 and α_2 ; α_1 gives birth to α_{11} , α_{12} and α_{13} , while α_2 gives birth to α_{21} and α_{22} , who give birth to \dots .

The indices are chosen in a natural way, as is often done, to keep track of the ancestry of any given individual on the tree. In fact, this orders the vertices of Γ .

Every edge e in the tree is assigned a random weight, $w(e)$, and we suppose that $\{w(e); e \in \text{Edge}(\Gamma)\}$ is an independent collection of random variables. In the percolation setting, $w(e) \in \{0, 1\}$; we can think of e as “open” when $w(e) = 1$, and “closed” when $w(e) = 0$.

We will assume that the weights $w(e)$ are i.i.d., to make our presentation simpler. Thus, there exists an inherent parameter, p , which is the probability $p = \mathbb{P}\{w(e) = 1\}$. To emphasize the dependence on the parameter p , we write P_p for \mathbb{P} .

Now, we define percolation on Γ .

If β is some vertex in Γ , we write $\alpha \leftrightarrow \beta$ for the event that for all edges, e , that link α to β , $w(e) = 1$. If A is a collection of vertices, we write $\alpha \leftrightarrow A$ for the event $\cup_{\beta \in A} \{\alpha \leftrightarrow \beta\}$.

Let $\partial\Gamma$ denote the collection of all vertices in Γ whose graph distance from α is maximal. (In Figure 2.1, $\partial\Gamma = \{\alpha_{ijk}\}$ where the indices i, j and k range over the values 1 – 4, as allowed by the the figure.) We then say that *percolation occurs on the finite tree* Γ , if $\alpha \leftrightarrow \partial\Gamma$.

Let us first look at a finite rooted tree Γ with root α , and let its depth be D .¹ [In Figure 2.1, $D = 3$.] We can label, from left to right, the vertices in $\partial\Gamma$ to get vertices $1, \dots, N$, where N is the number of vertices in $\partial\Gamma$. Thus, in Figure 2.1, we are relabeling $\alpha_{111}, \alpha_{112}, \dots, \alpha_{221}, \alpha_{222}$ as vertices 1 through $N = 13$. For each of these N vertices, we can define a D -dimensional random vector, X_i , that is comprised of all the random weights, $w(e)$, for e 's that link α to vertex i in $\partial\Gamma$. Then, it is easy to see that X_1, \dots, X_N is an inhomogeneous Markov chain, as long as we define $X_0 = x$ under the measure \mathbb{P}_x , where x takes its values in the space of all possible configurations of 1's and 0's corresponding to the rays from α to $\partial\Gamma$. In summary,

¹Recall that D is the number of edges needed to go from the root to the boundary.

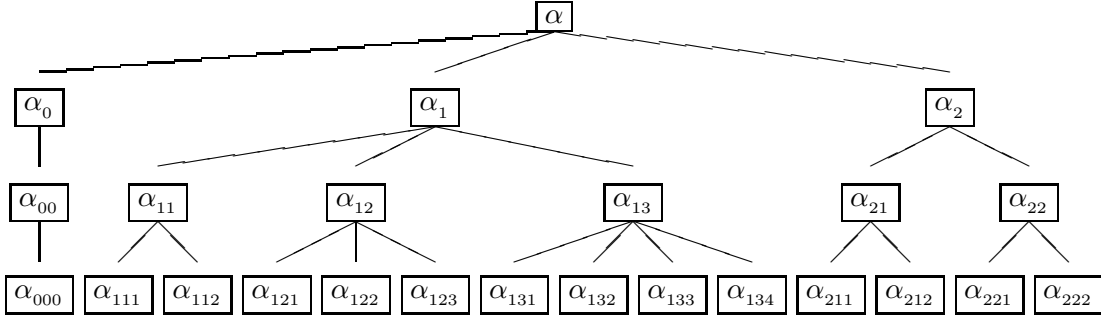


Figure 2.2: The ancestral tree of Figure 2.1, with an attached fictitious first ray

$x \in S = \{0, 1\}^D$. Pictorially, this means that for X to be an inhomogeneous Markov chain, we need to include an extra fictitious first ray that goes from α to $\partial\Gamma$ with no interference with the other rays.

Figure 2.2 shows what we do to Figure 2.1 in order to achieve this.

Now, it is easy to see that X is an inhomogeneous Markov chain on the state space $S = \{0, 1\}^D$.

We are interested in probability of percolation, which is

$$P_p\{\alpha \leftrightarrow \partial\Gamma\} = \mathbb{P}_x\{\exists n \in \{1, \dots, N\} : X_n = x\},$$

where $x = \overbrace{(1, \dots, 1)}^{D \text{ times}}$.

Next, we compute the transition probabilities of this chain. Note that for $x = (1, \dots, 1)$, as given, $p_{0,j}(x, x) = p^D$, since the above is just the probability that the ray leading to the j -th vertex on $\partial\Gamma$, counting from left, is all open. On the other hand, if $0 \leq k < j$, $p_{k,j}(x, x)$ is the probability that both rays leading to k and j are all open. This probability is $p_{k,j}(x, x) = p^{D-|k \wedge j|}$, where $k \wedge j$ stands for the greatest common ancestor of k and j ,² and where $|v|$ denotes the depth of any vertex v . In this way, we obtain

$$\frac{p_{0,j}(x, x)}{p_{k,j}(x, x)} = p^{-|k \wedge j|}.$$

Since $G_x = \partial\Gamma$, we can combine all this with Theorem 2.1 to obtain

$$\left[2 \inf_{\mu \in \mathcal{P}(\partial\Gamma)} \sum_{1 \leq k \leq j \leq N} p^{-|k \wedge j|} \mu_k \mu_j \right]^{-1} \leq P_p\{\alpha \leftrightarrow \partial\Gamma\} \leq \left[\inf_{\mu \in \mathcal{P}(\partial\Gamma)} \sum_{1 \leq k \leq j \leq N} p^{-|k \wedge j|} \mu_k \mu_j \right]^{-1}.$$

²This is a vertex on the tree, and is not to be mistaken with the number $k \wedge j$. For example, in Figure 2.1, $\alpha_{111} \wedge \alpha_{134} = \alpha_1$.

The notation $k \leq j$ is awkward, since it depends on a linear order of the boundary of Γ , which makes no sense when we take Γ to be an infinite tree. So, it is more convenient to use the inequalities:

$$\frac{1}{2} \sum_{1 \leq k \leq j \leq N} \sum_{1 \leq k \leq j \leq N} \leq \sum_{1 \leq k \leq j \leq N} \sum_{1 \leq k \leq j \leq N} \leq \sum_{1 \leq k \leq j \leq N} \sum_{1 \leq k \leq j \leq N},$$

to get

$$\left[2 \inf_{\mu \in \mathcal{P}(\partial\Gamma)} \sum_{k,j} \mathfrak{p}^{-|k \wedge j|} \mu_k \mu_j \right]^{-1} \leq P_{\mathfrak{p}}\{\alpha \leftrightarrow \partial\Gamma\} \leq 2 \left[\inf_{\mu \in \mathcal{P}(\partial\Gamma)} \sum_{k,j} \mathfrak{p}^{-|k \wedge j|} \mu_k \mu_j \right]^{-1}.$$

When the tree is an infinite tree, the boundary $\partial\Gamma$ still makes sense: it is defined as the collection of all infinite rays emanating from α . It should be checked that when Γ is a finite tree, this agrees with our previous definition. Moreover, the notion of percolation still makes sense, as well. Namely, we write $\alpha \leftrightarrow \partial\Gamma$ to mean that for all rooted finite subtrees $\Gamma' \subset \Gamma$ with root α , $\alpha \leftrightarrow \partial\Gamma'$. It turns out that one can take “hydrodynamical limits”, and let Γ become infinite. In this way, we get an improvement, due to Benjamini and Peres, of

Theorem 2.1 (Lyons) *Consider an infinite, locally finite, rooted tree Γ with root α as above. Then, for any $\mathfrak{p} \in (0, 1)$,*

$$\left[2 \inf_{\mu \in \mathcal{P}(\partial\Gamma)} \iint \mathfrak{p}^{-|\sigma \wedge \gamma|} \mu(d\sigma) \mu(d\gamma) \right]^{-1} \leq P_{\mathfrak{p}}\{\alpha \leftrightarrow \partial\Gamma\} \leq 2 \left[\inf_{\mu \in \mathcal{P}(\partial\Gamma)} \iint \mathfrak{p}^{-|\sigma \wedge \gamma|} \mu(d\sigma) \mu(d\gamma) \right]^{-1}.$$

In particular, there can be percolation iff there exists a probability measure μ , on $\partial\Gamma$, such that $\iint \mathfrak{p}^{-|\gamma \wedge \sigma|} \mu(d\sigma) \mu(d\gamma) < +\infty$.

SOMETHING TO TRY: Check that the first inequality holds even if we remove the constant 2. (Hint: compute the first two moments of $I = \int \mathbf{1}_{\{\alpha \leftrightarrow \sigma\}} \div P_{\mathfrak{p}}\{\alpha \leftrightarrow \sigma\} \mu(d\sigma)$, directly.

The above can be used to find critical percolation probabilities, as well as critical exponents, when the tree is regularly behaved. For example, if Γ is k -ary (or, more generally, radially symmetric), a μ that solves the optimization problem above is uniform measure on Γ , and this makes exact calculations possible; cf. Lyons (1990).

3 Brownian Motion

Next, we study another example of martingale techniques for clumping analysis by deriving a deep theorem of P. Lévy on the curve of a d -dimensional Brownian motion, $B = \{B(t); t \geq 0\}$, with $B(0) = 0$. Namely, we will show the following:

Theorem 3.1 (P. Lévy) *The random set $B(\mathbb{R}_+) \subset \mathbb{R}^d$ has positive d -dimensional Lebesgue's measure if and only if $d = 1$.*

When $d = 1$, this follows from the continuity of B , and from the fact that it is not a constant function. Thus, it suffices to reduce attention to $d \geq 2$ and show that for any interval $[L, U]$ where $0 < L < U < \infty$, $|B([L, U])| = 0$, almost surely. To see what this has to do with clumping, we first note that thank to Fubini's theorem,

$$\mathbb{E}\{|B([L, U])|\} = \int_{\mathbb{R}^d} \mathbb{P}\{a \in B([L, U])\} da, \quad (3.1)$$

where $|\bullet|$ refers to Lebesgue's measure. Thus, the question is “when is $\mathbb{P}\{a \in B([L, U])\} > 0$?” Now, it should not be surprising to see clumping at work: while $\mathbb{P}\{B(t) \approx a\}$ is small for any given t , $\mathbb{P}\{B(s) \approx a | B(t) \approx t\}$ is not. So, for $\{t \in [L, U] : B(t) = a\}$ to have any chance of being nonempty, it would have to be made up of clumps. Temporal homogeneity of B suggests that these clumps are, moreover, evenly distributed. All of this suggests looking at the random variable

$$J_\varepsilon(a) = \int_L^U \mathbf{1}_{\{|B(t)-a|\leq\varepsilon\}} ds, \quad (3.2)$$

where $a \in \mathbb{R}^d$ is fixed, and we think of ε as small.

Lecture 3

Proving Lévy's Theorem and Introducing Martingales

In this lecture, I will prove Theorem 3.1, and then, we will start our discussion of multiparameter martingales in earnest. For the latter part, I will pass out parts of my book, for the former ... please read on.

1 Proof of Lévy's Theorem

Here is one estimate that follows easily from the form of the Gaussian density function.

Lemma 1.1 *For any $a \in \mathbb{R}^d$, there exists positive and finite constants $C_1 = C_1(a, U)$ and $C_2 = C_2(a, U)$, such that for all $s \in [0, U]$,*

$$C_1 \left(\frac{\varepsilon}{\sqrt{s}} \wedge 1 \right)^d \leq \mathbb{P}\{|B(s) - a| \leq \varepsilon\} \leq C_2 \left(\frac{\varepsilon}{\sqrt{s}} \wedge 1 \right)^d.$$

Consequently, since L is *strictly* positive, there exists nontrivial $C_3 = C_3(a, L, U)$ and $C_4 = C_4(a, L, U)$, such that for all sufficiently small $\varepsilon > 0$,

$$C_3 \varepsilon^d \leq \mathbb{E}\{J_\varepsilon(a)\} \leq C_4 \varepsilon^d. \tag{1.1}$$

Next, we estimate the second moment of $J_\varepsilon(a)$, viz.,

$$\begin{aligned}
\mathbb{E}\{|J_\varepsilon(a)|^2\} &= 2 \int_L^U \int_s^U \mathbb{P}\{|B(s) - a| \leq \varepsilon, |B(t) - a| \leq \varepsilon\} dt ds \\
&\leq 2 \int_L^U \int_s^U \mathbb{P}\{|B(s) - a| \leq \varepsilon, |B(t) - B(s)| \leq 2\varepsilon\} dt ds \\
&\leq 2 \int_L^U \int_s^U \mathbb{P}\{|B(s) - a| \leq \varepsilon\} \mathbb{P}\{|B(t-s)| \leq 2\varepsilon\} dt ds \\
&\leq 2\mathbb{E}\{J_\varepsilon(a)\} \cdot \int_0^U \mathbb{P}\{|B(u)| \leq 2\varepsilon\} du \\
&\leq 2C_2C_4\varepsilon^d \int_0^U \left(\frac{2\varepsilon}{\sqrt{u}} \wedge 1\right)^d du.
\end{aligned}$$

Here is where things get interesting. When $d \geq 3$, the behavior of the above integral is like a constant times ε^2 . But, if $d = 2$, it behaves like $\varepsilon^2 \log(1/\varepsilon)$. Finally, if $d = 1$, the integral behaves like ε (times a constant). All considered, we get

$$\mathbb{E}\{|J_\varepsilon(a)|^2\} \leq C_5 \times \begin{cases} \varepsilon^2, & \text{if } d = 1 \\ \varepsilon^2 \log(1/\varepsilon), & \text{if } d = 2. \\ \varepsilon^{d+2}, & \text{if } d \geq 3 \end{cases}$$

Combine this with Eq. (1.1) and the Paley–Zygmund inequality [Lemma 1.1.2] to get

$$\begin{aligned}
\mathbb{P}\left\{\inf_{t \in [L, U]} |B(t) - a| \leq \varepsilon\right\} &\geq \mathbb{P}\{J_\varepsilon(a) > 0\} \\
&\geq \frac{|\mathbb{E}\{J_\varepsilon(a)\}|^2}{\mathbb{E}\{|J_\varepsilon(a)|^2\}} \\
&\geq \frac{C_3^2}{C_5} \times \begin{cases} 1, & \text{if } d = 1 \\ [\log(1/\varepsilon)]^{-1}, & \text{if } d = 2. \\ \varepsilon^{d-2}, & \text{if } d \geq 3 \end{cases} \quad (1.2)
\end{aligned}$$

I claim these bounds are sharp, up to multiplicative constants. Let \mathcal{F}_t denote the filtration of B , and assume it satisfies the usual conditions. Let us consider

$$M_t = \mathbb{E}\left\{\int_L^{L+U} \mathbf{1}_{\{|B(s)-a| \leq 2\varepsilon\}} ds \mid \mathcal{F}_t\right\}.$$

This is almost the same as $\mathbb{E}\{J_{2\varepsilon}(a) | \mathcal{F}_t\}$, but we have increased the upper limit of integration for some elbow room; you will see why shortly. Clearly,

$$M_t \geq \int_t^{L+U} \mathbb{P}\{|B(s) - a| \leq 2\varepsilon \mid \mathcal{F}_t\} ds, \quad \forall t \in [L, U].$$

You may be wondering about the null sets. If so, that is good. However, there is a general fact about Brownian motion that states that $\mathbb{E}\{Z|\mathcal{F}_t\}$ can be chosen to be continuous. So, the above holds simultaneously for all $t \in [L, U]$, outside one null set. We go one more step:

$$\begin{aligned}
M_t &\geq \int_t^{L+U} \mathbb{P}\{|B(s) - a| \leq 2\varepsilon \mid \mathcal{F}_t\} ds \times \mathbf{1}_{\{|B(t)-a|\leq\varepsilon\}}, & \forall t \in [L, U] \\
&\geq \int_t^{L+U} \mathbb{P}\{|B(s) - B(t)| \leq \varepsilon \mid \mathcal{F}_t\} ds \times \mathbf{1}_{\{|B(t)-a|\leq\varepsilon\}}, & \forall t \in [L, U] \\
&= \int_t^{L+U} \mathbb{P}\{|B(s - t)| \leq \varepsilon\} ds \times \mathbf{1}_{\{|B(t)-a|\leq\varepsilon\}}, & \forall t \in [L, U] \\
&\geq \int_0^U \mathbb{P}\{|B(u)| \leq \varepsilon\} du \times \mathbf{1}_{\{|B(t)-a|\leq\varepsilon\}}, & \forall t \in [L, U].
\end{aligned}$$

The point is that the above integral is nonrandom, and is an object we have seen before: when $d = 1$, it is of order ε^2 , when $d = 2$, it is of order $\varepsilon^2 \log(1/\varepsilon)$, and when $d \geq 3$, it is of order ε^{d+2} . So, there must exist some constant $C_6 = C_6(a, d, U, L)$, such that

$$M_t \geq C_6 \mathbf{1}_{\{|B(t)-a|\leq\varepsilon\}} \times \begin{cases} \varepsilon, & \text{if } d = 1 \\ \varepsilon^2 \log(1/\varepsilon), & \text{if } d = 2, \\ \varepsilon^{d+2}, & \text{if } d \geq 3 \end{cases} \quad \forall t \in [L, U].$$

(We needed the extra ‘‘elbow room’’ to get $\forall t \in [L, U]$ in the above.) Now, let $T = \inf\{t \in [L, U] : |B(t) - a| \leq \varepsilon\}$ to see that

$$\mathbb{E}\{M_T \mathbf{1}_{\{T < \infty\}}\} \geq C_6 \mathbb{P}\left\{ \inf_{t \in [L, U]} |B(t) - a| \leq \varepsilon \right\} \times \begin{cases} \varepsilon, & \text{if } d = 1 \\ \varepsilon^2 \log(1/\varepsilon), & \text{if } d = 2. \\ \varepsilon^{d+2}, & \text{if } d \geq 3 \end{cases}$$

Thanks to the boundedness of M , and by Doob’s optional stopping theorem, $\mathbb{E}\{M_T; T < \infty\} = \mathbb{E}\{J_\varepsilon(a)\}$, which is estimated by Eq. (1.1). Combining this and Eq. 1.2 we get

Theorem 1.2 For any $a \in \mathbb{R}^d$, there exists constants A_1 and A_2 , such that for all $\varepsilon > 0$ small,

$$A_1 \kappa(\varepsilon) \leq \mathbb{P}\left\{ \inf_{t \in [L, U]} |B(t) - a| \leq \varepsilon \right\} \leq A_2 \kappa(\varepsilon),$$

where

$$\kappa(\varepsilon) = \begin{cases} 1, & \text{if } d = 1 \\ [\log(1/\varepsilon)]^{-1}, & \text{if } d = 2. \\ \varepsilon^{d-2}, & \text{if } d \geq 3 \end{cases}$$

In particular, since B is a continuous random function, we can let $\varepsilon \rightarrow 0$ to see that when $d \geq 2$, $\mathbb{P}\{a \in B[L, U]\} = 0$ for all a , and we get Theorem 3.1 from integrating this $[da]$.

2 Review of Martingales

Recall that $\mathcal{F} = (\mathcal{F}_k)$ is a *filtration*, if \mathcal{F}_k 's are sigma-fields such that $\mathcal{F}_1 \subseteq \mathcal{F}_2 \subseteq \dots$.

A stochastic process $X = (X_k)$ is *adapted* to the filtration \mathcal{F} , if for each k , X_k is \mathcal{F}_k -measurable.

A process $M = (M_k)$ is a *martingale*, if

- for each k , $M_k \in L^1(\mathbb{P})$;
- M is adapted; and
- for all k , $\mathbb{E}\{M_{k+1} | \mathcal{F}_k\} = M_k$, a.s.

Check that the third part is equivalent to: for all $n \geq m$, $\mathbb{E}\{M_m | \mathcal{F}_n\} = M_m$, a.s.

Proof We can assume that $n > m$. Otherwise, there is nothing to prove. Since $\mathcal{F}_n \supseteq \mathcal{F}_m$, $\mathbb{E}\{M_m | \mathcal{F}_n\} = \mathbb{E}\{\mathbb{E}\{M_m | \mathcal{F}_{m-1}\} | \mathcal{F}_n\} = \mathbb{E}\{M_{m-1} | \mathcal{F}_n\}$. Proceed by induction. \square

Doob's Martingales are those of the form $M_k = \mathbb{E}\{Y | \mathcal{F}_k\}$, where $Y \in L^1(\mathbb{P})$.

A *stopping time* T is a $\mathbb{N} \cup \{\infty\}$ -valued random variable such that for all k , $(T \leq k) \in \mathcal{F}_k$.

The stopping time property always holds with respect to some filtration \mathcal{F} , of course.

If T is a stopping time, we define

$$\mathcal{F}_T = (A \in \vee_k \mathcal{F}_k : A \cap (T \leq k) \in \mathcal{F}_k, \forall k).$$

Notes (i) \mathcal{F}_T is a sigma-field; and (ii) both T and $X_T \mathbf{1}_{\{T < \infty\}}$ are \mathcal{F}_T -measurable.

Theorem 2.1 Suppose $Y \in L^1(\mathbb{P})$ and $M_k = \mathbb{E}\{Y | \mathcal{F}_k\}$ for a given filtration \mathcal{F} . If T is a stopping time,

$$M_T = \mathbb{E}\{Y | \mathcal{F}_T\}, \quad \text{on } (T < \infty).$$

Proof For all $A \in \mathcal{F}_T$,

$$\begin{aligned} \mathbb{E}\{M_T; A \cap (T = k)\} &= \mathbb{E}\{M_k; A \cap (T = k)\} \\ &= \mathbb{E}\{\mathbb{E}(Y | \mathcal{F}_k); A \cap (T = k)\} \\ &= \mathbb{E}\{\mathbb{E}(Y; A \cap (T = k) | \mathcal{F}_k)\}, & \text{since } A \cap (T = k) \in \mathcal{F}_k \\ &= \mathbb{E}\{Y; A \cap (T = k)\} \\ &= \mathbb{E}\{\mathbb{E}(Y | \mathcal{F}_T); A \cap (T = k)\}, & \text{since } A \cap (T = k) \in \mathcal{F}_T. \end{aligned}$$

Sum over all $k = 1, 2, \dots$ to see that

$$\mathbb{E}\{M_T \mathbf{1}_{\{T < \infty\}}; A\} = \mathbb{E}\{\mathbb{E}(Y | \mathcal{F}_T) \mathbf{1}_{\{T < \infty\}}; A\}, \quad \forall A \in \mathcal{F}_T.$$

Since $\mathbb{E}(Y | \mathcal{F}_T) \mathbf{1}_{\{T < \infty\}}$ and $M_T \mathbf{1}_{\{T < \infty\}}$ are both \mathcal{F}_T -measurable, this completes our proof; cf. “Notes” above for the latter remarks. \square

Theorem 2.2 (The Optional Stopping Theorem) Suppose $T_1 \leq T_2$ are bounded stopping times, and M is a martingale, both with respect to the same filtration \mathcal{F} . Then,

$$\mathbb{E}\{M_{T_2} | \mathcal{F}_{T_1}\} = M_{T_1}, \quad \text{a.s.}$$

Proof “Bounded” means that there exists a nonrandom $K > 0$, such that $T_1 \leq T_2 \leq K$. Note that

$$\forall j \leq K : \quad M_j = \mathbb{E}\{M_{K+1} | \mathcal{F}_j\}.$$

Thus, by Theorem 2.1,

$$M_{T_2} = \sum_{j \leq K} M_j \mathbf{1}_{\{T_2=j\}} = \mathbb{E}\{M_{K+1} | \mathcal{F}_{T_2}\}, \quad \text{a.s.}$$

Since $\mathcal{F}_{T_2} \supseteq \mathcal{F}_{T_1}$, $\mathbb{E}\{M_{T_2} | \mathcal{F}_{T_1}\} = \mathbb{E}\{M_{K+1} | \mathcal{F}_{T_1}\}$, a.s. Another appeal to Theorem 2.1 does the job.

3 Doob’s Maximal Inequalities

We can now state and prove Doob’s martingale version of Kolmogorov’s inequalities for random walks.

Theorem 3.1 Suppose M is a martingale with respect to a filtration \mathcal{F} . Then, for all $\lambda > 0$,

$$\mathbb{P}\{\max_{j \leq n} |M_j| \geq \lambda\} \leq \frac{1}{\lambda} \mathbb{E}\{|M_n|; \max_{j \leq n} |M_j| \geq \lambda\}.$$

Proof Let $T = \inf\{k : |M_k| \geq \lambda\}$ where $\inf \emptyset = \infty$. This is a stopping time, and $T \wedge n$ is a bounded stopping time.

By Theorem 2.2, $\mathbb{E}\{M_n | \mathcal{F}_{T \wedge n}\} = M_{T \wedge n}$. Applying Jensen's inequality, we deduce that a.s.,

$$\mathbb{E}\{|M_n| | \mathcal{F}_{T \wedge n}\} \geq |M_{T \wedge n}|.$$

Consequently,

$$\mathbb{E}\{|M_n|; T \leq n\} \geq \mathbb{E}\{|M_T|; T \leq n\} \geq \lambda \mathbb{P}\{T \leq n\},$$

giving the result, in a slightly different form. □

Corollary 3.2 (Doob) For all $p > 1$,

$$\mathbb{E}\{\max_{j \leq n} |M_j|^p\} \leq \left(\frac{p}{p-1}\right)^p \mathbb{E}\{|M_n|^p\},$$

while

$$\mathbb{E}\{\max_{j \leq n} |M_j|\} \leq \left(\frac{e}{e-1}\right) \{1 + \mathbb{E}\{|M_n| \ln_+ |M_n|\}\},$$

where $\ln_+(x) = \ln(x \vee 1)$.

Proof We use integration by parts: for any random variable $Z \geq 0$, and for all $p \geq 1$,

$$\mathbb{E}\{Z^p\} = p \int_0^\infty \lambda^{p-1} \mathbb{P}\{Z > \lambda\} d\lambda.$$

Apply this to $Z = \max_{j \leq n} |M_j|$ to see that

$$\begin{aligned} \mathbb{E}\{\max_{j \leq n} |M_j|^p\} &= p \int_0^\infty \lambda^{p-1} \mathbb{P}\{\max_{j \leq n} |M_j| > \lambda\} d\lambda \\ &\leq p \int_0^\infty \lambda^{p-2} \mathbb{E}\{|M_n|; \lambda \leq \max_{j \leq n} |M_j|\} d\lambda \\ &= p \mathbb{E}\left\{|M_n| \cdot \int_0^{\max_{j \leq n} |M_j|} \lambda^{p-2} d\lambda\right\} \quad (\text{Fubini' theorem}) \\ &= \left(\frac{p}{p-1}\right) \mathbb{E}\{|M_n| \cdot \max_{j \leq n} |M_j|^{p-1}\}. \end{aligned}$$

By Hölder's inequality, $\mathbb{E}\{|M_n| \cdot \max_{j \leq n} |M_j|^{p-1}\} \leq \|M_n\|_p \cdot \|\max_{j \leq n} |M_j|\|_p^{p-1}$. The result follows from this when $p > 1$. For the $p = 1$ case, see Chapter 1 of MPP. □

Lecture 4

Preliminaries on Ortho-Martingales

Before our general discussion of ortho-martingales, let us look at a simpler, more concrete, class of multi-parameter “martingales”; these first arose in Cairoli (1969), and have resurfaced in many works, including those of Cairoli and Walsh, several times since, and for a number of different reasons.

1 Two Parameters Doob-Type Ortho-Martingales

Suppose \mathcal{F}^1 and \mathcal{F}^2 are two one-parameter filtrations. To be sure that we are getting the indices right, we stress that this notation means, in particular, that $\mathcal{F}^1 = (\mathcal{F}_1^1, \mathcal{F}_2^1, \dots)$ and $\mathcal{F}^2 = (\mathcal{F}_1^2, \mathcal{F}_2^2, \dots)$ are two ordinary filtrations on the same underlying probability space.

By a *two-parameter Doob-type ortho-martingale*, we mean the process

$$M_{n,m} = \mathbb{E}\{Y \mid \mathcal{F}_n^1 \vee \mathcal{F}_m^2\}, \quad \forall n, m \geq 1,$$

where $Y \in L^1(\mathbb{P})$.

We should recognize that

$$\begin{aligned} M_{n,m} &= \mathbb{E}\{M_{n+1,m} \mid \mathcal{F}_n^1 \vee \mathcal{F}_m^2\} \\ &= \mathbb{E}\{M_{n,m+1} \mid \mathcal{F}_n^1 \vee \mathcal{F}_m^2\} \\ &= \mathbb{E}\{M_{n+i,m+j} \mid \mathcal{F}_n^1 \vee \mathcal{F}_m^2\}, \end{aligned} \tag{1.1}$$

for any $i, j \geq 0$. This is why we call M a *Doob-type ortho-martingale*, since an ortho-martingale is to be thought of as an *orthant-wise martingale*, i.e., a martingale in each parameter.

Lemma 1.1 For all $p > 1$, and for all integers $n, m \geq 1$,

$$\mathbb{E}\left\{\max_{\substack{i \leq n \\ j \leq m}} |M_{i,j}|^p\right\} \leq \left(\frac{p}{p-1}\right)^{2p} \mathbb{E}\{|Y|^p\}.$$

Proof If we fix $i, j \mapsto M_{i,j}$ is a 1-parameter Doob martingale; cf. Eq. (1.1). Therefore, by Doob's inequality (Corollary 3.2, Lecture 3),

$$\mathbb{E}\left\{\max_{j \leq m} |M_{i,j}|^p\right\} \leq \left(\frac{p}{p-1}\right)^p \mathbb{E}\{|M_{i,m}|^p\}, \quad \forall i \geq 1.$$

Another application of Eq. (1.1) reveals that $i \mapsto \max_{j \leq m} |M_{i,j}|^p$ is a 1-parameter submartingale (that I have not defined in the lectures, but you know about.) Therefore, we can apply Doob's inequality again to see that

$$\mathbb{E}\left\{\max_{i \leq n} \max_{j \leq m} |M_{i,j}|^p\right\} \leq \left(\frac{p}{p-1}\right)^p \mathbb{E}\left\{\max_{j \leq m} |M_{n,m}|^p\right\},$$

which, together with the previous display, proves

$$\mathbb{E}\left\{\max_{i \leq n} \max_{j \leq m} |M_{i,j}|^p\right\} \leq \left(\frac{p}{p-1}\right)^{2p} \mathbb{E}\{|M_{n,m}|^p\},$$

which has the desired result, thanks to Jensen's inequality. \square

What about the $p = 1$ case?

Lemma 1.2 For all $n, m \geq 1$,

$$\mathbb{E}\left\{\max_{\substack{i \leq n \\ j \leq m}} |M_{i,j}|\right\} \leq \left(\frac{e}{e-1}\right)^2 \left[2 + \mathbb{E}\{|M_{n,m}| \ln_+ |M_{n,m}|\}\right].$$

For a proof, see Ch. 1 of MPP.

Finally, we have

Lemma 1.3 For all $n, m \geq 1$ and all $\lambda > 0$,

$$\mathbb{P}\left\{\max_{\substack{i \leq n \\ j \leq m}} |M_{i,j}| \geq \lambda\right\} \leq \frac{1}{\lambda} \left(\frac{e}{e-1}\right) \left[1 + \mathbb{E}\{|M_{n,m}| \ln_+ |M_{n,m}|\}\right].$$

Proof Since $\max_{j \leq m} |M_{i,j}|$ is a submartingale for each $i \geq 1$, by Doob's maximal inequality,

$$\mathbb{P}\left\{\max_{\substack{i \leq n \\ j \leq m}} |M_{i,j}| \geq \lambda\right\} \leq \frac{1}{\lambda} \mathbb{E}\left\{\max_{j \leq m} |M_{n,j}|\right\},$$

and another application of Doob's maximal inequality does the job; cf. (Corollary 3.2, Lecture 3). \square

2 Two-Parameter Convergence Theorems

Continuing with our discussion of the previous section, we now address the question of when, if ever, $\lim M_{n,m}$ exists.

One can imagine many different notions of limits in the present 2-parameter setting. Here are two important ones:

- **Topological limits.** We say that M has a *topological limit* (at infinity) if with probability one, $\lim_{n,m \rightarrow \infty} M_{n,m}$ exists. That is, outside one null set, we can let n and m go to infinity in any way we like and still $M_{n,m}$ converges. This is, indeed, a topological limit, as can be argued simply as follows: Consider $\mathbb{R}^2 / \partial \mathbb{R}^2$, i.e., \mathbb{R}^2 where we identify all points on the axes with each other. Endow it with the relative topology, and then, one-point compactify it to see that $\lim_{n,m \rightarrow \infty}$ means (n, m) converges, in the latter topology, to the added compactification points ∞ , say.
- **Pathwise limits.** In contrast with topological convergence, pathwise convergence is a probabilistic notion. A collection of points $(i_1, j_1), \dots, (i_k, j_k)$ is an *increasing path* if $i_1 \leq \dots \leq i_k$ and $j_1 \leq \dots \leq j_k$. This extends to the case where some of the i_ℓ 's and j_ℓ 's are infinity, as well. With this in mind, we say that our two-parameter orthomartingale M has *pathwise limits* if for every increasing path $(i_1, j_1), (i_2, j_2), \dots$ there exists a null set outside which $\lim_{\ell \rightarrow \infty} M_{i_\ell, j_\ell}$ exists. Both the limit, here, as well as the null set may depend on the increasing path.

Remark 1 Check that any real-valued function $f(n, m)$ has topological limits iff $\lim_n f(n, m)$ (and $\lim_m f(n, m)$, resp.) exists *uniformly* in m (in n , resp.)

Remark 2 If M is non random, the existence of a topological limit at infinity is equivalent to the existence of pathwise limits at infinity (i.e., when the increasing paths go away from the axes in all directions.) For nonrandom processes, this is not true, in general, since there are uncountably many increasing paths.

Theorem 2.1 *The process M always has limits. Moreover, its pathwise limit at infinity is $\mathbb{E}\{Y \mid \mathcal{F}_\infty^1 \vee \mathcal{F}_\infty^2\}$, where $\mathcal{F}_\infty^\ell = \bigvee_{j \geq 1} \mathcal{F}_j^\ell$. Finally, it has a topological limit at infinity if $\mathbb{E}\{|Y| \ln_+ |Y|\} < \infty$.*

Proof The existence of pathwise limits is easy: let $(i_1, j_1), (i_2, j_2), \dots$ denote any increasing path to deduce that $\ell \mapsto M_{i_\ell, j_\ell}$ is a uniformly integrable, one-parameter martingale with respect to the one-parameter filtration $\mathcal{F}_{i_\ell}^1 \vee \mathcal{F}_{j_\ell}^2$. Thus, by Doob's theorem, it converges to $\mathbb{E}\{Y \mid \mathcal{F}_\infty^1 \vee \mathcal{F}_\infty^2\}$, a.s. in and $L^1(\mathbb{P})$.

Next, we suppose $Y \in L \ln_+ L$, and aim to prove topological convergence. Let $\Psi(x) = x \ln_+(x)$ ($x > 0$) and note that

$$\Psi(|x - y|) \leq |\Psi(x) - \Psi(y)|, \quad \forall x, y > 0.$$

This requires a few lines of calculations, but also follows from convexity consideration; cf. Ch. 1 of MPP.

The preceding discussion, together with Lebesgue's dominated convergence theorem, imply that when $Y \in L \ln_+ L$,

$$\lim_{i, j \rightarrow \infty} \mathbb{E}\{\Psi(|M_{i, j} - Y|)\} = 0.$$

In particular, there exists a sequence of constants, $c_1 \leq c_2 \leq \dots$, such that $\lim_{j \rightarrow \infty} c_j = \infty$, and

$$\mathbb{E}\{\Psi(c_j |M_{i, j} - Y|)\} \leq 1.$$

(Why?) Now, $c_j \{M_{i, j} - Y\}$ is also a two-parameter Doob-type ortho-martingale. Thus, we apply Lemma 1.3 to the latter and deduce that for all $\lambda > 0$,

$$\mathbb{P}\left\{\sup_{\substack{i \geq 1 \\ j \geq m}} |M_{i, j} - Y| \geq \lambda\right\} \leq \frac{1}{c_m \lambda} \sup_{j \geq m} \sup_{i \geq 1} \mathbb{E}\{\Psi(c_j |M_{i, j} - Y|)\} \leq \frac{1}{c_m \lambda}.$$

The above goes to zero, as $m \rightarrow \infty$, and this is enough to show that for all $\lambda > 0$,

$$\limsup_{m \rightarrow \infty} \sup_{\substack{i \geq 1 \\ j \geq m}} |M_{i, j} - Y| < \lambda, \quad \text{a.s.}$$

Since the above holds for all $\lambda > 0$, this proves our result. □

3 Further Discussion

In general, we have N one-parameter filtrations, $\mathcal{F}^1, \dots, \mathcal{F}^N$ and say that an N -parameter process is an *ortho-martingale* (with respect to these filtrations) if

- M_t is $\mathcal{F}_{t_1}^1 \vee \cdots \vee \mathcal{F}_{t_N}^N$ -measurable for all $t \in \mathbb{N}^N$;
- for all $t \in \mathbb{N}^N$, $M_t \in L^1(\mathbb{P})$; and
- whenever $t_i \geq s_i$ for all $i = 1, \dots, N$ (written $t \succcurlyeq s$),

$$\mathbb{E}\{M_t \mid \mathcal{F}_{s_1}^1 \vee \cdots \vee \mathcal{F}_{s_N}^N\} = M_s, \quad \text{a.s.}$$

We, then, have the following, whose proof is based on ideas that we have already seen in the simpler context of two-parameter Doob-type ortho-martingales.

Theorem 3.1 (Cairoli) Maximal Inequalities: *Let M be an N -parameter ortho-martingale. Then,*

$$\mathbb{E}\left\{\max_{s \preccurlyeq t} |M_s|^p\right\} \leq \left(\frac{p}{p-1}\right)^{Np} \mathbb{E}\{|M_t|^p\},$$

if $p > 1$. If $p = 1$,

$$\mathbb{E}\left\{\max_{s \preccurlyeq t} |M_s|\right\} \leq \left(\frac{e}{e-1}\right)^N \left[N + \mathbb{E}\{|M_t|[\ln_+ |M_t|]^N\}\right].$$

Convergence Theorems *If M_t is uniformly integrable for all $t \in \mathbb{N}^N$, M has pathwise limits a.s. and in $L^1(\mathbb{P})$. Moreover, its pathwise limit at infinity is $\mathbb{E}\{Y \mid \mathcal{F}_\infty^1 \vee \cdots \vee \mathcal{F}_\infty^N\}$. If $\sup_t \mathbb{E}\{|M_t|[\ln_+ |M_t|]^{N-1}\} < +\infty$, M has topological limits, a.s. and in $L^1(\mathbb{P})$.*

4 Random Walk Examples

4.1 Additive Random Walks

Suppose X and Y are two independent mean 0 random walks, and define the *additive random walk*, S , by

$$S_{n,m} = X_n + Y_m, \quad \forall n, m \geq 1.$$

Let \mathcal{F}^1 and \mathcal{F}^2 denote the filtrations generated by X and Y , respectively. That is,

$$\begin{aligned} \mathcal{F}_i^1 &= \sigma(X_\ell; \ell \leq i), \text{ and} \\ \mathcal{F}_j^2 &= \sigma(Y_\ell; \ell \leq j). \end{aligned}$$

Then, it is easy to see that S is a two-parameter ortho-martingale.

4.2 Multi-Parameter Random Walks

Suppose $X_{i,j}$ are i.i.d. mean zero random variables. The *random walk* associated with the X 's is

$$S_{n,m} = \sum_{\substack{i \leq n \\ j \leq m}} X_{i,j}, \quad \forall n, m \geq 1.$$

Now, we need some filtrations. Let

$$\begin{aligned} \mathcal{F}_i^1 &= \sigma(X_{\ell,j}; \ell \leq i, j \geq 1) \\ \mathcal{F}_j^2 &= \sigma(X_{i,\ell}; i \geq 1, \ell \leq j). \end{aligned}$$

Then, it is not hard, and very instructive, to check that for any $n, m \geq 1$ fixed: $\{S_{i,j}; i \leq n, j \leq m\}$ is a two-parameter Doob-type ortho-martingale with respect to these filtrations. In fact, $\frac{1}{nm}S_{n,m}$ is a reversed “ortho-martingale” (later) with respect to the reversed filtrations

$$\begin{aligned} \mathcal{R}_i^1 &= \sigma(X_{\ell,j}; \ell \geq i, j \geq 1) \\ \mathcal{R}_j^2 &= \sigma(X_{i,\ell}; i \geq 1, \ell \geq j). \end{aligned}$$

This, and the reversed analogue of Theorem 3.1 for ortho-martingales, together, prove the interesting half of *Smythe's law of large numbers*:

Theorem 4.1 (Smythe) *If $X_{1,1} \in L \ln_+ L$,*

$$\lim_{n,m \rightarrow \infty} \frac{S_{n,m}}{nm} = \mathbb{E}\{X_{1,1}\}, \quad \text{a.s. and in } L^1(\mathbb{P}).$$

Moreover, if $(n, m) \mapsto \frac{1}{nm}S_{n,m}$ has topological limits at infinity, then $X_{1,1} \in L \ln_+ L$.

See MPP and Cairoli and Dalang (1996) for complete pedagogical proofs, as well as related results .

Lecture 5

Ortho-Martingales and Intersections of Walks and Brownian Motion

Having motivated orthomartingales somewhat, we proceed in earnest, cookbook style.

1 Ortho-martingales: The Recipe

Ingredients

1. An N -parameter process $\{M_t; t \in \mathbb{N}^N\}$;
2. N one-parameter filtrations $\mathcal{F}^1, \dots, \mathcal{F}^N$, such that for each $i = 1, \dots, N$, $t^{(i)} \mapsto M_t$ is a martingale with respect to \mathcal{F}^i .

For instance, take $N = 2$: then M is a martingale (with respect to \mathcal{F}^1 and \mathcal{F}^2), if

- for all $j, i \mapsto M_{i,j}$ is a martingale for \mathcal{F}^1 ; and
- for all $i, j \mapsto M_{i,j}$ is a martingale for \mathcal{F}^2 .

Example (revisited) Let X^1 and X^2 be two *independent* mean zero random walks, and consider

$$\begin{aligned}\mathcal{F}_i^1 &= \sigma(X_1^1, \dots, X_i^1) \vee \sigma(X_1^2, X_2^2, \dots) \\ \mathcal{F}_j^2 &= \sigma(X_1^1, X_2^1, \dots) \vee \sigma(X_1^2, \dots, X_j^2).\end{aligned}$$

Then, $(n, m) \mapsto X_n^1 + X_m^2$ is a two-parameter ortho-martingale for \mathcal{F}^1 and \mathcal{F}^2 .

SOMETHING TO TRY: Show that $(n, m) \mapsto X_n^1 \cdot X_m^2$ is an ortho-martingale if X^1 and X^2 are independent martingales.

Example (revisited) Suppose $X_{i,j}$ are i.i.d. mean zero random variables and define the 2-parameter random walk S by

$$S_{n,m} = \sum_{i \leq n} \sum_{j \leq m} X_{i,j}.$$

Then, S is a 2-parameter ortho-martingale with respect to

$$\begin{aligned} \mathcal{F}_i^1 &= \sigma(X_{\ell,j} : \ell \leq i, j \geq 1) \\ \mathcal{F}_j^2 &= \sigma(X_{i,\ell} : i \geq 1, \ell \leq j). \end{aligned}$$

Henceforth, when $s, t \in \mathbb{R}^N$ and we write $s \preceq t$, we mean $s^{(i)} \leq t^{(i)}$ for each coordinate $i = 1, \dots, N$.

SOMETHING TO TRY: Check that when M is an ortho-martingale and when $s \preceq t$ are both in \mathbb{N}^N , then $\mathbb{E}\{|M_s|^p\} \leq \mathbb{E}\{|M_t|^p\}$ for $p > 1$. (Hint: Jensen's inequality.)

Now, we come to our maximal inequalities; they are proved working one parameter at a time, just as the Doob-type examples that were worked out earlier.

Theorem 1.1 (Cairol's Inequalities) *If M is an N -parameter ortho-martingale indexed by \mathbb{N}^N say. Then,*

$$\begin{aligned} \mathbb{E}\{\max_{s \preceq t} |M_s|^p\} &\leq C_{p,N} \mathbb{E}\{|M_t|^p\} \\ \mathbb{E}\{\max_{s \preceq t} |M_s|\} &\leq C_{1,N} \left[N + \mathbb{E}\{|M_t| \ln_+^N |M_t|\} \right] \\ \mathbb{P}\{\max_{s \preceq t} |M_s| \geq \lambda\} &\leq \frac{C_{1,N-1}}{\lambda} \left[(N-1) + \mathbb{E}\{|M_t| \ln^{N-1} |M_t|\} \right], \quad \forall \lambda > 0, \end{aligned}$$

where

$$C_{p,N} = \left(\frac{p}{p-1} \right)^{Np} \quad \text{and} \quad C_{1,N} = \left(\frac{e}{e-1} \right)^N.$$

There is a topological convergence theorem, as well. It is proved by working one parameter at a time, as in our example from last week, viz.,

Theorem 1.2 *If M is an ortho-martingale, and if $\sup_t \mathbb{E}\{|M_t| \ln_+^{N-1} |M_t|\} < +\infty$, then the topological limit $\lim_{t \rightarrow \infty} M_t$ exists, a.s. and in $L^1(\mathbb{P})$.*

The analogue of pathwise limits, at this level of generality, is tricky. Instead, we will introduce sectorial limits when $N = 2$. The remaining details can be found in MPP.

Henceforth, $N = 2$, in this section.

The two **sectorial limits**, when they exist, of a two-variable function f are defined by the ordered limiting operations: $\lim_{n \rightarrow \infty} \lim_{m \rightarrow \infty} f(n, m)$ and $\lim_{m \rightarrow \infty} \lim_{n \rightarrow \infty} f(n, m)$. In general, even when they exist, they need not be equal. When both sectorial limits exist and when they both agree, we write $\lim_{n, m \rightarrow \infty} f(n, m)$ for their common value, and call it **the sectorial limit** of f .

Theorem 1.3 *If M is a uniformly integrable ortho-martingale, $\lim_{t \rightsquigarrow \infty} M_t$ exists a.s. and in $L^2(\mathbb{P})$.*

Proof for $N = 2$ By uniform integrability, $M_{i,\infty} = \lim_{j \rightarrow \infty} M_{i,j}$ exists a.s. and in $L^1(\mathbb{P})$. We now make two claims:

Claim 1 $i \mapsto M_{i,\infty}$ is uniformly integrable.

Claim 2 $i \mapsto M_{i,\infty}$ is an \mathcal{F}^1 -submartingale.

To prove Claim 1, note that by Fatou's lemma, for all $\lambda > 0$,

$$\mathbb{E}\{|M_{i,\infty}|; |M_{i,\infty}| \geq \lambda\} \leq \liminf_{j \rightarrow \infty} \mathbb{E}\{|M_{i,j}|; |M_{i,j}| \geq \lambda\} \leq \sup_j \mathbb{E}\{|M_{i,j}|; |M_{i,j}| \geq \lambda\},$$

which goes to 0 as $\lambda \rightarrow \infty$, uniformly in i . Claim 2 is proved similarly. Indeed,

$$\mathbb{E}\{M_{i+1,\infty} | \mathcal{F}_i^1\} \leq \liminf_{j \rightarrow \infty} \mathbb{E}\{M_{i+1,j} | \mathcal{F}_i^1\} = \lim_{j \rightarrow \infty} M_{i,j} = M_{i,\infty}.$$

In particular, $M_{\infty,\infty}^1 = \lim_{i \rightarrow \infty} M_{i,\infty}$ exists. Of course, $M_{\infty,\infty}^1 = \lim_i \lim_j M_{i,j}$ is the first sectorial limit, which we now know exists. Similarly, $M_{\infty,j} = \lim_i M_{i,j}$ exists a.s. and in $L^1(\mathbb{P})$, which leads to the existence of the second sectorial limit, $M_{\infty,\infty}^2 = \lim_j \lim_i M_{i,j}$.

Our job, now, is to check that $M_{\infty,\infty}^1 = M_{\infty,\infty}^2$. But for any $k \geq 1$,

$$\begin{aligned} \mathbb{E}\{M_{\infty,\infty}^1 \mid \mathcal{F}_k^1\} &\leq \liminf_{i \rightarrow \infty} \liminf_{j \rightarrow \infty} \mathbb{E}\{M_{i,j} \mid \mathcal{F}_k^1\} \\ &= \lim_{i \rightarrow \infty} \lim_{j \rightarrow \infty} M_{k,j} \quad (\text{by the ortho-martingale property}) \\ &= M_{k,\infty}. \end{aligned}$$

Let $k \rightarrow \infty$ to see that

$$\mathbb{E}\{M_{\infty,\infty}^1 \mid \bigvee_k \mathcal{F}_k^1\} \leq M_{\infty,\infty}^2.$$

But $M_{\infty,\infty}^1$ is always $\bigvee_k \mathcal{F}_k^1$ -measurable. So, $M_{\infty,\infty}^1 \leq M_{\infty,\infty}^2$, which, by symmetry, shows that the two sectorial limits agree, a.s. \square

2 Applications to Intersections of Simple Walks

Our next goal is to apply the maximal inequality for ortho-martingales to a result about the intersections of independent Brownian motions. To illustrate some of the moment estimates, we first turn to a discrete version, which is simple to work out, although it does not really need ortho-martingale theory.

Let X^1 and X^2 be two simple symmetric random walks in \mathbb{Z}^d , and consider

$$J_{n,m} = \sum_{i \leq n} \sum_{j \leq m} \mathbf{1}\{X_i^1 = X_j^2\},$$

which is nothing other than the number of times the trajectories of the random functions X^1 and X^2 intersect.

How big is $J_{n,n}$? Its mean is easy to estimate. Recall, from the classical local limit theorem, that if S_n is the simple walk in \mathbb{Z}^d ,

$$\begin{aligned} \mathbb{P}\{S_{2n} = 0\} &\sim C_d (2n)^{-\frac{d}{2}}, \quad \text{as } n \rightarrow \infty, \text{ and} \\ \sup_{a \in \mathbb{Z}^d} \mathbb{P}\{S_k = a\} &\leq C_d k^{-\frac{d}{2}}, \quad \forall k \geq 1. \end{aligned} \tag{2.1}$$

Of course, $\mathbb{P}\{S_{2n+1} = 0\} = 0$, for all n , and $a_n \sim b_n$ means that $\lim_n a_n \div b_n = 1$.

Now, we can estimate $\mathbb{E}\{J_{n,m}\}$, since $X_i^1 + X_j^2$ has the same distribution as S_{i+j} . Thus,

$$\mathbb{E}\{J_{n,n}\} \sim C'_d \sum_{i \leq n} \sum_{j \leq n} (i+j)^{-\frac{d}{2}},$$

as $n \rightarrow \infty$. The above has the same asymptotics as $\int_1^n \int_1^n (x+y)^{-\frac{d}{2}} dx dy$ which can be computed explicitly. However, it is more elegant to observe that

$$\begin{aligned} \int_{[1,n]^2} (x+y)^{-\frac{d}{2}} dx dy &\asymp \int_{\substack{z \in \mathbb{R}^2: \\ 1 \leq \|z\| \leq cn}} \|z\|^{-\frac{d}{2}} dz \\ &= C \int_1^{cn} w^{-\frac{d}{2}+1} dw, \end{aligned}$$

where \asymp means ‘has the same rough asymptotics as’. More precisely, $f_n \asymp g_n$ means $0 < \liminf_n f_n/g_n \leq \limsup_n f_n/g_n < \infty$. The last line follows from calculating in polar coordinates. As usual, constants c, C_d , etc. are immaterial to our discussion. In fact, a little more care can be used to show that

$$\mathbb{E}\{J_{n,n}\} \sim C_d \times \begin{cases} 1, & \text{if } d > 4 \\ \ln n, & \text{if } d = 4 \\ n^{\frac{1}{2}}, & \text{if } d = 3. \\ n, & \text{if } d = 2 \\ n^{\frac{3}{2}}, & \text{if } d = 1 \end{cases} \quad (2.2)$$

We now make the following claim:

Lemma 2.1 *In all dimensions, there exists C_d such that for all $n \geq 1$,*

$$\mathbb{E}\{J_{n,n}^2\} \leq C_d |\mathbb{E}\{J_{n,n}\}|^2.$$

Admitting Lemma 2.1 for the moment, we combine it with (2.2) and the Paley-Zygmund Lemma (Lemma 1.2, Lecture 1) to get:

$$\inf_{n \geq 1} \mathbb{P}\{J_{n,n} \geq \frac{1}{2} \mathbb{E}(J_{n,n})\} > 0, \quad \forall d \geq 1.$$

On the other hand, Eq. (2.2) also shows us that $\mathbb{E}\{J_{n,n}\} \rightarrow \infty$ iff $n \leq 4$. Thus, we have obtained half of the following theorem, due to Dvoretzky, Erdős and Kakutani:

Theorem 2.2 *Two independent simple walk paths in \mathbb{Z}^d intersect infinitely often iff $d \leq 4$.*

As I mentioned earlier, modulo proving Lemma 2.1, we have shown half of this theorem; namely, that $d \leq 4$ implies infinite intersections ‘with positive probability’, which implies ‘with probability one’, by

the Hewitt–Savage 0-1 law. On the other hand, if $d > 5$, Eq. (2.1) shows us that $\limsup_n \mathbb{E}\{J_{n,n}\} < +\infty$. From this, we gather that $J_{\infty,\infty} < +\infty$, a.s. and we are done. Thus, Theorem 2.2 will follow from our

Proof of Lemma 2.1 Just expand the square, and use symmetry considerations to obtain

$$\begin{aligned} \mathbb{E}\{J_{n,n}^2\} &= \sum_{i=1}^n \sum_{j=1}^n \sum_{i'=1}^n \sum_{j'=1}^n \mathbb{P}\{X_i^1 = X_j^2, X_{i'}^1 = X_{j'}^2\} \\ &\leq 2 \sum_{1 \leq i < i' \leq n} \sum_{1 \leq j < j' \leq n} \mathbb{P}\{X_i^1 = X_j^2, X_{i'}^1 = X_{j'}^2\} + \\ &\quad + 2 \sum_{1 \leq i \leq i' \leq n} \sum_{1 \leq j' \leq j \leq n} \mathbb{P}\{X_i^1 = X_j^2, X_{i'}^1 = X_{j'}^2\}. \end{aligned}$$

In expanding the above, it helps to recognize that there are four cases to consider:

1. $i < i'$ and $j < j'$;
2. $i < i'$ but $j > j'$;
3. $i = i'$ but $j < j'$; and finally
4. $i = i'$ and $j = j'$.

Thus, we can write

$$\mathbb{E}\{J_{n,n}^2\} = 2(T_1 + T_2 + T_3 + T_4)$$

where

$$\begin{aligned} T_1 &= \sum_{1 \leq i < i' \leq n} \sum_{1 \leq j < j' \leq n} \mathbb{P}\{X_i^1 = X_j^2, X_{i'}^1 = X_{j'}^2\} \\ T_2 &= \sum_{1 \leq i < i' \leq n} \sum_{1 \leq j' < j \leq n} \mathbb{P}\{X_i^1 = X_j^2, X_{i'}^1 = X_{j'}^2\} \\ T_3 &= \sum_{1 \leq i \leq n} \sum_{1 \leq j < j' \leq n} \mathbb{P}\{X_i^1 = X_j^2 = X_{j'}^2\} \\ T_4 &= \sum_{1 \leq i \leq n} \sum_{1 \leq j \leq n} \mathbb{P}\{X_i^1 = X_j^2\}. \end{aligned}$$

Of course,

$$T_4 = \mathbb{E}\{J_{n,n}\}. \tag{2.3}$$

Furthermore,

$$\begin{aligned}
T_3 &= \sum_{1 \leq i \leq n} \sum_{1 \leq j < j' \leq n} \mathbb{P}\{X_i^1 = X_j^2\} \cdot \mathbb{P}\{X_{j'-j}^2 = 0\} \\
&\leq \mathbb{E}\{J_{n,n}\} \cdot \sum_{j=1}^n \mathbb{P}\{X_j^2 = 0\}.
\end{aligned} \tag{2.4}$$

For T_2 , we note that if $i < i'$ and $j' < j$,

$$\mathbb{P}\{X_i^1 = X_j^2, X_{i'}^1 = X_{j'}^2\} = \mathbb{P}\{X_i^1 = X_{j'}^2 + \xi_1, X_i^1 + \xi_2 = X_{j'}^2\}$$

where

$$\begin{aligned}
\xi_1 &= X_j^2 - X_{j'}^2, \text{ and} \\
\xi_2 &= X_{i'}^1 - X_i^1.
\end{aligned}$$

Thus, by the independence of (ξ_1, ξ_2) from $(X_i^1, X_{j'}^2)$,

$$\begin{aligned}
\mathbb{P}\{X_i^1 = X_j^2, X_{i'}^1 = X_{j'}^2\} &= \mathbb{P}\{X_i^1 = X_{j'}^2 + \xi_1, \xi_1 + \xi_2 = 0\} \\
&\leq \sup_{a \in \mathbb{Z}^d} \mathbb{P}\{X_i^1 = X_{j'}^2 + a\} \cdot \mathbb{P}\{\xi_1 + \xi_2 = 0\} \\
&= \sup_{a \in \mathbb{Z}^d} \mathbb{P}\{X_i^1 = X_{j'}^2 + a\} \cdot \mathbb{P}\{X_{i'-i}^1 + X_{j-j'}^2 = 0\} \\
&= \sup_{a \in \mathbb{Z}^d} \mathbb{P}\{X_i^1 = X_{j'}^2 + a\} \cdot \mathbb{P}\{X_{i'-i}^1 = X_{j-j'}^2\},
\end{aligned}$$

Now, by another appeal to the local limit theorem, $\mathbb{P}\{X_i^1 = X_{j'}^2 + a\} \leq C_d(i+j)^{-\frac{d}{2}} \leq C'_d \mathbb{P}\{X_i^1 = X_{j'}^2\}$, at least when $i+j$ is even; cf. (2.1). Therefore, by reshuffling the labels,

$$\begin{aligned}
T_2 &\leq \sum_{1 \leq i < i' \leq n} \sum_{1 \leq j' < j \leq n} \mathbb{P}\{X_i^1 = X_j^2, X_{i'}^1 = X_{j'}^2\} \\
&\leq C_d \sum_{1 \leq i < i' \leq n} \sum_{1 \leq j' < j \leq n} \mathbb{P}\{X_i^1 = X_j^2\} \cdot \mathbb{P}\{X_{i'-i}^1 = X_{j-j'}^2\} \\
&\leq C_d |\mathbb{E}\{J_{n,n}\}|^2.
\end{aligned} \tag{2.5}$$

Finally,

$$\begin{aligned}
T_1 &= \sum_{1 \leq i < i' \leq n} \sum_{1 \leq j < j' \leq n} \mathbb{P}\{X_i^1 = X_j^2\} \mathbb{P}\{X_{i'-i}^1 = X_{j'-j}^2\} \\
&\leq \frac{1}{2} [\mathbb{E}\{J_{n,n}\}]^2.
\end{aligned}$$

Combining this with Eq.'s (2.3), (2.4) and (2.5), we obtain

$$\mathbb{E}\{J_{n,n}^2\} \leq C_d [\mathbb{E}\{J_{n,n}\}]^2 + 2\mathbb{E}\{J_{n,n}\} \cdot \left[1 + 2 \sum_{j=1}^n \mathbb{P}\{X_j^1 = 0\}\right].$$

We use the local limit theorem, once more, to see that

$$\sum_{j=1}^n \mathbb{P}\{X_j^1 = 0\} \leq C_d \times \begin{cases} 1, & \text{if } d \geq 3 \\ \ln n, & \text{if } d = 2 \\ n^{\frac{1}{2}}, & \text{if } d = 1 \end{cases}. \quad (2.6)$$

The lemma follows from this. □

You should compare Theorem 2.2 with the recurrence theorem for Markov chains, specialized to this setting; cf. Theorem 1.1 of Lecture 1. Indeed, we note that by the latter theorem, and by the computations that lead to Eq. (2.6),

Theorem 2.3 *The simple walk in \mathbb{Z}^d is recurrent iff $d \leq 2$.*

Thus, it is the case that some transient walks intersect infinitely many times.

We could have studied this earlier on, since there are no multiparameter martingales needed. However, the continuous analogue of this theorem forces us to reconsider such a remark. Indeed, we shall prove the following theorem of Dvoretzky et al. next time, using the Cairoli-Walsh theory.

Theorem 2.4 *Let X and Y denote 2 independent d -dimensional Brownian motions, both starting at 0. Then, $X(\mathbb{R}_+) \cap Y(\mathbb{R}_+) = \{0\}$ iff $d \geq 4$.*

In the course of our proof of Lévy theorem, we needed a probability estimate (cf. Theorem 1.2 of Lecture 2) that has the following two-parameter analogue, due to Aizenmann and Simon.

Theorem 2.5 (Aizenmann and Simon) *There exists C_1 and C_2 , such that for all $\varepsilon \in (0, 1)$,*

$$C_1 \kappa(\varepsilon) \leq \mathbb{P}\left\{ \inf_{1 \leq s, t \leq 2} |X(s) - Y(t)| \leq \varepsilon \right\} \leq C_2 \kappa(\varepsilon),$$

where

$$\kappa(\varepsilon) = \begin{cases} \varepsilon^{d-4}, & \text{if } d \geq 5 \\ [\ln(1/\varepsilon)]^{-1}, & \text{if } d = 4. \\ 1, & \text{if } d \leq 3 \end{cases}$$

Lecture 6

Intersections of Brownian Motion, Multiparameter Martingales

Before proceeding, I wish to state two open problems; one of them was mentioned in Lecture 3.

Open Problem 1. Suppose $X_{i,j}$ are i.i.d. centered random variables for $i, j \geq 1$, and consider the 2-parameter random walk

$$S_{n,m} = \sum_{i \leq n} \sum_{j \leq m} X_{i,j}, \quad \forall n, m \geq 1.$$

By standardizing them, we can assume, without loss of generality, that $\mathbb{E}[X_{i,j}] = 0$ and $\mathbb{E}[X_{i,j}^2] = 1$. Then, it is possible to show that the following law of the iterated logarithm holds:

$$\limsup_{n,m \rightarrow \infty} \frac{S_{n,m}}{\sqrt{4nm \log \log nm}} = 1,$$

at least as long as $X_{1,1} \in L^2 \ln L$; cf. Ch. 4 of MPP. On the other hand, $X_{1,1} \in L^2$ is clearly not enough. In fact, $X_{1,1} \in \frac{L^2 \ln L}{\ln \ln L}$ is necessary for

$$\limsup_{n,m \rightarrow \infty} \frac{|S_{n,m}|}{\sqrt{nm \log \log nm}} < +\infty,$$

(*loc. cit.*) Is there a necessary and sufficient moment-type condition for the LIL?

Open Problem 2. By working harder, one can extend Theorem 2.2 of Lecture 5 as follows: if S^1 and S^2 are two independent, *symmetric*, and transient random walks on \mathbb{Z}^d ,

$$\#\{S^1(\mathbb{Z}_+) \cap S^2(\mathbb{Z}_+)\} = +\infty \iff \sum_{i,j \geq 1} \mathbb{P}\{S_i^1 = S_j^2\} = +\infty.$$

See Ch. 3 of MPP. Can symmetry be dropped? Are there *any* nonsymmetric walks that can be analyzed?

1 Proof of Theorem 2.2

We start, as before, by considering the first two moments of

$$J_\varepsilon = \int_1^2 \int_1^2 \mathbf{1}_{\{|X_s - Y_t| \leq \varepsilon\}} ds dt.$$

Note that $X_s - Y_t$ is a normal vector with mean zero and variance matrix $\sqrt{s+t}$ times the identity matrix. Note also that $1 \leq \sqrt{s+t} \leq 2$, for all $s, t \in [1, 2]$. Thus, by Lemma 1.1 of Lecture 3, we can find C_1 and C_2 such that for all $\varepsilon \in (0, 1)$,

$$C_1 \varepsilon^d \leq \mathbb{E}\{J_\varepsilon\} \leq C_2 \varepsilon^d.$$

Furthermore,

$$\begin{aligned} \mathbb{E}\{J_\varepsilon^2\} &= \iint_{[1,2]^2} \iint_{[1,2]^2} \mathbb{P}\{|X_s - Y_t| \leq \varepsilon, |X_{s'} - Y_{t'}| \leq \varepsilon\} ds dt ds' dt' \\ &= 2(T_1 + T_2), \end{aligned}$$

where

$$\begin{aligned} T_1 &= \iint_{1 \leq s \leq s' \leq 2} \iint_{1 \leq t \leq t' \leq 2} \mathbb{P}\{|X_s - Y_t| \leq \varepsilon, |X_{s'} - Y_{t'}| \leq \varepsilon\} ds dt ds' dt', \text{ and} \\ T_2 &= \iint_{1 \leq s' \leq s \leq 2} \iint_{1 \leq t \leq t' \leq 2} \mathbb{P}\{|X_s - Y_t| \leq \varepsilon, |X_{s'} - Y_{t'}| \leq \varepsilon\} ds dt ds' dt'. \end{aligned}$$

Clearly,

$$\begin{aligned} T_1 &\leq \iint_{1 \leq s \leq s' \leq 2} \iint_{1 \leq t \leq t' \leq 2} \mathbb{P}\{|X_s - Y_t| \leq \varepsilon\} \cdot \mathbb{P}\{|X_{s'-s} - Y_{t'-t}| \leq 2\varepsilon\} ds dt ds' dt' \\ &\leq C_3 \iint_{1 \leq s \leq s' \leq 2} \iint_{1 \leq t \leq t' \leq 2} \varepsilon^d \cdot \left(\frac{\varepsilon}{\sqrt{|s' - s| + |t' - t|}} \wedge 1 \right)^d ds dt ds' dt'. \end{aligned}$$

We have used Lemma 1.1 of Lecture 3 once more. Now, this shows that

$$\begin{aligned}
T_1 &\leq C'_3 \varepsilon^d \int_{\substack{x \in \mathbb{R}^2: \\ \|x\| \leq c}} \left(\frac{\varepsilon}{\|x\|^{\frac{1}{2}} \wedge 1} \right)^d dx \\
&= C''_3 \varepsilon^d \int_0^c \left(\frac{\varepsilon}{r^{\frac{1}{2}}} \wedge 1 \right)^d r dr \\
&\leq C''_3 \varepsilon^d \times \begin{cases} \varepsilon^d, & \text{if } d \leq 3 \\ \varepsilon^2 \ln_+(\frac{1}{\varepsilon}), & \text{if } d = 4. \\ \varepsilon^2, & \text{if } d \geq 5 \end{cases}
\end{aligned}$$

We will obtain the same estimates for T_2 :

$$\begin{aligned}
T_2 &= \iint_{1 \leq s' \leq s \leq 2} \iint_{1 \leq t \leq t' \leq 2} \mathbb{P}\{|X_{s'} + \xi_1 - Y_t| \leq \varepsilon, |X_{s'} + \xi_2 - Y_t| \leq \varepsilon\} ds dt ds' dt' \\
&\leq \iint_{1 \leq s' \leq s \leq 2} \iint_{1 \leq t \leq t' \leq 2} \mathbb{P}\{|X_{s'} + \xi_1 - Y_t| \leq \varepsilon, |\xi_1 + \xi_2| \leq 2\varepsilon\} ds dt ds' dt',
\end{aligned}$$

where

$$\xi_1 = (X_s - X_{s'}) \quad \text{and} \quad \xi_2 = (Y_{t'} - Y_t).$$

Note that $X_s - Y_t$ is independent of (ξ_1, ξ_2) . Thus,

$$\begin{aligned}
T_2 &\leq \iint_{1 \leq s' \leq s \leq 2} \iint_{1 \leq t \leq t' \leq 2} \sup_{a \in \mathbb{R}^d} \mathbb{P}\{|X_{s'} + a - Y_t| \leq \varepsilon\} \mathbb{P}\{|\xi_1 + \xi_2| \leq 2\varepsilon\} ds dt ds' dt' \\
&\leq C_4 \varepsilon^d \times \begin{cases} \varepsilon^d, & \text{if } d \leq 3 \\ \varepsilon^2 \ln_+(\frac{1}{\varepsilon}), & \text{if } d = 4, \\ \varepsilon^2, & \text{if } d \geq 5 \end{cases}
\end{aligned}$$

by yet another application of Lemma 1.1 of Lecture 3. Combining what we have, we obtain

$$\mathbb{E}\{J_\varepsilon\} \asymp \varepsilon^d \quad \text{and} \quad \mathbb{E}\{J_\varepsilon^2\} \leq C \varepsilon^d \times \begin{cases} \varepsilon^d, & \text{if } d \leq 3 \\ \varepsilon^4 \ln_+(\frac{1}{\varepsilon}), & \text{if } d = 4. \\ \varepsilon^4, & \text{if } d \geq 5 \end{cases} \quad (1.1)$$

This is good enough to prove that

$$\mathbb{P}\left\{ \inf_{1 \leq s, t \leq 2} |X_s - Y_t| \leq \varepsilon \right\} \geq C \kappa(\varepsilon),$$

which is the desired lower bound. We now show this is sharp.

Consider the two 1-parameter filtrations

$$\mathcal{F}_s^1 = \sigma(X_r; r \leq s) \vee \sigma(Y_u; u \geq 0) \quad \text{and} \quad \mathcal{F}_t^2 = \sigma(Y_r; r \leq t) \vee \sigma(X_u; u \geq 0).$$

We now estimate

$$\begin{aligned} M_{s,t} &= \mathbb{E}\{J_\varepsilon \mid \mathcal{F}_s^1 \cap \mathcal{F}_t^2\} \geq \int_s^2 \int_t^2 \mathbb{P}\{|X_u - Y_v| \leq \varepsilon \mid \mathcal{F}_s^1 \cap \mathcal{F}_t^2\} dv du \cdot \mathbf{1}_{\{|X_s - Y_t| \leq \frac{\varepsilon}{2}\}} \\ &\geq \int_0^{2-s} \int_0^{2-t} \mathbb{P}\{|X_{u-s} - Y_{v-t}| \leq \frac{\varepsilon}{2}\} dv du \cdot \mathbf{1}_{\{|X_s - Y_t| \leq \frac{\varepsilon}{2}\}}. \end{aligned}$$

Thus, as long as $1 \leq s, t \leq \frac{3}{2}$,

$$M_{s,t} \geq \int_0^{\frac{1}{2}} \int_0^{\frac{1}{2}} \mathbb{P}\{|X_{u-s} - Y_{v-t}| \leq \frac{\varepsilon}{2}\} dv du \cdot \mathbf{1}_{\{|X_s - Y_t| \leq \frac{\varepsilon}{2}\}}.$$

Now, the above double integral is bounded below by a constant multiple of

$$\int_{\substack{z \in \mathbb{R}^2: \\ \|z\| \leq c}} \left(\frac{\varepsilon}{\|z\|^{\frac{1}{2}}} \wedge 1 \right)^d dz \geq C_d \frac{\mathbb{E}\{J_\varepsilon^2\}}{\varepsilon^d}.$$

So,

$$\mathbb{E}\left\{ \sup_{1 \leq s, t \leq \frac{3}{2}} M_{s,t}^2 \right\} \geq C_d \left[\frac{\mathbb{E}\{J_\varepsilon^2\}}{\varepsilon^d} \right]^2 \mathbb{P}\left\{ \inf_{1 \leq s, t \leq \frac{3}{2}} |X_s - Y_t| \leq \frac{\varepsilon}{2} \right\}.$$

On the other hand, Cairoli's maximal inequality assures us that $\mathbb{E}\{\sup_{s,t} M_{s,t}^2\} \leq 16\mathbb{E}\{J_\varepsilon^2\}$, which, together with Eq. (1.1) yields

$$\mathbb{P}\left\{ \inf_{1 \leq s, t \leq \frac{3}{2}} |X_s - Y_t| \leq \frac{\varepsilon}{2} \right\} \leq C\kappa(\varepsilon).$$

By a change of scale, the infimum can be taken over $s, t \in [1, 2]$ at little cost, and we are done. \square

2 Multiparameter Martingales

Consider an N -parameter process $\{M_t; t \in \mathbb{N}^N\}$, and an N -parameter sequence of sigma-fields $\mathcal{F} = \{\mathcal{F}_t; t \in \mathbb{N}^N\}$.

We say that \mathcal{F} is a *filtration*, if

$$s \preccurlyeq t \implies \mathcal{F}_s \subseteq \mathcal{F}_t.$$

Perhaps it is best to recall that $s \preccurlyeq t$ if and only if $s_i \leq t_i$ for all $i = 1, \dots, N$. We say that M is *adapted* to the N -parameter filtration \mathcal{F} if for all $t \in \mathbb{N}^N$, M_t is \mathcal{F}_t -measurable. Finally, M is a *martingale* with respect to the filtration \mathcal{F} if (i) M is adapted to \mathcal{F} ; (ii) $M_t \in L^1(\mathbb{P})$ for all $t \in \mathbb{N}^N$; and (iii) almost surely,

$$s \preccurlyeq t \implies \mathbb{E}\{M_t | \mathcal{F}_s\} = M_s.$$

In general, multiparameter martingales have not much structure to speak of. However, they do, if the filtration in question has some conditional independence properties.

2.1 Commutation

We say that the N -parameter filtration \mathcal{F} is *commuting*, if for all $s, t \in \mathbb{N}^N$, and for all \mathcal{F}_t -measurable bounded random variables Y ,

$$\mathbb{E}\{Y | \mathcal{F}_s\} = \mathbb{E}\{Y | \mathcal{F}_{s \wedge t}\},$$

where $s \wedge t$ is the point in \mathbb{N}^N whose i th coordinate is $s_i \wedge t_i$. The typical N -parameter filtration does *not* commute!

Theorem 2.1 *If \mathcal{F} is a commuting N -parameter filtration, for all random variables $Z \in L^1(\mathbb{P})$, and for all $t \in \mathbb{N}^N$,*

$$\mathbb{E}\{Z | \mathcal{F}_t\} = \mathbb{E}\left[\dots \mathbb{E}\left(\mathbb{E}\{Z | \mathcal{F}_{t_1}^1\} \mid \mathcal{F}_{t_2}^2 \right) \dots \mid \mathcal{F}_{t_N}^N \right].$$

In particular, we shall see, as a consequence, that any martingale with respect to a commuting \mathcal{F} is an ortho-martingale! This, in turn, implies maximal inequalities, convergence theorems, etc.

Proof It suffices to prove this for bounded Z (why?). We will do this when $N = 2$.

By Doob's 1-parameter martingale convergence theorem, and by Lebesgue's dominated convergence theorem,

$$\begin{aligned} \mathbb{E}\left[\mathbb{E}\{Z | \mathcal{F}_i^1\} \mid \mathcal{F}_j^2 \right] &= \mathbb{E}\left[\lim_{k \rightarrow \infty} \mathbb{E}\{Z | \mathcal{F}_{i,k}\} \mid \mathcal{F}_j^2 \right] \\ &= \lim_{k \rightarrow \infty} \mathbb{E}\left[\mathbb{E}\{Z | \mathcal{F}_{i,k}\} \mid \mathcal{F}_j^2 \right] \\ &= \lim_{k \rightarrow \infty} \lim_{\ell \rightarrow \infty} \mathbb{E}\left[\mathbb{E}\{Z | \mathcal{F}_{i,k}\} \mid \mathcal{F}_{\ell,j} \right], \end{aligned}$$

where all of the convergences are taking place in $L^1(\mathbb{P})$. Recall that \mathcal{F} is commuting and $Y = \mathbb{E}\{Z | \mathcal{F}_{i,k}\}$ is $\mathcal{F}_{i,k}$ -measurable and bounded. This implies that for any other $\ell, j \geq 1$, a.s.,

$$\begin{aligned} \mathbb{E}\{Y | \mathcal{F}_{\ell,j}\} &= \mathbb{E}\{Y | \mathcal{F}_{(i,k) \wedge (\ell,j)}\} \\ &= \mathbb{E}\{Y | \mathcal{F}_{i \wedge \ell, k \wedge j}\} \\ &= \mathbb{E}\{Z | \mathcal{F}_{i \wedge \ell, k \wedge j}\}. \end{aligned}$$

Thus, for every $i, j \geq 0$ and for all bounded random variables Z ,

$$\mathbb{E}[\mathbb{E}\{Z | \mathcal{F}_i^1\} | \mathcal{F}_j^2] = \lim_{k \rightarrow \infty} \lim_{\ell \rightarrow \infty} \mathbb{E}[Z | \mathcal{F}_{i \wedge \ell, k \wedge j}], \quad \text{a.s.}$$

This is clearly equal to $\mathbb{E}\{Z | \mathcal{F}_{i,j}\}$, and the result follows. \square

The notion of commutation is equivalent to conditional independence. Recall that two sigma-fields \mathcal{G}_1 and \mathcal{G}_2 are *conditionally independent, given \mathcal{G}* , if for all bounded \mathcal{G}_i -measurable random variables Y_i ($i = 1, 2$),

$$\mathbb{E}\{Y_1 Y_2 | \mathcal{G}\} = \mathbb{E}\{Y_1 | \mathcal{G}\} \cdot \mathbb{E}\{Y_2 | \mathcal{G}\}.$$

Theorem 2.2 For a given filtration $\mathcal{F} = (\mathcal{F}_t; t \in \mathbb{N}_0^N)$, the following are equivalent:

- (i) \mathcal{F} is commuting; and
- (ii) for all $s, t \in \mathbb{N}_0^N$, \mathcal{F}_t and \mathcal{F}_s are conditionally independent, given $\mathcal{F}_{s \wedge t}$.

Proof Suppose for all $t \in \mathbb{N}_0^N$, Y_t is a bounded \mathcal{F}_t -measurable random variable. Then,

$$\mathbb{E}[Y_t Y_s | \mathcal{F}_{t \wedge s}] = \mathbb{E}[Y_t \mathbb{E}\{Y_s | \mathcal{F}_t\} | \mathcal{F}_{t \wedge s}] = \mathbb{E}[Y_t \mathbb{E}\{Y_s | \mathcal{F}_{t \wedge s}\} | \mathcal{F}_{t \wedge s}], \quad \text{a.s.}$$

Thus, (i) \Rightarrow (ii). Conversely, supposing that (ii) holds,

$$\begin{aligned} \mathbb{E}[Y_t Y_s] &= \mathbb{E}\left[\mathbb{E}\{Y_t Y_s | \mathcal{F}_{t \wedge s}\}\right] \\ &= \mathbb{E}\left[\mathbb{E}\{Y_t | \mathcal{F}_{t \wedge s}\} \cdot \mathbb{E}\{Y_s | \mathcal{F}_{t \wedge s}\}\right] \\ &= \mathbb{E}\left[\mathbb{E}\{Y_t | \mathcal{F}_{t \wedge s}\} \cdot Y_s\right]. \end{aligned}$$

Since $\mathcal{F}_{t \wedge s} \subset \mathcal{F}_s$ and the above holds for all bounded, \mathcal{F}_s -measurable random variables Y_s , $\mathbb{E}[Y_t | \mathcal{F}_{t \wedge s}] = \mathbb{E}[Y_t | \mathcal{F}_s]$, a.s.. This shows that (ii) implies (i) and hence, (ii) is equivalent to (i). \square

Example If $X_{i,j}$ are independent random variables ($i, j \geq 1$), define $\mathcal{F}_{n,m} = \sigma(X_{i,j}; i \leq n, j \leq m)$ to see that \mathcal{F} is commuting. In particular, if we also knew that the X 's have mean zero, $S_{n,m} = \sum_{i \leq n} \sum_{j \leq m} X_{i,j}$ is a martingale, as well as an ortho-martingale. What about additive walks? \square

2.2 Back to Martingales

We conclude our discussion of martingales, under commutation, by linking them to ortho-martingales.

Theorem 2.3 *Suppose \mathcal{F} is a commuting N -parameter filtration. Then, M is a martingale with respect to \mathcal{F} if and only if it is an ortho-martingale with respect to the marginal filtrations of \mathcal{F} .*

Proof Once again, we only really need to do this for $N = 2$. First, suppose M is an ortho-martingale, and consider $i \leq n$ and $j \leq m$. Clearly,

$$\mathbb{E}\{M_{n,m}|\mathcal{F}_{i,j}\} = \mathbb{E}(\mathbb{E}\{M_{n,m}|\mathcal{F}_i^1\} | \mathcal{F}_{i,j}) = \mathbb{E}\{M_{i,m}|\mathcal{F}_{i,j}\}.$$

Now, project m to j to see that orthomartingale implies martingale. In fact, for this half, commutation is not needed.

Conversely, suppose M is a martingale with respect to a commuting filtration. We will show it is an ortho-martingale for the marginal filtrations. By Doob's convergence theorem,

$$\mathbb{E}\{M_{i+1,j}|\mathcal{F}_i^1\} = \lim_{k \rightarrow \infty} \mathbb{E}\{M_{i+1,j}|\mathcal{F}_{i,k}\}.$$

Thanks to commutation, the above is $\lim_k \mathbb{E}\{M_{i+1,j}|\mathcal{F}_{i,j}\} = M_{i,j}$. Similarly, $\mathbb{E}\{M_{i,j+1}|\mathcal{F}_j^2\} = M_{i,j}$, and we are done. \square

Thus, in the presence of commutation, we have maximal inequalities, convergence theorems, etc.

Lecture 7

Capacity, Energy and Dimension

We now come to the second part of these lectures which has to do with “exceptional sets”. The most obvious class of exceptional sets are those of measure 0, where the measure is some nice one. As an example, consider a compact set $E \subset \mathbb{R}^d$. One way to construct its Lebesgue measure is as follows: cover E by small boxes, compute the volume of the cover, and then optimize over all the covers. That is,

$$|E| = \lim_{\varepsilon \rightarrow 0^+} \inf \left\{ \sum_i [\text{diam}(E_i)]^d : E_1, E_2, \dots \text{ closed boxes of diameter } \leq \varepsilon \text{ with } \cup_i E_i \supseteq E \right\}.$$

Here, we are computing the diameter of the box as twice its ℓ^1 -radius; i.e., it is the length of any side. This is equivalent to the usual definition of Lebesgue’s measure, although it is long out of fashion in standard analysis courses.

1 Hausdorff Dimension and Measures

The first class of exceptional sets that we can discuss are those of Lebesgue’s measure 0, of course. But, this is too crude for differentiating amongst very thin sets. For example, consider the rationals \mathbb{Q} , as well as Cantor’s tertiary set \mathbf{C} . While they are both measure 0 sets, \mathbf{C} is uncountable, whereas \mathbb{Q} is not. We would like a concrete way of saying that \mathbf{C} is larger than \mathbb{Q} , and perhaps measure how much larger, as well. There are many ways of doing this, and we will choose a route that is useful for our probabilistic needs. First, note that for any $\alpha \geq 0$, we can define the analogue of $|E|$ as above. Namely, define

$$\mathcal{H}_\alpha(E) = \lim_{\varepsilon \rightarrow 0^+} \inf \left\{ \sum_i [\text{diam}(E_i)]^\alpha : E_1, E_2, \dots \text{ closed boxes of diameter } \leq \varepsilon \text{ with } \cup_i E_i \supseteq E \right\}.$$

This makes sense even if $\alpha \leq 0$.

The set function \mathcal{H}_α is called the α -dimensional *Hausdorff measure*. This terminology is motivated by the following, which is proved by using the method given to us by Carathéodory:

Theorem 1.1 *The set function \mathcal{H}_α is an outer measure on Borel subsets of \mathbb{R}^d . For all $\alpha > d$, $\mathcal{H}_\alpha(E) = 0$ identically. On the other hand, when $\alpha \leq d$ is an integer, $\mathcal{H}_\alpha(E)$ equals the α -dimensional Lebesgue's measure of Borel set E .*

Hausdorff dimensions provide us with a more refined sense of how big a set is. Note that for any compact (or even Borel, say) set E , there is *always* a critical α such that for all $\beta < \alpha$, $\mathcal{H}_\beta(E) = 0$, while for all $\beta > \alpha$, $\mathcal{H}_\beta(E) = +\infty$. This is an easy calculation. But it leads to the following important measure-theoretic notion of dimension:

$$\dim(E) = \inf\{\alpha : \mathcal{H}_\alpha(E) = 0\} = \sup\{\alpha : \mathcal{H}_\alpha(E) = +\infty\}.$$

This is the *Hausdorff dimension* of E .

How does one compute the Hausdorff dimension of a set? You typically proceed by establishing an upper bound, as well as a lower bound. The first step is not hard: just find a “good” covering E_i of diameter less than ε , and compute $\sum_i [\text{diam}(E_i)]^\alpha$. Here is one way to get an upper bound systematically; other ways abound.

Suppose we are interested in computing the Hausdorff dimension of a given compact set $E \subset [0, 1]^d$. Fix a *real* number $n \geq 1$, and define $E_j = [\frac{j}{n}, \frac{j+1}{n}]$, for integers $0 \leq j \leq n$. Then, it is clear that the diameter of each E_j is no more than $\frac{2}{n}$, while $\cup_j E_j \supset E$. So,

$$\mathcal{H}_\alpha(E) \leq \left(\frac{2}{n}\right)^\alpha \mathcal{N}_n(E),$$

where $\mathcal{N}_n(E) = \sum_{0 \leq j \leq n} \mathbf{1}\{I_{j,n} \cap E \neq \emptyset\}$ is the number of times the intervals $I_{j,n}$ contains portions of E . Therefore, if we can find α such that $\limsup_n n^{-\alpha} \mathcal{N}_n(E) < +\infty$, we have $\dim(E) \leq \alpha$.¹ Incidentally, the minimal α such that $\limsup_n n^{-\alpha} \mathcal{N}_n(E) < +\infty$ is the so-called *upper Minkowski (or box) dimension* of E . If we write the latter as $\dim_M(E)$, we have shown that

$$\dim(E) \leq \dim_M(E). \tag{1.1}$$

If we replace E_j by a d -dimensional box of the form $[\frac{j_1}{n}, \frac{j_1+1}{n}] \times \dots \times [\frac{j_d}{n}, \frac{j_d+1}{n}]$ and repeat the procedure, we obtain the upper Minkowski dimension in d dimensions, and Eq. (1.1) remains to hold.

We now use this to obtain an upper bound for the tertiary Cantor set \mathbf{C} . First, let us recall the following iterative construction of \mathbf{C} : let $\mathbf{C}_0 = [0, 1]$. Now, remove the middle third to obtain $\mathbf{C}_1 = [0, \frac{1}{3}] \cup [\frac{2}{3}, 1]$. Next, remove the middle thirds of each of the two subintervals to get $\mathbf{C}_2 = [0, \frac{1}{9}] \cup [\frac{2}{9}, \frac{3}{9}] \cup [\frac{6}{9}, \frac{7}{9}] \cup [\frac{8}{9}, 1]$, and so on. In this way, you have a decreasing sequence of compact subsets of $[0, 1]$, and, as such, $\mathbf{C} = \cap_n \mathbf{C}_n$ is a nontrivial compact subset of $[0, 1]$. At the n th level of construction, \mathbf{C}_n is comprised of 2^n intervals of length

¹We do not require n to be an integer here.

3^{-n} . Therefore, $|\mathbf{C}_n| = (\frac{2}{3})^n \rightarrow |\mathbf{C}| = 0$. On the other hand, we just argued that there are 2^n boxes, of diameter no greater than (in fact, equal to) 3^{-n} that cover \mathbf{C} . Therefore, we have shown that $\mathcal{N}_{3^{-n}}(\mathbf{C}) = 2^n$. In particular, for any $\alpha > \log_3(2)$, $\limsup_{m \rightarrow \infty} (3^{-m})^{-\alpha} \mathcal{N}_{3^{-m}}(E) = \lim_{m \rightarrow \infty} 3^{-m\alpha} 2^m = 0$. So that, after a little work, we get $\dim_M(E) \leq \log_3(2)$. In fact, it is easy to see, by the same reasoning, that $\dim_M(E) = \log_3(2)$. In any event, we obtain the following:

$$\dim(\mathbf{C}) \leq \log_3(2) = \frac{\ln 2}{\ln 3}. \quad (1.2)$$

We will show that this is sharp in that the above inequality is an equality. But first, a question: *why not stick to Minkowski dimension?* It is certainly easier to compute than Hausdorff dimension, and at first sight, more natural. To answer this, try computing $\dim_M(\mathbb{Q})$, or \dim_M of any other dense subset of $[0, 1]^d$ for that matter! You will see that the answer is 1! On the other hand, it is not hard to show that $\dim(E) = 0$ if E is countable, for then we can write $E = \{r_i\}$ and note that $\{r_i\}$ is a cover of E with diameter less than ε . This seemingly technical difference is really a big one.

Now, to the lower bound for $\dim(\mathbf{C})$. Obtaining lower bound on Hausdorff dimension is, in principle, very hard, since you have to work uniformly over all covers. What makes things difficult is that there are *alot* of potential covers!

The ingenious idea behind obtaining lower bounds is due to O. Frostman who found it in his Ph.D. thesis in the 1935! Namely,

Theorem 1.2 (Frostman's lemma) *Suppose we knew that the compact set E carries a probability measure μ that is Hölder-smooth in the following sense: there exists $\alpha > 0$ and a constant C such that for all $r \in (0, 1)$,*

$$\mu(\mathcal{B}(y, r)) \leq Cr^\alpha,$$

for μ -almost all y , where $\mathcal{B}(y, r)$ is the ℓ^∞ -ball of radius r about $y \in \mathbb{R}^d$. Then, $\dim(E) \geq \alpha$.

There is a converse to this that we will only need once, and will not prove, as a result; for a proof, see Appendix C of MPP.

Theorem 1.3 (Frostman's Lemma (continued)) *Suppose $\dim(E) \geq \alpha > 0$. Then, for each $\beta < \alpha$, there exists $\mu \in \mathcal{P}(E)$ such that*

$$\sup_{x \in \mathbb{R}^d} \sup_{r \in (0, 1)} \frac{\mu\{\mathcal{B}(x, r)\}}{r^\beta} < +\infty.$$

Proof I will prove this when instead of μ -almost all x , the lemma holds for all x . The necessary modifications to prove the general case are technical but not hard.

Fix $\varepsilon \in (0, 1)$, and consider any cover E_1, E_2, \dots of diameter $\leq \varepsilon$. Note that

$$1 = \mu(E) \leq \sum_i \mu(E_i) \leq C \sum_i [\text{diam}(E_i)]^\alpha.$$

Optimize over all such covers, and let $\varepsilon \rightarrow 0$, to see that $1 \leq 2C\mathcal{H}_\alpha(E)$. The theorem follows, since this shows that for any $\beta < \alpha$, $\mathcal{H}_\beta(E) = +\infty$. (To prove in the general case, note that if $\mu(E_j)$ is not less than $C[\text{diam}(E_j)]^\alpha$, we can cover E_j by at most 2^d compact intervals $F_{j,1}, \dots, F_{j,2^d}$ of diameter less than twice that of E_j , such that $\mu(F_{j,k}) \leq C[\text{diam}(F_{j,k})]^\alpha \leq 2^\alpha C[\text{diam}(E_j)]^\alpha$. Thus, $\mu(E_j) \leq 2^{\alpha+d} C[\text{diam}(E_j)]^\alpha$, which is good enough. \square)

We use this to complete our proof of the following.

Proposition 1.4 *If \mathbf{C} denotes the tertiary Cantor set, $\dim(\mathbf{C}) = \frac{\ln 2}{\ln 3}$.*

Proof In light of what we have already done, we only need to verify the lower bound on dimension. We do this by finding a sufficiently smooth measure on \mathbf{C} . Our choice is more or less obvious and is found iteratively as follows: construct the smoothest possible probability measure μ_n on \mathbf{C}_n and “take limits”. Now, the smoothest and flattest probability measure on \mathbf{C}_n is the uniform measure, μ_n . It is easy to see that for all $x \in [0, 1]$,

$$\mu_n([x - 3^{-n}, x + 3^{-n}]) \leq 2^{-n} = (3^{-n})^{\ln 2 / \ln 3}. \quad (1.3)$$

This is suggestive, but we need to work a little bit more. To do so, we next note that the μ_n 's are nested: We write $\mathbf{C}_n = \cup_{i=1}^{2^n} I_{i,n}$ where $I_{i,n}$ is an interval of length 3^{-n} . The nested property of the μ_n 's is the following, which can be checked by induction:

$$\forall n \geq m, \forall j = 1, \dots, 2^m : \quad \mu_n(I_{j,m}) = \mu_m(I_{j,m}) = 2^{-m}.$$

Standard weak convergence theory guarantees us of the existence of a probability measure μ_∞ on the compact set \mathbf{C} such that for all $m \geq 1$ and all $j = 1, \dots, 2^m$,

$$\mu_\infty(I_{j,m}) = \mu_m(I_{j,m}) = 2^{-m}.$$

Moreover, Eq. (1.3) extends to μ_∞ . Namely, for all $x \in [0, 1]$ and all $n \geq 0$,

$$\mu_\infty([x - 3^{-n}, x + 3^{-n}]) \leq (3^{-n})^{\ln 2 / \ln 3}.$$

Now, if $r \in (0, 1)$, we can find $n \geq 0$ such that $3^{-n-1} \leq r \leq 3^{-n}$. Therefore,

$$\sup_x \mu_\infty([x-r, x+r]) \leq \sup_x \mu_\infty([x-3^{-n}, x+3^{-n}]) \leq (3^{-n})^{\ln 2 / \ln 3} \leq (3r)^{\ln 2 / \ln 3}.$$

So, we have found a probability measure μ_∞ on \mathbf{C} , that satisfies the condition of Frostman's lemma with $C = 3^{\ln 2 / \ln 3} = 2$ and $\alpha = \ln 3 / \ln 3$. This completes our proof. \square

2 Energy and Capacity

Suppose μ is a probability measure on some given compact set $E \subset \mathbb{R}^d$. We will write this as $\mu \in \mathcal{P}(E)$, and define for any measurable function $f : E \times E \rightarrow \mathbb{R}_+ \cup \{\infty\}$,

$$\mathcal{E}_f(\mu) = \iint f(x, y) \mu(dx) \mu(dy).$$

This is the *energy* of μ with respect to the gauge function f ; it is always defined although it may be infinite. The following energy forms are of use to us:

$$\text{Energy}_\alpha(\mu) = \iint |x - y|^{-\alpha} \mu(dx) \mu(dy),$$

where $|x| = \max_{1 \leq j \leq d} |x_j|$ for concreteness, although any other Euclidean norm will do just as well. This is the so-called α -dimensional *Bessel–Riesz energy* of μ . The question, in the flavor of the previous section, is *when does a set E carry a probability measure of finite energy?* To facilitate the discussion, we define the *capacity* of a set E by

$$\mathcal{C}_f(E) = \left[\inf_{\mu \in \mathcal{P}(E)} \mathcal{E}_f(\mu) \right]^{-1}, \quad \text{and in particular,}$$

$$\text{Cap}_\alpha(E) = \left[\inf_{\mu \in \mathcal{P}(E)} \text{Energy}_\alpha(\mu) \right]^{-1}.$$

The above is Gauss' principle of minimum energy. Next, we argue that there is a minimum energy measure called the equilibrium measure. Moreover, its potential is essentially constant, and the constant is the energy.

Theorem 2.1 (Equilibrium Measure) *Suppose E is a compact set in \mathbb{R}^d such that for some $\alpha > 0$, $\text{Cap}_\alpha(E) > 0$. Then, there exists $\mu \in \mathcal{P}(E)$, such that*

$$\text{Energy}_\alpha(\mu) = [\text{Cap}_\alpha(E)]^{-1}.$$

Moreover, for μ -almost all x ,

$$\int |x - y|^{-\alpha} \mu(dy) = \text{Energy}_\alpha(\mu).$$

Proof By definition, there exists a sequence of probability measures μ_n , all supported on E , such that (i) they have finite energy; and (ii) for all $n \geq 1$, $(1 + \frac{1}{n})[\text{Cap}_\alpha(E)]^{-1} \geq \text{Energy}_\alpha(\mu_n) \geq [\text{Cap}_\alpha(E)]^{-1}$. Let μ be any subsequential limit of the μ_n 's. Since $\mu \in \mathcal{P}(E)$ as well, $\text{Energy}_\alpha(\mu) \geq [\text{Cap}_\alpha(E)]^{-1}$. We aim to show the converse holds too. By going to a subsequence n' along which $\mu_{n'}$ converges weakly to μ , we see that for any $r_0 > 0$,

$$\iint_{|x-y| \geq r_0} |x-y|^{-\alpha} \mu(dx) \mu(dy) = \lim_{n' \rightarrow \infty} \iint_{|x-y| \geq r_0} |x-y|^{-\alpha} \mu_{n'}(dx) \mu_{n'}(dy) \leq [\text{Cap}_\alpha(E)]^{-1}.$$

Let $r_0 \downarrow 0$ and use the dominated convergence theorem to deduce the first assertion. For the second assertion, i.e., that the minimum energy principle is actually achieved for some probability measure.

Lecture 8

Frostman's Theorem, Hausdorff Dimension and Brownian Motion

1 Frostman's Theorem (Continued)

Now, consider

$$\Upsilon_\eta = \left\{ x \in E : \int |x - y|^{-\alpha} \mu(dy) < (1 - \eta) \text{Energy}_\alpha(\mu) \right\}, \quad \eta \in (0, 1).$$

We wish to show that $\mu(\Upsilon_\eta) = 0$ for all $\eta \in (0, 1)$. If this is not the case for some $\eta \in (0, 1)$, then, consider the following

$$\zeta(\bullet) = \frac{\mu(\bullet \cap \Upsilon_\eta)}{\mu(\Upsilon_\eta)}.$$

Evidently, $\zeta \in \mathcal{P}(E)$, and has finite energy. Define

$$\lambda_\varepsilon = (1 - \varepsilon)\mu + \varepsilon\zeta, \quad \varepsilon \in (0, 1).$$

Then, λ_ε is also a probability measure on E , and it, too, has finite energy. In fact, writing $\lambda_\varepsilon = \mu - \varepsilon(\mu - \zeta)$, a little calculation shows that

$$\text{Energy}_\alpha(\lambda_\varepsilon) = \text{Energy}_\alpha(\mu) + \varepsilon^2 \text{Energy}_\alpha(\mu - \zeta) - 2\varepsilon \iint |x - y|^{-\alpha} \mu(dx) [\mu(dy) - \zeta(dy)].$$

(The energy of $\mu - \zeta$ is defined as if $\mu - \zeta$ were a positive measure.)

Since μ minimizes energy, the above is greater than or equal to $\text{Energy}_\alpha(\mu)$. Thus,

$$\varepsilon^2 \text{Energy}_\alpha(\mu - \zeta) \geq 2\varepsilon \iint |x - y|^{-\alpha} \mu(dx) [\mu(dy) - \zeta(dy)].$$

Divide by ε and let $\varepsilon \rightarrow 0$ to see that

$$\text{Energy}_\alpha(\mu) \leq \iint |x - y|^{-\alpha} \mu(dx) \zeta(dy).$$

But by the definition of Υ_η , the right hand side is no more than $(1 - \eta)\text{Energy}_\alpha(\mu)$, which contradicts the assumption that $\mu(\Upsilon_\eta) > 0$. In other words,

$$\int |x - y|^{-\alpha} \mu(dy) \geq \text{Energy}_\alpha(\mu), \quad \mu\text{-a.s.}$$

It suffices to show the converse inequality. But this is easy. Indeed, suppose

$$\mathfrak{G}\mu(x) = \int |x - y|^{-\alpha} \mu(dy) \geq (1 + \eta)\text{Energy}_\alpha(\mu),$$

on a set of positive μ -measure. The function $x \mapsto \mathfrak{G}\mu(x)$ is the α -dimensional potential of the measure μ . We could integrate $[d\mu]$ to get the desired contradiction, viz.,

$$\begin{aligned} \text{Energy}_\alpha(\mu) &= \int_{\Theta_\eta} \mathfrak{G}\mu(x) \mu(dx) + \int_{\Theta_\eta^c} \mathfrak{G}\mu(x) \mu(dx) \\ &\geq (1 + \eta)\text{Energy}_\alpha(\mu) \cdot \mu(\Theta_\eta) + \int_{\Theta_\eta^c} \mathfrak{G}\mu(x) \mu(dx), \end{aligned}$$

where $\Theta_\eta = \{x : \mathfrak{G}\mu(x) \geq (1 + \eta)\text{Energy}_\alpha(\mu)\}$. Therefore, by Theorem 2.1 on equilibrium measure,

$$\begin{aligned} \text{Energy}_\alpha(\mu) &\geq \text{Energy}_\alpha(\mu) \left[(1 + \eta)\mu(\Theta_\eta) + \mu(\Theta_\eta^c) \right] \\ &= \text{Energy}_\alpha(\mu) \left[1 + \eta\mu(\Theta_\eta) \right], \end{aligned}$$

which is a contradiction, unless $\mu(\Theta_\eta) = 0$. This concludes our proof. \square

SOMETHING TO TRY: The α -dimensional Bessel–Riesz energy defines a Hilbertian pre-norm. Indeed, define $\mathcal{M}_\alpha(E)$ to be the collection of all measures of finite α -dimensional Bessel–Riesz energy on E . On this, define the inner product,

$$\langle \mu, \nu \rangle = \iint |x - y|^{-\alpha} \mu(dx) \nu(dy).$$

Check that this defines a positive-definite bilinear form on $\mathcal{M}_\alpha(E)$ if $\alpha \in (0, d)$. From this, conclude that for all $\mu, \nu \in \mathcal{M}_\alpha(E)$, $\langle \mu, \nu \rangle^2 \leq \text{Energy}_\alpha(\mu) \cdot \text{Energy}_\alpha(\nu)$. This fills a gap in the above proof.

The *capacitary dimension* of a compact set $E \subset \mathbb{R}^d$ is defined as

$$\dim_c(E) = \sup \{ \alpha : \text{Cap}_\alpha(E) > 0 \} = \inf \{ \alpha : \text{Cap}_\alpha(E) = 0 \}.$$

Theorem 1.1 (Frostman's Theorem) *Capacitary and Hausdorff dimensions are one and the same.*

Proof Here is one half of the proof: we will show that if there exists $\alpha > 0$ and a probability measure μ on E , such that $\text{Energy}_\alpha(\mu) < +\infty \Rightarrow \dim(E) \geq \alpha$. This shows that $\dim_c(E) \leq \dim(E)$, which is half the theorem.

By Theorem 2.1, we can assume without loss of generality that μ is an equilibrium measure. In particular,

$$\mu(\mathcal{B}(x, r)) \leq r^\alpha \int |x - y|^{-\alpha} \mu(dy) = r^\alpha \text{Energy}_\alpha(\mu),$$

μ -almost everywhere. Frostman's lemma (Theorem 1.2) shows that $\dim(E) \geq \alpha$, as needed.

For the other half, we invoke the second half of Frostman's theorem (Theorem 1.3) to produce for each $\beta < \dim(E)$ a probability measure $\mu \in \mathcal{P}(E)$, such that

$$\mu(\mathcal{B}(x, r)) \leq Cr^\beta, \quad \forall x \in \mathbb{R}^d, r \in (0, 1).$$

But if D denotes the diameter of E ,

$$\begin{aligned} \text{Energy}_\gamma(\mu) &= \sum_{j=0}^{\infty} \iint_{2^{-j-1}D \leq |x-y| \leq 2^{-j}D} |x-y|^{-\gamma} \mu(dx) \mu(dy) \\ &\leq \sum_{j=0}^{\infty} 2^{(j+1)\gamma} D^{-\gamma} \sup_{x \in \mathbb{R}^d} \mu(\mathcal{B}(x, 2^{-j}D)) \\ &\leq C2^\gamma D^{\beta-\gamma} \sum_{j=0}^{\infty} 2^{j\gamma} 2^{-j\beta}, \end{aligned}$$

which sums if $\gamma < \beta$. Thus, we have shown that for all $\gamma < \dim(E)$, $\text{Cap}_\gamma(E) > 0$, i.e., $\dim_c(E) \geq \gamma$ for all $\gamma < \dim(E)$, which completes the proof. \square

2 The Brownian Curve

Next, we roll up our sleeves and compute the Hausdorff dimension of a few assorted and interesting random fractals that arise from Brownian considerations. Our goal is to illustrate the methods and ideas rather than the final word on this subject.

Throughout, $B = \{B_t; t \geq 0\}$ denotes Brownian motion in \mathbb{R}^d . That is, a Gaussian process in \mathbb{R}^d such

that $B_0 = 0$, and

$$\begin{aligned} \mathbb{E}\{B_t^i\} &= 0 & \forall t \geq 0, i = 1, \dots, d \\ \mathbb{E}\{B_s^i B_t^j\} &= \begin{cases} s \wedge t & \text{if } i = j \\ 0 & \text{otherwise} \end{cases} \end{aligned}$$

Recall also that B is a strong Markov process. Then, we have already shown that

$$B \text{ hits points iff } d = 1, \text{ i.e., } \exists t > 0 : B_t = 0 \iff d = 1.$$

This follows from our proof of Lévy's theorem (Theorem 3.1, Lecture 2.) In particular, note that when $d = 1$, the Brownian curve has full Lebesgue measure, and also full dimension. On the other hand, when $d \geq 2$, the Brownian curve has zero Lebesgue measure, despite the following result.

Theorem 2.1 *If B denotes d -dimensional Brownian motion, where $d \geq 2$, $\dim B(\mathbb{R}_+) = 2$, a.s.*

Proof We do this in two parts. First, we show that $\dim B(\mathbb{R}_+) \leq 2$ (*the upper bound*), and then we show that $\dim B(\mathbb{R}_+) \geq 2$ (*the lower bound*). In any event, recall that $d \geq 2$.

Proof of the upper bound Recall from Theorem 1.2, Lecture 3, that for any interval $I \subset \mathbb{R}^d$,

$$\mathbb{P}\{B[1, 2] \cap I \neq \emptyset\} \leq c\kappa(|I|), \text{ where } \kappa(\varepsilon) = \begin{cases} \varepsilon^{d-2}, & \text{if } d \geq 3 \\ \ln_+(\frac{1}{\varepsilon}), & \text{if } d = 2 \end{cases}$$

A careful inspection of the proof shows that the constant c depends only on M , as long as $I \subseteq [-M, M]^d$. Consider I_1, \dots, I_{n^d} cubes of side $\frac{1}{n}$, such that (i) $I_i^\circ \cap I_j^\circ = \emptyset$ if $i \neq j$; and (ii) $\cup_{j=1}^{n^d} I_j = [0, 1]^d$. Based on these, define

$$E_j = \begin{cases} I_j, & \text{if } I_j \cap B[1, 2] \neq \emptyset \\ \emptyset, & \text{otherwise} \end{cases}$$

Note that E_1, \dots, E_{n^d} is a $(\frac{1}{n})$ -cover of $B[1, 2] \cap [0, 1]^d$. Thus,

$$\mathcal{H}_\alpha(B[1, 2] \cap [0, 1]^d) \leq \liminf_{n \rightarrow \infty} \sum_{j=1}^{n^d} n^{-\alpha} \mathbf{1}_{\{I_j \cap B[1, 2] \neq \emptyset\}}.$$

Consequently, as long as $\alpha > 2$,

$$\mathbb{E}\{\mathcal{H}_\alpha(B[1, 2] \cap [0, 1]^d)\} \leq c \liminf_{n \rightarrow \infty} \sum_{j=1}^{n^d} n^{-\alpha} \kappa(\frac{1}{n}) = c \liminf_{n \rightarrow \infty} n^{d-\alpha} \kappa(\frac{1}{n}) = 0.$$

In particular, $\dim(B[1, 2] \cap [0, 1]^d) \leq 2$, a.s. Similarly, $\dim(B[a, b] \cap [-n, n]^d) \leq 2$, a.s. for any $0 < a < b$ and $n > 0$. Let $n \uparrow \infty$, $a \downarrow 0$ and $b \uparrow \infty$, all along rational sequences to deduce that $\dim B(\mathbb{R}_+) \leq 2$, a.s. This uses the easily verified fact that whenever $A_1 \subseteq A_2 \subseteq \dots$ are compact, and if $\mathcal{H}_\alpha(A_j) = 0$, then $\mathcal{H}_\alpha(\cup_j A_j) = 0$.

Proof of the lower bound For the converse, we will show that $\dim B[1, 2] \geq 2$, and do this by appealing to Frostman's theorem (Theorem 1.1, Lecture 7). To do so, we need to define a probability, or at least a finite, measure on the Brownian curve. The most natural measure that lives on the curve of $\{B_s; 1 \leq s \leq 2\}$ is the occupation measure:

$$\mathbb{O}(E) = \int_1^2 \mathbf{1}_{\{B_s \in E\}} ds.$$

With this in mind, note that for any $\alpha > 0$,

$$\text{Energy}_\alpha(\mathbb{O}) = \iint |x - y|^{-\alpha} \mathbb{O}(dx) \mathbb{O}(dy) = \int_1^2 \int_1^2 |B_s - B_t|^{-\alpha} ds dt.$$

By Frostman's theorem, it suffices to show that $\mathbb{E}\{\text{Energy}_\alpha(\mathbb{O})\} < +\infty$ for all $0 < \alpha < 2$. But this is easy. Indeed, note that

$$\mathbb{E}\{\text{Energy}_\alpha(\mathbb{O})\} = 2 \int_1^2 \int_s^2 \mathbb{E}\{|B_{t-s}|^{-\alpha}\} ds dt = 2 \int_1^2 \int_s^2 |t-s|^{-\frac{\alpha}{2}} ds dt \times \mathbb{E}\{|Z|^{-\alpha}\},$$

where Z is a d -dimensional vector of i.i.d. standard normals. Since $\alpha < 2$, the double integral is finite. It suffices to show that $\mathbb{E}\{|Z|^{-\alpha}\} < +\infty$. But

$$\begin{aligned} \mathbb{E}\{|Z|^{-\alpha}\} &= \int_0^\infty \mathbb{P}\{|Z|^{-\alpha} > \lambda\} d\lambda \\ &\leq 1 + \int_1^\infty \mathbb{P}\{|Z|^{-\alpha} > \lambda\} d\lambda \\ &= 1 + \alpha \int_0^1 \mathbb{P}\{|Z| < u\} u^{-\alpha-1} du && (u = \lambda^{-\frac{1}{\alpha}}) \\ &= 1 + \alpha \int_0^1 [\mathbb{P}\{|Z_1| \leq u\}]^d u^{-\alpha-1} du. \end{aligned}$$

But $\mathbb{P}\{|Z_1| \leq u\} = (2\pi)^{-\frac{1}{2}} \int_{-u}^u e^{-\frac{1}{2}\lambda^2} d\lambda \leq u$. Hence, using $d \geq 2$,

$$\mathbb{E}\{|Z|^{-\alpha}\} \leq 1 + \alpha \int_0^1 u^{d-\alpha-1} du \leq 1 + \alpha \int_0^1 u^{1-\alpha} du,$$

which is finite, as promised. □

Lecture 9

Potential Theory of Brownian Motion and Stable Processes

1 Transient Brownian Motion

Classically, probabilistic potential theory has been concerned with connections between Newtonian (and more general potentials) and hitting probabilities for Markov processes. This is a rich theory that we have already been introduced to in earlier lectures in the following form: d -dimensional Brownian motion hits points iff $d = 1$. In fact, we did this by going alot further. Namely, we estimated the hitting probability of a small ball, where by hitting probability we mean something like the quantity $\mathbb{P}\{B[1, 2] \cap \mathcal{B}(x, \varepsilon) \neq \emptyset\}$. The classical connections run deep: harmonic and excessive functions, removable singularity for various elliptic PDE's on domains, etc.

To see immediate connections between Brownian motion and Bessel–Riesz potentials (here, Newtonian potentials), consider Brownian motion in \mathbb{R}^d where $d \geq 3$, and define the potential operator \mathcal{U} as

$$\mathcal{U}f(x) = \mathbb{E}_x \left\{ \int_0^\infty f(B_s) ds \right\} = \mathbb{E} \left\{ \int_0^\infty f(B_s + x) ds \right\}.$$

We compute this as follows: for all measurable $f : \mathbb{R}^d \rightarrow \mathbb{R}_+$,

$$\begin{aligned} \mathcal{U}f(x) &= \int_{\mathbb{R}^d} f(y) \int_0^\infty \frac{e^{-\frac{1}{2s}\|x-y\|^2}}{(2\pi s)^{\frac{d}{2}}} ds dy \\ &= \frac{\Gamma(\frac{d}{2} - 1)}{2\pi^{\frac{d}{2}}} \int_{\mathbb{R}^d} \frac{f(y)}{\|x - y\|^{d-2}} dy. \end{aligned}$$

Thus, if $\mathfrak{G}f$ denotes the $(d - 2)$ -dimensional Newtonian (and/or Bessel–Riesz) potential of f defined by $\mathfrak{G}f(x) = \int f(y)|x - y|^{2-d} dy$, it follows that $\mathcal{U}f \asymp \mathfrak{G}f$.

Now, let $E \subset \mathbb{R}^d$ denote a compact set whose hitting time is $\tau(E) = \inf\{s \geq 0 : B_s \in E\}$, and consider a sub-probability density $f : \mathbb{R}^d \rightarrow \mathbb{R}_+$ such that $f = 0$ off of E (OK, provided that E° has positive Lebesgue's measure.) Then,

$$\begin{aligned} \mathcal{U}f(x) &= \mathbb{E}_x \left\{ \int_{\tau(E)}^{\infty} f(B_s) ds; \tau(E) < +\infty \right\} \\ &= \mathbb{E}_x \left\{ \int_0^{\infty} f(B'_s + B_{\tau(E)}) ds; \tau(E) < +\infty \right\}, \end{aligned}$$

where B' is an independent copy of B . This holds by the strong Markov property. Hence,

$$\begin{aligned} \mathcal{U}f(x) &= \mathbb{E}_x \{ \mathcal{U}f(B_{\tau(E)}); \tau(E) < +\infty \} \\ &= \mathbb{E}_x \{ \mathcal{U}f(B_{\tau(E)}) \mid \tau(E) < +\infty \} \cdot \mathbb{P}_x \{ \tau(E) < +\infty \} \\ &= \int_{\mathbb{R}^d} \mathcal{U}f(x) \mu(dx) \cdot \mathbb{P}_x \{ \tau(E) < +\infty \}, \end{aligned}$$

where $\mu(A) = \mathbb{P}_x \{ B_{\tau(E)} \in A \mid \tau(E) < \infty \}$. Since E° has positive Lebesgue's measure, and since B has continuous samples, $\mu \in \mathcal{P}(E)$, and the above holds for all pdf's f on E . Clearly, $\mu(dx) \ll dx$. Suppose that E is nice enough that $f = d\mu/dx$ is a probability density function. Note, then, that $\mathcal{U}f(x) = c \int_{\mathbb{R}^d} \|x - y\|^{2-d} \mu(dy)$ and $\int \mathcal{U}f d\mu = c \iint \|x - y\|^{2-d} \mu(dx) \mu(dy)$ (for the same c), which yields:

$$\mathbb{P}_x \{ \tau(E) < +\infty \} \asymp \frac{\int_{\mathbb{R}^d} \|x - y\|^{2-d} \mu(dy)}{\text{Energy}_{d-2}(\mu)}.$$

Next, suppose that the starting point x is strictly outside E , so that for all $y \in E$, $\|x - y\| \asymp 1$. This shows that

$$\mathbb{P}_x \{ B(\mathbb{R}_+) \cap E \neq \emptyset \} \asymp \frac{1}{\text{Energy}_{d-2}(\mu)},$$

for some probability measure $\mu \in \mathcal{P}(E)$. In particular,

$$\mathbb{P}_x \{ B(\mathbb{R}_+) \cap E \neq \emptyset \} \leq c \text{Cap}_{d-2}(E). \quad (1.1)$$

It is relatively easy to argue that, in fact, the above is essentially optimal. Indeed, note that if $J(h) = \int_0^\infty h(B_s) ds$, when h is any pdf on E ,

$$\begin{aligned} \mathbb{E}_x \{ J(h) \} &= \int_{\mathbb{R}^d} \int_0^\infty h(y) \frac{e^{-\frac{1}{2s} \|x-y\|^2}}{(2\pi s)^{\frac{d}{2}}} ds dy \\ &\geq c, \end{aligned}$$

since $\sup_{y \in E} \|x - y\| \leq C$ in the above, and since $\int h dy = 1$. On the other hand, similar calculations reveal that $\mathbb{E}_x \{ |J(h)|^2 \} \leq C \text{Energy}_{d-2}(h)$. Thus, by the Paley-Zygmund inequality,

$$\mathbb{P}_x \{ \tau(E) < +\infty \} \geq c \left[\inf_{h=\text{pdf on } E} \text{Energy}_{d-2}(h) \right]^{-1}.$$

The above describes an absolutely continuous capacity that can be shown to coincide with $\text{Cap}_{d-2}(E)$. Modulo this last step and the part about “ E being nice”, we have shown the following which is part of a greater theorem of S. Kakutani.

Theorem 1.1 (Kakutani) *When $d \geq 3$, Brownian motion in \mathbb{R}^d can hit a compact set E iff $\text{Cap}_{d-2}(E) > 0$.*

Combining this with Frostman’s theorem, we get:

$$\begin{aligned} \dim(E) > d - 2 &\implies B \text{ can hit } E \\ \dim(E) < d - 2 &\implies B \text{ cannot hit } E. \end{aligned} \tag{1.2}$$

For $d = 1, 2$, the above proof breaks down, since our definition of $\mathcal{U}f$ is typically infinite; this is due to the neighborhood recurrence of Brownian motion in \mathbb{R}^d , $d \leq 2$. However, the basic principle is still correct, as long as we do something about times near $+\infty$. The classical way to do this is to introduce an independent mean 1 exponential random variable ϵ and define the 1-potential operator

$$\mathcal{U}_1 f(x) = \mathbb{E}_x \left\{ \int_0^\epsilon f(B_s) ds \right\}.$$

The reason for the choice of the exponential law is that B stopped at ϵ is still a strong Markov process. Also, check that

$$\mathcal{U}_1 f(x) = \mathbb{E}_x \left\{ \int_0^\infty e^{-s} f(B_s) ds \right\}.$$

Now, proceed as in the proof of Kakutani’s theorem above, but pay attention to the cases $d = 1$ and $d = 2$ separately; we also need to replace $\mathcal{U}f$ by $\mathcal{U}_1 f$ everywhere. This will lead to the complete

Theorem 1.2 (Kakutani’s Theorem) *Brownian motion in any dimension d hits a compact set E iff $\text{Cap}_{d-2}(E) > 0$.*

The result is trivial when $d = 1$ and $\text{Cap}_{-1}(E) = 1$ for all E . So, the content is in dimension 2, where Brownian motion hits only (and all) compact sets of positive logarithmic capacity. In particular, Eq. (1.2) holds in all dimensions $d \geq 2$.

2 Additive Brownian Motion

Suppose X and Y are two independent d -dimensional Brownian motions. We have already encountered the problem of deciding when $X(\mathbb{R}_+) \cap Y(\mathbb{R}_+) \neq \emptyset$. One way to interpret this is by the following identification:

$X(\mathbb{R}_+) \cap Y(\mathbb{R}_+) \neq \emptyset$ the origin in \mathbb{R}^d . One can ask, more generally, what types of sets can Z hit? By symmetry, we can rewrite Z (in law) as

$$Z_{s,t} = X_s + Y_t, \quad \forall s, t \geq 0,$$

and call it *additive Brownian motion* (with 2 parameters). More generally, we can define N -parameter additive Brownian motion Z as

$$Z_t = \sum_{j=1}^N B_{t_j}^j, \quad \forall t \in \mathbb{R}_+^N,$$

where B^1, \dots, B^N are iid d -dimensional Brownian motions. We can refer to Z as (N, d) -additive Brownian motion to keep the dimensions straight.

Theorem 2.1 (Hirsch and Song; Kh and Shi) (N, d) -additive Brownian motion can hit a compact set E iff $\text{Cap}_{d-2N}(E) > 0$.

Sketch of Proof We will sketch an argument that shows that if Z is (N, d) -additive Brownian motion, $\mathbb{P}\{Z[1, 2]^N \cap E \neq \emptyset\} > 0$ iff $\text{Cap}_{d-2N}(E) > 0$. Going from $[1, 2]^N$ to \mathbb{R}_+^N is standard, since the same argument shows that for any cube $[a, b] = \prod_{j=1}^N [a_j, b_j]$ with $a_j > 0$, $\mathbb{P}\{Z[a, b] \cap E \neq \emptyset\} > 0$ iff $\text{Cap}_{d-2N}(E) > 0$. Thus, we can let $a_j \downarrow 0$ and $b_j \uparrow \infty$ to finish. With this in mind, consider

$$J(f) = \int_{[1, 2]^N} f(Z_s) ds.$$

Then, one shows that if f_ε is a pdf on the ε -enlargement E^ε of E , $\inf_{\varepsilon > 0} \mathbb{E}\{J(f_\varepsilon)\} \geq c$ and $\mathbb{E}\{|J(f_\varepsilon)|^2\} \leq c \text{Energy}_{d-4}(f_\varepsilon)$. In both of these estimates, c depends on the outer radius of E only. But $\{J(f_\varepsilon) > 0\}$ implies $\{Z[1, 2]^N \cap E^\varepsilon \neq \emptyset\}$. Hence, $\mathbb{P}\{Z[1, 2]^N \cap E^\varepsilon \neq \emptyset\} \geq c \text{Cap}_{d-2N}(E^\varepsilon) \geq c \text{Cap}_{d-2N}(E)$. Let $\varepsilon \rightarrow 0$ and use the compactness of E , together with the continuity of Z .

For the harder converse, we will prove that $\mathbb{P}\{Z[1, \frac{3}{2}]^N \cap E \neq \emptyset\} > 0$ implies $\text{Cap}_{d-2N}(E) > 0$. Henceforth, we assume

$$\mathbb{P}\{Z[1, \frac{3}{2}]^N \cap E \neq \emptyset\} > 0, \tag{2.1}$$

and let \mathcal{F}_t denote the σ -field generated by $\{Z_r; r \preccurlyeq t\}$. Note that \mathcal{F} is commuting, and consider the N -parameter martingale $M(f)$ given by

$$M_t(f) = \mathbb{E}\{J(f) | \mathcal{F}_t\}, \quad \forall t \in \mathbb{R}_+^N.$$

Clearly, for all $t \in [1, \frac{3}{2}]^N$,

$$\begin{aligned} M_t(f) &\geq \int_{t \preceq s \preceq (2, \dots, 2)} \mathbb{E}\{f(Z_s) \mid \mathcal{F}_t\} ds \\ &= \int_{t \preceq s \preceq (2, \dots, 2)} \mathbb{E}\{f(Z'_{s-t} + Z_t \mid \mathcal{F}_t)\} ds, \end{aligned}$$

where Z' is an independent copy of Z . Therefore,

$$M_t(f) \geq \int_{[0, \frac{1}{2}]^N} \mathbb{E}\{f(Z'_r + Z_t) \mid \mathcal{F}_t\} dr, \quad \forall t \in [1, \frac{3}{2}] \cap \mathbb{Q}_+^N.$$

Now, suppose $T^\varepsilon \in \mathbb{Q}_+^N \cup \{\infty\}$ is any random variable such that

- (i) $T^\varepsilon \neq \infty$ iff $Z[1, \frac{3}{2}]^N \cap E^\varepsilon \neq \emptyset$ where E^ε is the closed ε -enlargement of E ;
- (ii) on $\{T^\varepsilon \neq \infty\}$, $Z_{T^\varepsilon} \in E^\varepsilon$.

The previous inequality for $M_t(f)$ implies,

$$M_{T^\varepsilon} \mathbf{1}_{\{T^\varepsilon \neq \infty\}} \geq \int_{[0, \frac{1}{2}]^N} \mathbb{E}\{f(Z'_r + Z_{T^\varepsilon}) \mid Z_{T^\varepsilon}, T^\varepsilon \neq \infty\} dr \cdot \mathbf{1}_{\{T^\varepsilon \neq \infty\}}.$$

Define, for simplicity, the operator $\mathcal{V}f(x) = \int_{[0, \frac{1}{2}]^N} \mathbb{E}\{f(Z_r + x)\} dr$, to see that

$$M_{T^\varepsilon} \mathbf{1}_{\{T^\varepsilon \neq \infty\}} \geq \mathcal{V}f(Z_{T^\varepsilon}) \cdot \mathbf{1}_{\{T^\varepsilon \neq \infty\}}.$$

Thanks to Eq. (2.1), $\mu_\varepsilon \in \mathcal{P}(E^\varepsilon)$, where

$$\mu_\varepsilon(\bullet) = \mathbb{P}\{Z_{T^\varepsilon} \in \bullet \mid T^\varepsilon \neq \infty\}.$$

Hence,

$$\begin{aligned} \mathbb{E}\left\{\sup_{t \in \mathbb{Q}_+^N} |M_t(f)|^2\right\} &\geq \mathbb{E}\left\{[\mathcal{V}f(Z_{T^\varepsilon})]^2 \cdot \mathbf{1}_{\{T^\varepsilon \neq \infty\}}\right\} \\ &= \mathbb{E}\left\{[\mathcal{V}f(Z_{T^\varepsilon})]^2 \mid T^\varepsilon \neq \infty\right\} \cdot \mathbb{P}\{T^\varepsilon \neq \infty\} \\ &\geq \left[\mathbb{E}\{\mathcal{V}f(Z_{T^\varepsilon}) \mid T^\varepsilon \neq \infty\}\right]^2 \cdot \mathbb{P}\{T^\varepsilon \neq \infty\} \\ &= \left[\int \mathcal{V}f(x) \mu_\varepsilon(dx)\right]^2 \cdot \mathbb{P}\{T^\varepsilon \neq \infty\}. \end{aligned}$$

By Cairoli's inequality, and by the mentioned second moment estimate for $J(f)$, $\mathbb{E}\{\sup_{t \in \mathbb{Q}_+^N} |M_t(f)|^2\} \leq 4^N \sup_t \mathbb{E}\{|M_t(f)|^2\} \leq c \text{Energy}_{d-2N}(f)$. This leads us to

$$c \text{Energy}_{d-2N}(f) \geq \left[\int \mathcal{V}f(x) \mu_\varepsilon(dx)\right]^2 \cdot \mathbb{P}\{T^\varepsilon \neq \infty\}.$$

Now, let $\varphi_\eta = (2\eta)^{-d} \mathbf{1}_{B(0,\eta)}$ be an approximation to the identity, and let $f = \mu_\varepsilon \star \varphi_\eta$. Then, f is a pdf on $E^{\eta+\varepsilon}$ and converges weakly to μ_ε as $\eta \downarrow 0$. One can show that $\text{Energy}_{d-2N}(f) \rightarrow \text{Energy}_{d-2N}(\mu_\varepsilon)$, as $\eta \rightarrow 0$, as well. We now wish to at least argue why $\int \mathcal{V}f d\mu_\varepsilon \rightarrow \text{Energy}_{d-2N}(\mu_\varepsilon)$ as $\eta \rightarrow 0$, as well. If so, we can deduce that

$$\mathbb{P}\{T^\varepsilon \neq \infty\} \leq c \div \text{Energy}_{d-2N}(\mu_\varepsilon) \leq c \text{Cap}_{d-2N}(E^\varepsilon),$$

which is our result but with E replaced by E^ε . Letting $\varepsilon \rightarrow 0$ will lead everything to converge to their proper limit. That is, $\mathbb{P}\{T^\varepsilon \neq \infty\} \rightarrow \mathbb{P}\{Z[1, \frac{3}{2}]^N \cap E \neq \emptyset\}$ (easy to see this), and $\text{Cap}_{d-2N}(E^\varepsilon) \rightarrow \text{Cap}_{d-2N}(E)$ (harder to show.) It remains to identify the limit as $\eta \rightarrow 0$ of $\int \mathcal{V}f d\mu_\varepsilon$. But

$$\begin{aligned} \int \mathcal{V}f d\mu_\varepsilon &= \int_{\mathbb{R}^d} \int_{[0, \frac{1}{2}]^N} \mathbb{E}\{f(Z_t + x)\} dt \mu_\varepsilon(dx) \\ &= \iint \int_{[0, \frac{1}{2}]^N} p_t(x - y) dt f(y) dy \mu_\varepsilon(dx), \end{aligned}$$

where $p_t(x) = (2\pi \sum_{i=1}^N t_i)^{-\frac{d}{2}} \exp(-\frac{\|x\|^2}{2 \sum_{i=1}^N t_i})$. Now, let $\eta \rightarrow 0$ to see that $f(y) dy \simeq \mu_\varepsilon(dy)$, which does the job. \square

3 Additive Stable Processes

Recall that an \mathbb{R}^d -valued process $X = \{X_t; t \geq 0\}$ is an *isotropic stable process* of index $\alpha \in (0, 2]$ if

St-1. For each t , $\mathbb{E}\{e^{i\xi \cdot X_t}\} = \exp(-\frac{1}{2}\|\xi\|^\alpha)$;

St-2. $X_0 = 0$, a.s.; and

St-3. for each $t \geq 0$, $s \mapsto X_{t+s} - X_t$ is a copy of X that is independent of $\{X_u; u \leq t\}$.

The condition that $\alpha \in (0, 2]$ is forced on us by the above conditions, and when $\alpha = 2$, X is Brownian motion.

To perform potential-theoretic calculations, we only need estimates for the pdf of X_t ; all else is done as in Brownian motion. Here are the requisite facts. All can be found in MPP; Chapters 8, and 10, together with proofs. Throughout, $p_t(x) = \mathbb{P}\{X_t \in dx\}/dx$.

By the Fourier inversion formula,

$$p_t(x) = (2\pi)^{-d} \int_{\mathbb{R}^d} e^{-i\xi \cdot x} e^{-\frac{1}{2}\|\xi\|^\alpha} d\xi.$$

This can be used to show to that

- $(t, x) \mapsto p_t(x)$ is continuous and strictly positive on compact subsets of $(0, \infty) \times \mathbb{R}^d$;

- **(scaling)** $p_t(x) = t^{-\frac{d}{\alpha}} p_1(x/t^{\frac{1}{\alpha}})$;
- **(unimodality)** $p_t(x) \leq p_t(0) = Ct^{-\frac{d}{\alpha}}$, for all $x \in \mathbb{R}^d$;
- **(isotropy)** $p_t(x) = p_t(y)$ if $\|x\| = \|y\|$;
- **(Blumenthal-Gettoor's asymptotics)** as $\|z\| \rightarrow \infty$, $p_1(z) \sim C\|z\|^{-(d+\alpha)}$.

An N -parameter *additive stable process* Z is defined by

$$Z_t = \sum_{j=1}^N X_{t_j}^j, \quad \forall t \in \mathbb{R}_+^N,$$

where X^1, \dots, X^N are i.i.d. isotropic stable processes of index $\alpha \in (0, 2]$. Then, the above together with the Brownian methods outlined earlier, can be used in conjunction to prove the following:

Theorem 3.1 (Hirsch and Song; MPP Ch. 11) *If Z denotes an (N, d) -additive stable process of index $\alpha \in (0, 2]$, then Z hits a compact set $E \subset \mathbb{R}^d$ iff $\text{Cap}_{d-\alpha N}(E) > 0$.*

This is attractive, since it relates Cap_β to a stochastic process for every $\beta \in \mathbb{R}_+$; when $N = 1$, this is classical, but only connects probability to Cap_β where $\beta \in [d - 2, d]$. One can use this to show also the following.

Theorem 3.2 (MPP Ch. 11) *Suppose Z is as above. Then, with probability one,*

$$\dim Z(\mathbb{R}_+^N) = d \wedge \alpha N.$$

To prove this, we only need the following, which is, in fact, a consequence of Theorem 3.1:

$$\mathbb{P}\{Z[1, 2]^N \cap \mathcal{B}(x, \varepsilon) \neq \emptyset\} \asymp C \times \begin{cases} \varepsilon^{d-\alpha N}, & \text{if } d > \alpha N \\ \{\log(1/\varepsilon)\}^{-1}, & \text{if } d = \alpha N \\ 1, & \text{if } d < \alpha N \end{cases}$$

In fact, the constants depend only on M , and the above holds uniformly for all $x \in [-M, M]^d$.

4 Application to Stochastic Codimension

The preceding has a remarkable consequence about a large class of random sets. We say that a random set $X \subset \mathbb{R}^d$ has *codimension* β , if β is the critical number such that for all compact sets $E \subset \mathbb{R}^d$ with $\dim(E) > \beta$, $\mathbb{P}\{X \cap E \neq \emptyset\} > 0$, while for all compact sets $F \subset \mathbb{R}^d$ with $\dim(F) < \beta$, $\mathbb{P}\{X \cap F \neq \emptyset\} = 0$. The notion of codimension was coined in this way in Kh-Shi '99, but the essential idea has been around in the works of Taylor '65, Lyons '99, Peres '95, ...

When it does exist, the codimension of a random set is a nonrandom number.

Warning: Not all random sets have a codimension.

The following is a fancy example for the above.

Theorem 4.1 (Kh-Peres-Xiao) *Let B denote Brownian motion, and consider*

$$F(\lambda) = \left\{ t \geq 0 : \limsup_{\varepsilon \rightarrow 0} \frac{|B_{t+\varepsilon} - B_t|}{\sqrt{2\varepsilon \log(1/\varepsilon)}} = \lambda \right\}.$$

Then, for any compact set $E \subset \mathbb{R}_+$, $\mathbb{P}\{F(\lambda) \cap E \neq \emptyset\} = 0$ if $\dim_p(E) < \lambda^2$, while it is 1 if $\dim_p(E) > \lambda^2$, where \dim_p denotes "packing dimension."

All that you need to know of packing dimension, here, is that it is not Hausdorff dimension although $\dim_p \geq \dim$. In fact, there are compact sets $E \subset \mathbb{R}_+$ such that $\dim_p(E) = 1$ while $\dim(E) = 0$. Multi-dimensional examples are also possible.

As examples of random sets that *do* have codimension, we mention the following consequence of Theorem 3.1:

Corollary 4.2 *If Z denotes an (N, d) -additive stable process of index $\alpha \in (0, 2]$, $\text{codim}(Z[1, 2]^N) = d - \alpha N$.*

We now wish to use Theorem 3.1 to prove the following result. In the present form, it is from MPP Ch. 11, but for $d = 1$, it is from Kh-Shi '99.

Theorem 4.3 (MPP Ch. 11) *If X is a random set in \mathbb{R}^d that has codimension $\beta \in (0, d)$,*

$$\dim(X) = d - \text{codim}(X), \quad \text{a.s.}$$

That is, in the best of circumstances,

$$\dim(X) + \text{codim}(X) = \text{topological dimension.}$$

A note of warning: if X is not compact, $\dim(X)$ can be defined by $\sup_{n \geq 1} \dim(\overline{X \cap [-n, n]^d})$.

The proof depends on the following result that can be found in the works of Yuval Peres '95, but with percolation proofs.

Lemma 4.4 (Peres' lemma) *For each $\beta \in (0, d)$, there exists a random set Λ_β , whose codimension is β . Moreover, $\dim(\Lambda_\beta) = d - \beta$, almost surely.*

Proof Let $\Lambda_\beta = Z(\mathbb{R}_+^N)$, where Z is an (N, d) -additive stable process. The result follows from Corollary 4.2 and Theorem 3.2. \square

Proof of Theorem 4.3 By localization, we may assume that X is a.s. compact. Let $\Lambda_\beta = \cup_{i=1}^{\infty} \Lambda_\beta^i$, where $\Lambda_\beta^1, \Lambda_\beta^2, \dots$ are iid copies of the sets in Peres' lemma, and are all totally independent of our random set X . Then, by Peres' lemma and by the lemma of Borel–Cantelli,

$$\mathbb{P}\{\Lambda_\beta \cap X \neq \emptyset \mid X\} = \begin{cases} 0, & \text{on } \{\dim(X) < \beta\} \\ 1, & \text{on } \{\dim(X) > \beta\} \end{cases}.$$

On the other hand, by the very definition of codimension,

$$\mathbb{P}\{\Lambda_\beta \cap X \neq \emptyset \mid \Lambda_\beta\} = \begin{cases} 0, & \text{if } \text{codim}(X) > d - \beta = \dim(\Lambda_\beta) \\ > 0, & \text{if } \text{codim}(X) < d - \beta \end{cases}.$$

Take expectations of the last two displays to see that for any $\beta \in (0, d)$,

$$\begin{aligned} \text{codim}(X) < d - \beta &\implies \dim(X) \geq \beta, \text{ a.s.} \\ \text{codim}(X) > d - \beta &\implies \dim(X) \leq \beta, \text{ a.s.} \end{aligned}$$

This easily proves our theorem.

□

Lecture 10

Brownian Sheet and Kahane's Problem

The (N, d) -Brownian sheet $B = \{B_t; t \in \mathbb{R}_+^N\}$ is an \mathbb{R}^d -valued, N -parameter Gaussian process with i.i.d. coordinate processes, B^1, \dots, B^d , each of which has the covariance:

$$\mathbb{E}\{B_s^1 B_t^1\} = \prod_{\ell=1}^N (s_\ell \wedge t_\ell), \quad \forall s, t \in \mathbb{R}_+^N.$$

For instance, consider $N = 2$ and write $B_{u,v}$ for the sheet. Then,

- for each fixed u , $v \mapsto u^{-\frac{1}{2}} B_{u,v}$ is a Brownian motion;
- for each fixed v , $u \mapsto v^{-\frac{1}{2}} B_{u,v}$ is a Brownian motion;
- for each $c > 0$ fixed, the process $B_{ce^{-v}, ev}$ has d i.i.d. coordinates each of which has covariance,

$$\mathbb{E}\{B_{ce^{-v}, ev}^1 B_{ce^{-u}, eu}^1\} = c^2 \exp(-|u - v|).$$

That is, $v \mapsto B_{ce^{-v}, ev}$ is an Ornstein–Uhlenbeck process.

You can also find all manners of time-changes of Brownian motion within Brownian sheet.

1 Local Structure and Potential Theory

1.1 Independent Increments

If $t \succ s$, $B_t - B_s$ is independent of $\{B_u; u \preccurlyeq s\}$. Since all is Gaussian, we check this by computing covariances, all the time assuming that $d = 1$, viz., for all $u \preccurlyeq s$,

$$\mathbb{E}\{(B_t - B_s)B_u\} = \prod_{\ell=1}^N (t_\ell \wedge u_\ell) - \prod_{i=1}^N (s_\ell \wedge u_\ell) = 0.$$

1.2 Incremental Laws

Whenever $t \succcurlyeq s$, $B_t - B_s$ is a vector of d i.i.d. centered Gaussians with variance

$$\mathbb{E}\{(B_t^1 - B_s^1)^2\} = \prod_{j=1}^N t_j + \prod_{k=1}^N s_k - 2 \prod_{i=1}^N (t_i \wedge s_i) = \prod_{j=1}^N t_j - \prod_{k=1}^N s_k.$$

In particular, if s is fixed, the above, for all $t \succcurlyeq s$ that are close to s , is $\simeq \sum_{j=1}^k |t_j - s_j|$. This is done by Taylor expansions. In particular, locally, $B_t - B_s$ has the same law as (N, d) additive Brownian motion.

Armed with the above estimate, one can then prove

Theorem 1.1 (Kh and Shi) *If E is a compact set in \mathbb{R}^d , $\mathbb{P}\{B(\mathbb{R}_+^N) \cap E \neq \emptyset\} > 0$ iff $\text{Cap}_{d-2N}(E) > 0$. In fact, for any $M > 0$, there exists c_1 and c_2 , such that for all compact $E \subset [-M, M]^d$,*

$$c_1 \text{Cap}_{d-2N}(E) \leq \mathbb{P}\{B[1, 2]^N \cap E \neq \emptyset\} \leq c_2 \text{Cap}_{d-2N}(E).$$

An immediate consequence of this is that $\text{codim} B(\mathbb{R}_+^N) = d - 2N$. Thus, essentially by Theorem 4.3 of Lecture 9, we have

Corollary 1.2 *With probability one,*

$$\dim B(\mathbb{R}_+^N) = 2N \wedge d.$$

(Essentially refers to the fact that the $2N \geq d$ case needs to be handled separately, but in the latter cases, it is not hard to show directly that the dimension is d .)

2 Kahane's Problem

We come to the last portion of these lectures, which is on a class of problems that I call Kahane's problem, due to the work of J.-P. Kahane in this area.

Kahane's problem for a random field X is: "when does $X(E)$ have positive Lebesgue's measure?" I will work the details out for Brownian motion, where things are easier. The problem for the Brownian sheet was partly solved by Kahane (cf. his '86 book) and completely solved by Kh. '99 in case $N = 2$. Recent work of Kh. and Xiao '01 has completed the solution to Kahane's problem and a class of related problems, and we hope to write this up at some point. Here is the story for Brownian motion, where we work things out more or less completely. The story for Brownian sheet is more difficult, and I will say some words about

the details later.

Theorem 2.1 (Kahane; Hawkes) *If B denotes Brownian motion in \mathbb{R} , and if $E \subset \mathbb{R}_+$ is compact, then*

$$\mathbb{E}\{|B(E)|\} > 0 \iff \text{Cap}_{\frac{d}{2}}(E) > 0.$$

You can interpret this as a statement about hitting probabilities for the level sets of Brownian motion, viz.,

$$\int_{\mathbb{R}^d} \mathbb{P}\{B^{-1}\{a\} \cap E \neq \emptyset\} da > 0 \iff \text{Cap}_{\frac{d}{2}}(E) > 0.$$

I will prove the following for Brownian motion. It clearly implies the above theorem upon integration.

Theorem 2.2 *Suppose $E \subset [1, 2]$ is compact, and fix $M > 0$. Then, there exists c_1 and c_2 such that for all $|a| \leq M$,*

$$c_1 \text{Cap}_{\frac{1}{2}}(E) \leq \mathbb{P}\{a \in B(E)\} \leq c_2 \text{Cap}_{\frac{1}{2}}(E).$$

Proof Without loss of any generality, we may and will assume that $E \subseteq [0, 1]$.

For any $\mu \in \mathcal{P}(E)$ and for all $a \in \mathbb{R}$, define

$$J_\varepsilon^a(\mu) = (2\varepsilon)^{-1} \int_0^\infty \mathbf{1}_{\{|B_s - a| \leq \varepsilon\}} \mu(ds).$$

Then, for every $M > 0$, there exists c such that

$$\begin{aligned} \inf_{\varepsilon \in (0,1)} \inf_{a \in [-M,M]} \mathbb{E}\{J_\varepsilon^a(\mu)\} &\geq c, \text{ and} \\ \sup_{a \in \mathbb{R}} \sup_{\varepsilon \in (0,1)} \mathbb{E}\{|J_\varepsilon^a(\mu)|^2\} &\leq \text{Energy}_{\frac{d}{2}}(\mu). \end{aligned} \tag{2.1}$$

Now, we apply Paley–Zygmund inequality:

$$\begin{aligned} \mathbb{P}\{a \in B(E)\} &\geq \mathbb{P}\{J_\varepsilon^a(\mu) > 0\} \\ &\geq \frac{|\mathbb{E}\{J_\varepsilon^a(\mu)\}|^2}{\mathbb{E}\{|J_\varepsilon^a(\mu)|^2\}} \\ &\geq \frac{c}{\text{Energy}_{\frac{d}{2}}(\mu)}. \end{aligned}$$

Since this holds for all $\mu \in \mathcal{P}(E)$, we obtain the desired lower bound.

Let $\{\mathcal{F}_t\}_{t \geq 0}$ denote the filtration of B and consider the martingale

$$M_t^{a,\varepsilon}(\mu) = \mathbb{E}\{J_\varepsilon^a(\mu) \mid \mathcal{F}_t\}, \quad \forall t \geq 0.$$

Clearly,

$$\begin{aligned} M_t^{a,\varepsilon}(\mu) &\geq (2\varepsilon)^{-1} \int_{s \geq t} \mathbb{P}\{|B_s - a| \leq \varepsilon \mid \mathcal{F}_t\} \mu(ds) \cdot \mathbf{1}_{\{|B_t - a| \leq \frac{\varepsilon}{2}\}} \\ &\geq (2\varepsilon)^{-1} \int_{s \geq t} \mathbb{P}\{|B_{s-t}| \leq \frac{1}{2}\varepsilon\} \mu(ds) \cdot \mathbf{1}_{\{|B_t - a| \leq \frac{\varepsilon}{2}\}} \\ &\geq c\varepsilon^{-1} \int_{s \geq t} \left(\frac{\varepsilon}{\sqrt{s-t}} \wedge 1 \right) \mu(ds) \cdot \mathbf{1}_{\{|B_t - a| \leq \frac{\varepsilon}{2}\}} \end{aligned}$$

Let $\sigma_\varepsilon = \inf\{s \in E : |B_s - a| \leq \frac{1}{2}\varepsilon\}$. This is a stopping time and on $\{\sigma_\varepsilon < \infty\}$,

$$M_{\sigma_\varepsilon}^{a,\varepsilon}(\mu) \geq c\varepsilon^{-1} \int_{s \geq \sigma_\varepsilon} \left[\frac{\varepsilon}{\sqrt{s - \sigma_\varepsilon}} \wedge 1 \right] \mu(ds),$$

since all bounded Brownian martingales are continuous. Now, we choose μ carefully: **WLOG** $\mathbb{P}\{\sigma_\varepsilon < \infty\} > 0$ which implies that $\mu_\varepsilon \in \mathcal{P}(E)$, where

$$\mu_\varepsilon(\bullet) = \mathbb{P}\{\sigma_\varepsilon \in \bullet \mid \sigma_\varepsilon < \infty\}.$$

Thus, by the optional stopping theorem,

$$\begin{aligned} 1 &\geq \mathbb{E}\{M_{\sigma_\varepsilon}^{a,\varepsilon}(\mu_\varepsilon); \sigma_\varepsilon < \infty\} \\ &\geq c\varepsilon^{-1} \iint_{s \geq t} \left(\frac{\varepsilon}{\sqrt{s-t}} \wedge 1 \right) \mu_\varepsilon(ds) \mu_\varepsilon(dt) \cdot \mathbb{P}\{\sigma_\varepsilon < \infty\} \\ &\geq \frac{c}{2}\varepsilon^{-1} \iint \left(\frac{\varepsilon}{\sqrt{s-t}} \wedge 1 \right) \mu_\varepsilon(ds) \mu_\varepsilon(dt) \cdot \mathbb{P}\{\sigma_\varepsilon < \infty\} \\ &= \frac{c}{2} \iint \left(\frac{1}{\sqrt{s-t}} \wedge \frac{1}{\varepsilon} \right) \mu_\varepsilon(ds) \mu_\varepsilon(dt) \cdot \mathbb{P}\{\sigma_\varepsilon < \infty\}. \end{aligned}$$

Fix $\delta_0 > 0$ and from the above deduce that for all ε small,

$$1 \geq \frac{c}{2} \iint_{|s-t| \geq \delta_0} |s-t|^{-\frac{1}{2}} \mu_\varepsilon(ds) \mu_\varepsilon(dt) \cdot \mathbb{P}\{\inf_{t \in E} |B_t - s| \leq \frac{1}{2}\varepsilon\}.$$

Let $\varepsilon \rightarrow 0$, and invoke Prohorov's theorem to get $\mu \in \mathcal{P}(E)$ such that

$$\mathbb{P}\{a \in B(E)\} \leq \frac{2}{c} \left[\iint_{|s-t| \geq \delta_0} |s-t|^{-\frac{1}{2}} \mu(ds) \mu(dt) \right]^{-1}.$$

Let $\delta_0 \downarrow 0$ to finish. □

To prove the general result, one needs the properties of the process B around a time point t . There are 2^N different notions of 'around', one for each quadrant centered at t , and this leads to 2^N different N -parameter martingales, each of which is a martingale with respect to a commuting filtration, but each filtration is indeed a filtration with respect to a different partial order. The details are complicated enough for $N = 2$ and can be found in my paper in the *Transactions of the AMS* (1999). When $N > 2$, the details are more complicated still and will be written up in the future. I will end with a related

CONJECTURE: Suppose X is an (N, d) symmetric stable sheet with index $\alpha \in (0, 2]$ (see below.) Then, for any compact $E \subset \mathbb{R}_+^N$, $\mathbb{E}\{|X(E)|\} > 0$ iff $\text{Cap}_{\frac{d}{\alpha}}(E) > 0$.

At the moment, this seems entirely out of the reach of the existing theory, but the analogous result for additive stable processes, and much more, holds (joint work with Xiao—will write up later.)

To finish: $\{X_t; t \in \mathbb{R}_+^N\}$ is an (N, d) symmetric stable sheet if it has i.i.d. coordinates and the first coordinate has the representation $X_t^1 = \int \mathbf{1}_{\{0 \preccurlyeq s \preccurlyeq t\}} \mathbb{X}(ds)$, where \mathbb{X} is a totally scattered random measure such that for every nonrandom measurable $A \subset \mathbb{R}_+^N$, $\mathbb{E}\{\exp[i\xi \mathbb{X}(A)]\} = \exp(-\frac{1}{2}|A| \|\xi\|^\alpha)$. (Scattered means that for nonrandom measurable A and A' in \mathbb{R}_+^N , if $A \cap A' = \emptyset$, $\mathbb{X}(A)$ and $\mathbb{X}(A')$ are independent.)

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